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MECHANICS

OF SOLIDS

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MECHANICS OF SOLIDS **CHESOLIDS**

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Preface

Mechanics of Solids is an important course for all engineering students by which they develop analytical skill. In this course, laws of mechanics are applied to parts of bodies and skill is developed to get solution to engineering problems maintaining continuity of the parts.

The author has clearly explained theories involved and illustrated them by solving a number of engineering problems. Neat diagrams are drawn and solutions are given without skipping any step. SI units and standard notations as suggested by Indian Standard Code are used throughout. The author has made this book to suit the latest syllabus of Gujarat Technical University. **Example 18**

Mechanics of Solida is an important course for all engineering students which they develop analytical skill. In this course, laws of mechanics are applied to part solidois and ekill is developed to get solut

Author hopes, the students and teachers of Gujarat Technical University will receive this book whole-heartedly as most of the earlier books of the author have been received by the students and teachers all over India.

The suggestions and corrections, if any, are most welcome.

The author acknowledges the efforts of M/s. New Age International Publishers in bringing out this book in nice form. He also acknowledges the opportunity given by AICTE for associating him with B.U.B. Engineering College, Hubli.

—Author

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1

Introduction to Mechanics of Solids

The state of rest and the state of motion of the bodies under the action of different forces has engaged the attention of mathematicians and scientists for many centuries. The branch of physical science that deal with the state of rest or the state of motion of bodies is termed as **mechanics.** Starting from the analysis of rigid bodies under gravitational force and application of simple forces the mechanics has grown into the analysis of complex structures like multistorey buildings, aircrafts, space crafts and robotics under complex system of forces like dynamic forces, atmospheric forces and temperature forces. **Introduction**

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Archemedes (287–212 BC), Galileo (1564–1642), Sir Issac Newton (1642–1727) and Einstein (1878–1955) have contributed a lot to the development of mechanics. Contributions by Varignon, Euler, and D. Alemberts are also substantial. The mechanics developed by these researchers may be grouped as

- (*i*) Classical mechanics/Newtonian mechanics
- (*ii*) Relativistic mechanics
- (*iii*) Quantum mechanics/Wave mechanics.

Sir Issac Newton, the principal architect of mechanics, consolidated the philosophy and experimental findings developed around the state of rest and state of motion of the bodies and putforth them in the form of three laws of motion as well as the law of gravitation. The mechanics based on these laws is called **Classical mechanics or Newtonian mechanics.**

Albert Einstein proved that Newtonian mechanics fails to explain the behaviour of high speed (speed of light) bodies. He putfourth the theory of **Relativistic mechanics.**

Schrödinger (1887–1961) and Broglie (1892–1965) showed that Newtonian mechanics fails to explain the behaviour of particles when atomic distances are concerned. They putforth the theory of **Quantum mechanics.**

Engineers are keen to use the laws of mechanics to actual field problems. Application of laws of mechanics to field problems is termed as **Engineering mechanics.** For all the problems between atomic distances to high speed distances there are various engineering problems for which Newtonian mechanics has stood the test of time and hence is the mechanics used by engineers.

The various bodies on which engineers are interested to apply laws of mechanics may be classified as

(*i*) Solids and

(*ii*) Fluids.

The bodies which do not change their shape or size appreciably when the forces are applied are termed as **Solids** while the bodies which change their shape or size appreciably even when small forces are applied are termed as **Fluids.** Stone, steel, concrete etc. are the example of solids while water, gases are the examples of fluids.

In this book application of Newtonian mechanics to solids is dealt with.

1.1 BASIC TERMINOLOGIES IN MECHANICS

The following are the terms basic to the study of mechanics, which should be understood clearly.

Mass

The quantity of the matter possessed by a body is called mass. The mass of a body will not change unless the body is damaged and part of it is physically separated. If the body is taken out in a space craft, the mass will not change but its weight may change due to the change in gravitational force. The body may even become weightless when gravitational force vanishes but the mass remain the same.

Time

The time is the measure of succession of events. The successive event selected is the rotation of earth about its own axis and this is called a day. To have convenient units for various activities, a day is divided into 24 hours, an hour into 60 minutes and a minute into 60 seconds. Clocks are the instruments developed to measure time. To overcome difficulties due to irregularities in the earths rotation, the unit of time is taken as second which is defined as the duration of 9192631770 period of radiation of the cesium-133 atom. In this book application of Newtonian mechanics to solids is dealt with.

1.1 **BASIC TERMINOLOGIES IN MECHANICS**

The following are the terms basic to the study of mechanics, which should be understoot

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Space

The geometric region in which study of body is involved is called space. A point in the space may be referred with respect to a predetermined point by a set of linear and angular measurements. The reference point is called the *origin* and the set of measurements as *coordinates*. If the coordinates involved are only in mutually perpendicular directions, they are known as *cartesian coordination*. If the coordinates involve angles as well as the distances, it is termed as *Polar Coordinate System.*

Length

It is a concept to measure linear distances. The diameter of a cylinder may be 300 mm, the height of a building may be 15 m, the distance between two cities may be 400 km.

Actually metre is the unit of length. However depending upon the sizes involved micro, milli or kilo metre units are used for measurements. A metre is defined as length of the standard bar of platinum-iradium kept at the International Bureau of weights and measures. To overcome the difficulties of accessibility and reproduction now metre is defined as 1690763.73 wavelength of krypton-86 atom.

Continuum

A body consists of several matters. It is a well known fact that each particle can be subdivided into molecules, atoms and electrons. It is not possible to solve any engineering problem by treating a body as conglomeration of such discrete particles. The body is assumed to be a continuous distribution of matter. In other words the body is treated as continuum.

Rigid Body

A body is said to be rigid, if the relative positions of any two particles do not change under the action of the forces acting on it. In Fig. 1.1 (*a*), point *A* and *B* are the original positions in a body. After the application of forces F_1 , F_2 , F_3 , the body takes the position as shown in Fig. 1.1(*b*). *A*^{\prime} and *B*′ are the new positions of *A* and *B*. If the body is treated as rigid, the relative position of *A*′*B*′ and *AB* are the same *i.e*.

$$
A'B' = AB
$$

Many engineering problems can be solved by assuming bodies rigid

Particle

A particle may be defined as an object which has only mass and no size. Theoretically speaking such a body cannot exist. However in dealing with problems involving distances considerably larger compared to the size of the body, the body may be treated as a particle, without sacrificing accuracy.

For example:

- A bomber aeroplane is a particle for a gunner operating from the ground.
- A ship in mid sea is a particle in the study of its relative motion from a control tower.
- In the study of movement of the earth in celestial sphere, earth is treated as a particle.

Force

Force is an important term used in solid mechanics. Newton's first law states that everybody continues in its state of rest or of uniform motion in a straight line unless it is compelled by an external agency acting on it. This leads to the definition of force as '*force is an external agency which changes or tends to change the state of rest or uniform linear motion of the body*'.

 Magnitude of force is defined by Newton's second law. It states that the rate of change of momentum of a body is directly proportional to the impressed force and it takes place in the direction of the force acting on it. Noting that rate of change of velocity is acceleration, and the product of mass and velocity is momentum we can derive expression for the force as given below:

From Newton's second law of motion

Force ∞ rate of change of momentum

∝ rate of change of (mass × velocity)

Since mass do not change,

Force ∞ mass \times rate of change of velocity ∝ mass × acceleration $F \propto m \times a$...(1.1) $= k \times m \times a$

where F is the force, m is the mass and a is the acceleration and k is the constant of proportionality.

In all the systems, unit of force is so selected that the constant of the proportionality becomes unity. For example, in S.I. system, unit of force is Newton, which is defined as the force that is required to move one kilogram (kg) mass at an acceleration of 1 m/sec². From X a
 $F = mx \times a$

where *F* is the cone, mix the mass and a is the exceleration and *k* is the constant of proportionality

and all the systems, unit of force is so extended that the constant of the proportionality

tin

∴ One newton = 1 kg mass \times 1 m/sec² Thus $k = 1$ $F = m \times a$...(1.2)

However in MKS acceleration used is one gravitational acceleration (9.81 m/sec^2) on earth surface) and unit of force is defined as kg-wt.

Thus

It may be noted that in usage kg-wt is often called as kg only.

Characteristics of a Force

It may be noted that a force is completely specified only when the following four characteristics are specified

- Magnitude
- Point of application
- Line of action
- Direction.

In Fig. 1.2, *AB* is a ladder kept against a wall. At point *C*, a person weighing 600 N is standing. The force applied by the person on the ladder has the following characters:

- magnitude is 600 N
- the point of application is *C* which is at 2 m from *A* along the ladder
- the line of action is vertical
- the direction is downward.

It may be noted that in the figure

- magnitude is written near the arrow
- the line of arrow shows the line of application
- the arrow head shows the point of application
- the direction of arrow represents the direction of the force.

 \sim $\hat{\mathcal{E}}$

A

600 N

C

B

CONTROL COMMANDER

1.2 UNITS

Length (L), mass (M) and time (S) are the fundamental units used in mechanics. The units of all other quantities may be expressed in terms of these basic units. The three commonly used systems are

- Metre, Kilogram, Second (MKS)
- Centimetre, Gram, Second (CGS)
- Foot, Pound, Second (FPS).

The systems are named after the units used to define the fundamental quantities length, mass and time. Using these basic units, the units of other quantities can be found. For example in MKS the units for various quantities are

S.I. Units

Presently the whole world is in the process of switching over to SI-system of units. SI units stands for the System International *d*′ units or International System of units. As in MKS units in SI also the fundamental units are metre for length, kilogram for mass and second for time. The difference between MKS and SI system arises mainly in selecting the unit of force. In MKS unit of force is kg-wt while in SI units it is newton. As we have already seen one kg-wt is equal to 9.81 newtons.

The prefixes used in SI when quantities are too big or too small are shown in Table 1.1.

Table 1.1. Prefixes in SI Units

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1.3 SCALAR AND VECTOR QUANTITIES

Various quantities used in mechanics may be grouped into scalars and vectors. A quantity is said to be scalar, if it is completely defined by its magnitude alone. Examples of scalars are length, area, time and mass.

A quantity is said to be vector if it is completely defined only when its magnitude as well as direction are specified. The example of vectors are displacement, velocity, acceleration, momentum, force etc.

1.4 COMPOSITION AND RESOLUTION OF VECTORS

The process of finding a single vector which will have the same effect as a set of vectors acting on a body is known as composition of vectors. The resolution of vectors is exactly the opposite process of composition *i.e.,* it is the process of finding two or more vectors which will have the same effect as that of a vector acting on the body.

Parallelogram Law of Vectors

The parallelogram law of vectors enables us to determine the single vector called resultant vector which can replace the two vectors acting at a point with the same effect as that of the two vectors. This law was formulated based on exprimental results on a body subjected to two forces. This law can be applied not only to the forces but to any two vectors like velocities, acceleration, momentum etc. Though stevinces employed it in 1586, the credit of presenting it as a law goes to Varignon and Newton (1687). This law states that *if two forcer (vectors) acting simultaneously on a body at a point are represented in magnitude and directions by the two adjacent sides of a parallelogram, their resultant is represented in magnitude and direction by the diagonal of the parallelogram which passes thorough the point of intersection of the two sides representing the forces (vectors).*

In the Fig. 1.3, the force $F_1 = 4$ units and the force $F_2 = 3$ unit are acting on a body at a point *A*. To get the resultant of these forces, according to this law, construct the parallelogram *ABCD* such that *AB* is equal to 4 units to the linear scale and *AC* is equal to 3 units. Then according to this law, the diagonal *AD* represents the resultant in magnitude and direction. Thus the resultant of the forces F_1 and F_2 is equal to the units corresponding to *AD* in the direction α to F_1 .

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Triangle Law of Vectors

Referring to Fig. 1.3 (*b*), it can be observed that the resultant *AD* may be obtained by constructing the triangle *ABD*. Line *AB* is drawn to represent F_1 and *BD* to represent F_2 . Then *AD* should represent the resultant of F_1 and F_2 . Thus we have derived the triangle law of forces from the fundamental law of parallelogram. The *Triangle Law of Forces (vectors)* may be stated as *if two forces (vectors) acting on a body are represented one after another by the sides of a triangle, their resultant is represented by the closing side of the triangle taken from the first point to the last point.*

Polygon Law of Forces (Vectors)

If more than two forces (vectors) are acting on a body, two forces (vectors) at a line can be combined by the triangle law, and finally resultant of all forces (vectors) acting on the body may be obtained.

A system of four concurrent forces acting on a body are shown in Fig. 1.4. *AB* represents *F*¹ and *BC* represent F_2 . Hence according to triangle law of forces *AC* represents the resultant of F_1 and F_2 , say R_1 .

If *CD* is drawn to represent F_3 , then from the triangle law of forces *AD* represents the resultant of R_1 and F_3 . In other words, *AD* represents the resultant of F_1 , F_2 and F_3 . Let it be called as R_2 .

Similarly the logic can be extended to conclude that *AE* represents the resultant of F_1 , F_2 , F_3 and $F₄$. The resultant *R* is represented by the closing line of the polygon *ABCDE* in the direction form *A* to *E*. Thus we have derived the polygon law of the forces (vectors) and it may be stated as *if a number of concurrent forces (vectors) acting simultaneously on a body are represented in magnitude and direction by the sides of a polygon, taken in a order, then the resultant is represented in magnitude and direction by the closing side of the polygon, taken from the first point to the last point.*

Analytical Method of Composition of Two Vectors

Parallelogram law, triangle law and polygonal law of vectors can be used to find the resultant graphically. This method gives a clear picture of the work being carried out. However the main disadvantage is that it needs drawing aids like pencil, scale, drawing sheets. Hence there is need for analytical method.

Consider the two forces F_1 and F_2 acting on a particle as shown in Fig 1.5(*a*). Let the angle between the two forces be θ. If parallelogram *ABCD* is drawn as shown in Fig. 1.5(*b*) with *AB* respresenting F_1 and AD representing F_2 to some scale, according to parallelogram law of forces *AC* represents the resultant *R*. Drop perpendicular *CE* to *AB*.

Fig. 1.5

The resultant *R* of F_1 and F_2 is given by

$$
R = AC = \sqrt{AE^2 + CE^2} = \sqrt{(AB + BE)^2 + CE^2}
$$

$$
\mathbf{F}_{2}
$$
\n
$$
\mathbf{F}_{3}
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$$
\mathbf{F}_{4}
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$$
\mathbf{F}_{5}
$$
\n(a)
\n**Fig. 1.5**\n
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\mathbf{F}_{6}
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\mathbf{F}_{8}
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\mathbf{F}_{9}
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Since, $\sin^2 \theta + \cos^2 \theta = 1$.

The inclination of resultant to the direction of F_1 is given by α , where

$$
\tan \alpha = \frac{CE}{AE} = \frac{CE}{AB + BE} = \frac{F_2 \sin \theta}{F_1 + F_2 \cos \theta}
$$

Hence $\alpha = \tan^{-1} \frac{F_2}{\alpha}$ $F_1 + F_2$ 2 $_1$ + $_2$ sin cos θ + $F_2 \cos \theta$

Particular cases:

1. When
$$
\theta = 90^{\circ}
$$
 [Ref. Fig. 1.6a],
\n
$$
R = \sqrt{F_1^2 + F_2^2}
$$
\n2. When $\theta = 0^{\circ}$ [Ref. Fig. 1.6b],
\n
$$
R = \sqrt{F_1^2 + 2F_1F_2^2 + F_2^2} = F_1 + F_2
$$
\n3. When $\theta = 180^{\circ}$ [Ref. Fig. 1.6c],
\n
$$
R = \sqrt{F_1^2 - 2F_1F_2 + F_2^2} = F_1 - F_2
$$

Resolution of Vectors

Since the resolution of vectors is exactly opposite process of composition of vectors, exactly the opposite process of composition can be employed to get the resolved components of a given force.

In Fig. 1.7(*a*), the given force *F* is resolved into two components making angles α and β with *F*. In Fig. 1.7(*b*) the force *F* is resolved into its rectangular components F_x and F_y . In Fig. 1.7(*c*), the force *F* is resolved into its four components F_1 , F_2 , F_3 and F_4 . It may be noted that all component forces act at the same point as the given force. Resolution

of forces into its rectangular components is more useful in solving the problems in mechanics. In this case, if the force *F* makes angle θ with *x*-axis, from Fig. 1.7(*a*), it is clear that

 $F_x = F \cos \theta$ and $F_y = F \sin \theta$.

Example 1.1. A boat is rowed at a velocity of 20 km/hour across a river. The velocity of stream is 8 km/hour. Determine the resultant velocity of the boat.

Solution: Taking downstream direction as *x* and direction across the river as *y*, it is given that

$$
V_x = 8 \text{ km/hour}
$$
\n
$$
V_y = 20 \text{ km/hour}
$$
\n
$$
\mathbf{v} = \sqrt{8^2 + 20^2} = 21.54 \text{ km/hour}
$$
\n
$$
\mathbf{\alpha} = \tan^{-1} \frac{V_y}{V_x} = \tan^{-1} \frac{20}{8} = 68.20^\circ, \text{ as shown in Fig. 1.8}
$$
\n
$$
\sum_{\substack{Q \text{ odd} \\ Q \text{ odd}}}^{\infty} \frac{V_x = 8 \text{ km/hour}}{S_x^2}
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\sum_{\substack{Q \text{ odd} \\ Q \text{ odd}}}^{\infty} \frac{V_x = 8 \text{ km/hour}}{S_x^2}
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Example. 1.2. The guy wire of the electrical pole shown in Fig. 1.9(a) makes 60° to the horizontal and is carrying a force of 60 kN. Find the horizontal and vertical components of the force.

Solution: Figure 1.9(*b*) shows the resolution of force $F = 20$ kN into its components in horizontal and vertical components. From the figure it is clear that

> $F_x = F \cos 60^\circ = 20 \cos 60^\circ = 10 \text{ kN (to the left)}$ $F_y = F \sin 60^\circ = 20 \sin 60^\circ = 17.32 \text{ kN (downward)}$

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Example 1.3. A black weighing W = 10 kN is resting on an inclined plane as shown in Fig. 1.10(a). Determine its components normal to and parallel to the inclined plane.

Solution: The plane makes an angle of 20° to the horizontal. Hence the normal to the plane makes an angles of 70° to the horizontal *i.e.*, 20° to the vertical [Ref. Fig. 1.10(*b*)]. If *AB* represents the given force *W* to some scale, *AC* represents its component normal to the plane and *CB* represents its component parallel to the plane.

Thus from ∆ *ABC*, Component normal to the plane $= AC$

From the above example, the following points may be noted:

1. Imagine that the arrow drawn represents the given force to some scale.

2. Travel from the tail to head of arrow in the direction of the coordinates selected.

- 3. Then the direction of travel gives the direction of the component of vector.
- 4. From the triangle of vector, the magnitudes of components can be calculated.

Example 1.4. The resultant of two forces, one of which is double the other is 260 N. If the direction of the larger force is reversed and the other remain unaltered, the magnitude of the resultant reduces to 180 N. Determine the magnitude of the forces and the angle between the forces. **Solution:** Let the magnitude of the smaller force be *F*. Hence the magnitude of the larger force is 2*F*.

Thus $F_1 = F$ and $F_2 = 2F$

Let θ be the angle between the two forces.

∴ From the condition 1, we get

$$
R = \sqrt{F_1^2 + 2F_1F_2 \cos \theta + F_2^2} = 260
$$

i.e.,

$$
F^2 + 2F (2F) \cos \theta + (2F)^2 = 260^2
$$

$$
5F^2 + 4F^2 \cos \theta = 67600
$$
...(i)

From condition 2, we get $F_1^2 + 2F_1F_2 \cos(180 + \theta) + F_2^2 = 180$ $F^2 - 2F(2F) \cos \theta + (2F)^2 = 32400$...(*ii*) Adding equation (*i*) and (*ii*), we get $10F^2 = 100000$ ∴ $F = 100 N$ Hence **F**₁ = **F** = 100 N; **F**₂ = 2**F** = 200 N Substituting the values of F_1 and F_2 in eqn (*i*), we get, $5(100)^2 + 4(100)^2 \cos \theta = 67600$ $∴ \cos θ = 0.44$ or $\theta = 63.9^\circ$ *Example 1.5. Two forces* F_1 *and* F_2 *are acting at point A as shown in Fig. 1.11. The angle between the two forces is 50°. It is found that the resultant R is 500 N and makes angles 20°* with the force F_1 as shown in the figure. Determine the forces F_1 and F_2 . *Solution:* Let ∆*ABC* be the triangle of forces drawn to some scale. In this $\angle BAC = \alpha = 20^{\circ}$ $\angle ABC = 180 - 50 = 130^{\circ}$ ∴ ∠*ACB* = 180 – (20 + 130) = 30° Applying sine rule to ∆ *ABC*, we get *AB BC* $\frac{12}{\sin 30^{\circ}} = \frac{26}{\sin 20^{\circ}} = \frac{560}{\sin 130^{\circ}}$ 500 $\sin 130^\circ$ $AB = 326.35$ N and $BC = 223.24 \text{ N}$. Thus $F_1 = AB = 326.35 N$ and **F₂** = BC = 223.24 N 50° $F₂$ $A \t B$ C $\sqrt{\frac{50^{\circ}}{\alpha}} = 20^{\circ}$ R **Fig. 1.11** Adding equation (i) and (ii), we get
 $F_1 = V = 100 \text{ N}$
 $F_2 = 2F = 200 \text{ N}$

Hence
 $F_1 = 100 \text{ N}$; $F_2 = 2F = 200 \text{ N}$

Substituting the values of F_1 and F_2 in eqn (i), we get
 $S(100)^2 + 4(100)^2 \cos \theta = 67600$

or

Example 1.6. The resultant of two forces $F_1 = 400$ *N and F2 = 260 N acting at point A is 520 N. Determine the angle between the two forces and the angle between the resultant and force F1.*

Solution: Let *ABC* be the triangle of forces as shown in Fig. 1.12. θ be the angle between F_1 and F_2 , and α be the angle between resultant and F_1

Using the relation

$$
R = \sqrt{F_1^2 + F_2^2 + 2F_1F_2 \cos \theta},
$$

we get,

$$
520^2 = 400^2 + 260^2 + 2 \times 400 \times 260 \times \cos \theta
$$

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Example 1.7. Fig. 1.13 shows a particular position of 200 mm connecting rod AB and 80 mm long crank BC. At this position, the connecting rod of the engine experience a force of 3000 N on the crank pin at B. Find its

- *(a) horizontal and vertical component*
- *(b) component along BC and normal to it.*

Solution: The force of 3000 N acts along line *AB*. Let *AB* make angle α with horizontal. Then, obviously 200 sin $\alpha = 80 \sin 60^\circ$

∴ $\alpha = 20.268^\circ$

Referring to Fig. 1.13(*b*), we get **Horizontal component** = 3000 cos 20.268° = **2814.2 N Vertical component** $= 3000 \sin 20.268^\circ = 1039.2 \text{ N}$ Components along and normal to crank: The force makes angle $\alpha + 60^{\circ} = 20.268 + 60 = 80.268^{\circ}$ with crank. ∴ **Component along crank** = 3000 cos 80.268° = **507.1 N Component normal to crank** = 3000 sin 80.268° = **2956.8 N**

IMPORTANT FORMULAE

1. Resultant of two vectors can be obtained by solving the triangle of forces.

2. If V_1 and V_2 are the two vectors at angle 'θ' between them, then the resultant is

$$
R = \sqrt{V_1^2 + V_2^2 + 2V_1V_2 \cos \theta}
$$

and acts at ' α ' to V_1 vector, where

$$
\tan \alpha = \frac{V_2 \sin \alpha}{V_1 + V_2 \cos \alpha}
$$

Vectors may be forces, velocities, momentum etc.

3. If a force makes angle θ with *x*-axis, then its components are

$$
F_x = F \cos \theta
$$

$$
F_y = F \sin \theta.
$$

4. If a body weighing *W* rests on an inclined plane, its components normal to and parallel to the plane are

> $F_n = W \cos \theta$, a thrust on the plane. $F_t = W \sin \theta$, down the plane.

THEORY QUESTIONS

- **1.** Explain the following terms:
	-

(*i*) Space (*ii*) Continuum

-
- (*iii*) Particle (*iv*) Rigid body.
- **2.** Explain the term 'Force' and list its characteristics.
- **3.** Distinguish between
	- (*i*) MKS and SI units
	- (*ii*) Scalars and vectors.
- **4.** State and explain parallelogram law of vectors.
- **5.** State parallelogram law of vector and derive triangle and polygonal law of vectors.

PROBLEMS FOR EXERCISE

- **1.** The resultant of two forces one of which is 3 times the other is 300 N. When the direction of smaller force is reversed, the resultant is 200 N. Determine the two forces and the angle between them. **[Ans.** F₁ = 80.6 N, F₂ = 241.8 N, θ = 50.13°]
- **2.** A rocket is released from a fighter plane at an angle upward 20° to the vertical with an acceleration of 8 m/sec². The gravitational acceleration is 9.1 m/sec² downward. Determine the instantaneous acceleration of the rocket when it was fired. [Ans. 9.849 m/sec², $\theta = 49.75^{\circ}$ to vertical] Vectors may be forces, velocities, momentum etc.

3. If a force makes angle θ with zaris, then its components are
 $F_x = F$ os θ

4. If a hody weighing W rests on an inclined plane, its components normal to and parall

2

Fundamentals of Statics

In this chapter principles of statics is explained and their applications to concurrent and non-concurrent force system in plane is illustrated by solving several engineering problems.

2.1 PRINCIPLES OF STATICS

The statics is based on the following principles of mechanics:

- 1. Newton's laws of mechanics
- 2. Law of transmissibility
- 3. Parallelogram law of forces
- 4. Principles of physical independence
- 5. Principles of superposition.

2.1.1 Newton's Laws of Mechanics

As already discussed in first chapter, Newton's first law gave definition of the force and second law gave basis for quantifying the force. There are two more Newton's laws:

- *a*. Newton's Third Law
- *b*. Newton's Law of Gravitation

These laws are explained in this article.

(a) Newton's Third Law

It states that for every action there is an equal and opposite reaction. Consider the two bodies in contact with each other. Let one body apply a force *F* on another. According to this law the second body develops a reactive force *R* which is equal in magnitude to force *F* and acts in the line same as *F* but in the opposite direction. Figure 2.1 shows the action of a ball on the floor and the reaction of floor to this action. In Fig. 2.2 the action of a ladder on the wall and the floor and the reactions from the wall and the floor are shown.

(b) Newton's Law of Gravitation

It states that everybody attracts the other body. The force of attraction between any two bodies is directly proportional to their masses and inversely proportional to the square of the distance between them. Thus the force of attraction between the bodies of mass m_1 and mass m_2 at distance '*d*' between them as shown in Fig. 2.3 is

$$
F = G \frac{m_1 m_2}{d^2} \tag{2.1}
$$

 $m₂$

where *G* is the constant of proportionality and is known as constant of gravitation.

1 $m₁$

From eqn. 2.1,

$$
G = \frac{Fd^2}{m_1m_2}
$$

Hence unit of
$$
G = \frac{N \times m^2}{kg \times kg} = \text{Nm}^2/\text{kg}^2
$$

It has been proved by experiments that the value of $G = 6.673 \times 10^{-11} \text{ Nm}^2/\text{kg}^2$. Thus if two bodies one of mass 10 kg and the other of 5 kg are at a distance of 1 m, they exert a force

$$
F = \frac{6.673 \times 10^{-11} \times 10 \times 5}{1^2} = 33.365 \times 10^{-10} \text{ N}
$$

 F $F \leftarrow$ 2

F

d **Fig. 2.3**

on each other.

Similarly 1 kg-mass on earth surface experiences a force of

$$
F = \frac{6.673 \times 10^{-11} \times 1 \times 5.96504 \times 10^{24}}{(6371 \times 10^3)^2} = 9.80665 \text{ N}
$$

Since, mass of earth = 5.96504×10^{24} kg

and radius of earth = 6371×10^3 m.

This force of attraction is always directed towards the centre of earth.

In common usage the force exerted by a earth on a body is known as weight of the body. Thus weight of 1 kg-mass on/near earth surface is 9.80665 N, which is approximated as 9.81 N for all practical problems. Compared to this force the force exerted by two bodies on each other is negligible. Thus in statics:

a. Weight of a body = mg

b. Its direction is towards the centre of the earth, in other words, vertically downward.

c. The force of attraction between the other two objects on the earth is negligible.

2.1.2 Law of Transmissibility

According to this law *the state of rest or motion of the rigid body is unaltered*, *if a force acting on the body is replaced by another force of the same magnitude and direction but acting anywhere on the body along the line of action of the replaced force.*

Let *F* be the force acting on a rigid body at point *A* as shown in Fig. 2.4. According to this law, this force has the same effect on the state of body as the force *F* applied at point *B*, where *AB* is in the line of force *F*.

In using law of transmissibility it should be carefully noted that it is applicable only if the body can be treated as rigid. Hence if we are interested in the study of internal forces developed in a body, the deformation of body is to be considered and hence this law cannot be applied in such studies.

2.1.3 Parallelogram Law of Forces

This has been already explained in chapter 1 along with the derived laws *i.e.*, triangle and polygonal law.

2.1.4 Principles of Physical Independence of Forces

It states that *the action of a force on a body is not affected by the action of any other force on the body.*

2.1.5 Principles of Superposition of Forces

It states that *the net effect of a system of forces on a body is same as the combined* of individual forces acting on the body. Since a system of forces in equilibrium do not have any effect on a rigid body this principle is stated in the following form also: 'The effect of a given system of forces on a rigid body is not changed by adding or subtracting another system of forces in equilibrium.'

2.2 SYSTEM OF FORCES

When several forces of different magnitude and direction act upon a body, they constitute a *system of forces.* If all the forces in a system lie in a single plane, it is called a *coplanar force system.* If the line of action of all the forces in a system pass through a single point it is called a *concurrent force system.* In a system of *parallel forces* all the forces are parallel to each other. If the line of action of all forces lie along a single line then it is called a *collinear force system.* Various system of forces, their characteristics and examples are given in Table 2.1 below.

Table 2.1 System of Forces

2.3 MOMENT OF A FORCE

Moment of a force about a point is the measure of its rotational effect. *Moment* is defined as the product of the magnitude of the force and the perpendicular distance of the point from the line of action of the force. The point about which the moment is considered is called *moment centre* and the perpendicular distance of the point from the line of action of the force is called *moment arm.* Referring to Fig. 2.5, if d_1 is the perpendicular distance of point 1 from the line of action of force *F*, the moment of *F* about point 1 is given by

$$
M_1 = F d_1
$$
...(2.2)
Similarly, moment about point 2 is given by

$$
M_2 = F d_2
$$
...(2.3)

If the moment centre 3 lies on the line of action of the force *F*, the moment arm is zero and hence,

$$
M_3 = F \times 0 = 0 \tag{2.4}
$$

Thus, it may be noted that if a point lie on the line of action of a force, the moment of the force about that point is zero.

The moment of a force has got direction also. In Fig. 2.5 it may be noted that M_1 is *clockwise* and M_2 is *anticlockwise*. To find the direction of the moment, imagine that the line of action of the force is connected to the point by a rigid rod pinned at the point and is free to move around the point. The direction of the rotation indicates the direction of the moment. **EVALUATE: EVALUATE: EVALUATE:** $M_1 = F A_1$

Similarly, moment about point 2 is given by

If the moment centre 3 lies on the line of action of the force *F*, the moment ann is

hence,

Thus, it may be nuted that if a p

If the force is taken in newton unit and the distance in millimetre, the unit of moment will be N-mm. Commonly used units of moment in engineering are kN-m, N-m, kN-mm and N-mm.

2.4 VARIGNON'S THEOREM

French mathematician Varignon (1654–1722) gave the following theorem which is also known as principle of moments:

The algebraic sum of the moments of a system of coplanar forces about a moment centre in their plane is equal to the moment of their resultant force about the same moment centre.

Proof: Referring to Fig. 2.6 let *R* be the resultant of forces F_1 and F_2 and *B* the moment centre. Let *d*, d_1 and d_2 be the moment arms of the forces, *R*, F_1 and F_2 , respectively from the moment centre *B*. Then in this case, we have to prove that:

$$
Rd = F_1 \, d_1 + F_2 \, d_2
$$

Join *AB* and consider it as *y* axis and draw *x* axis at right angles to it at *A* [Fig. 2.6(*b*)]. Denoting by θ the angle that *R* makes with *x* axis and noting that the same angle is formed by perpendicular to *R* at *B* with AB_1 , we can write:

$$
Rd = R \times AB \cos\theta
$$

= AB \times (R \cos\theta)
= AB \times R_x ...(a)

where R_x denotes the component of R in x direction.

Fig. 2.5

Fig. 2.6

F.

Similarly, if F_{1x} and F_{2x} are the components of F_1 and F_2 , in *x* direction, respectively, then

and
$$
F_1 d_1 = AB F_{1x}
$$
 ...(b)
 $F_2 d_2 = AB F_{2x}$...(c)

(b)

x

$$
F_1 d_1 + F_2 d_2 = AB (F_{1x} + F_{2x})
$$

= AB × R_x ...(d)

From equation (*a*) and (*d*), we get

A

From Eqns. (*b*) and (*c*)

(a)

$$
Rd = F_1 d_1 + F_2 d_2
$$

If a system of forces consists of more than two forces, the above result can be extended as given below:

Let F_1 , F_2 , F_3 and F_4 be four concurrent forces and *R* be their resultant. Let d_1 , d_2 , d_3 , d_4 and *a* be the distances of line of action of forces F_1 , F_2 , F_3 , F_4 and *R*, respectively from the moment centre *O*, [Ref. Fig 2.7].

If R_1 is the resultant of F_1 and F_2 and its distance from *O* is a_1 , then applying Varignon's theorem:

$$
R_1 \, a_1 = F_1 \, d_1 + F_2 \, d_2
$$

If R_2 is the resultant of R_1 and F_3 (and hence of F_1 , F_2 and F_3) and its distance from *O* is a_2 , then applying Varignon's theorem:

$$
R_2 a_2 = R_1 a_1 + F_3 a_3
$$

= $F_1 d_1 + F_2 d_2 + F_3 d_3$

Fig. 2.7

Now considering R_2 and F_4 , we can write:

 $Ra = R_2 a_2 + F_4 d_4$

Since *R* is the resultant of R_2 and F_4 (i.e. F_1 , F_2 , F_3 and F_4). ∴ $Ra = F_1d_1 + F_2d_2 + F_3d_3 + F_4d_4$...(2.5)

Thus, the moment of the resultant of a number of forces about a moment centre is equal to the sum of the moments of its component forces about the same moment centre.

Example 2.1. Find the moment of 100 N force acting at B about point A as shown in Fig. 2.8.

Solution: 100 N force may be resolved into its horizontal components as 100 cos 60° and vertical component 100 sin 60°. From Varignon's theorem, moment of 100 N force about the point A is equal to sum of the moments of its components about A.

Taking clockwise moment as positive,

$$
M_A = 100 \cos 60^\circ \times 500 - 100 \sin 60^\circ \times 400
$$

= 25, 000 - 34, 641.02
= - 9641.02 N-mm
= 9641.016 N-mm Anticlockwise.

Example 2.2. What will be the y intercept of the 5000 N force if its moment about A is 8000 N-m in Fig. 2.9.

Solution: 5000 N force is shifted to a point *B* along its line of action (law of transmissibility) and it is resolved into its *x* and *y* components (F_x and F_y as shown in Fig. 2.9).

$$
F_x = 5000 \cos \theta = 5000 \times \frac{4}{5} = 4000 \text{ N}
$$

and

$$
F_y = 5000 \sin \theta = 5000 \times \frac{3}{5} = 3000 \text{ N}.
$$

By Varignon's theorem, moment of 5000 N force about *A* is equal to moment of its component forces about the same point.

 $8000 = 4000 \times y + 3000 \times 0$ ∴ *y* = **2 m.**

2.5 COUPLE

Two parallel forces equal in magnitude and opposite in direction and separated by a definite distance are said to form a couple. The sum of the forces forming a couple is zero, since they are equal and opposite, which means the translatory effect of the couple is zero.

An interesting property can be observed if we consider rotational effect of a couple about any point. Let the magnitude of the forces forming the couple be *F* and the perpendicular distance between the two forces be *d.* Consider the moment of the two forces constituting a couple about point 1 as shown in Fig. 2.10(*a*). Let the moment be M_1 then,

$$
M_1 = Fd_1 + Fd_2
$$

= $F (d_1 + d_2) = Fd$

Now, consider the moment of the forces about point 2 which is outside the two forces as shown in Fig. 2.10(*b*). Let M_2 be the moment.

Then,

Thus, moment of a couple about any point is the same. Now we can list the following characteristics of a couple:

– A couple consists of a pair of equal and opposite parallel forces which are separated by a definite distance;

– The translatory effect of a couple on the body is zero;

– The rotational effect (moment) of a couple about any point is a constant and it is equal to the product of the magnitude of the forces and the perpendicular distance between the two forces.

Since the only effect of a couple is a moment and this moment is the same about any point, the effect of a couple is unchanged if:

- The couple is rotated through any angle;
- The couple is shifted to any other position;
- The couple is replaced by another pair of forces whose rotational effect is the same.

2.6 TRANSFER OF A FORCE TO PARALLEL POSITION

It will be advantageous to resolve a force acting at a point on a body into a force acting at some other suitable point on the body and a couple. In Fig. 2.11(*a*) F is a force acting on a body at A .

Now it can be shown that *F* at *A* may be resolved into force *F* at *B* and a couple of magnitude $M = F \times d$, where *d* is the perpendicular distance of *B* from the line of action of *F* through *A*.

By applying equal and opposite forces *F* at *B* the system of forces is not disturbed. Hence the system of forces in Fig. 2.11(*b*) is the same as the system given in Fig. 2.11(*a*). Now the original force *F* at *A* and the opposite force *F* at *B* form a couple of magnitude *Fd*. The system in Fig. 2.11(*b*) can be replaced by the system shown in Fig. 2.11(*c*). Thus, the given force F at \overline{A} is replaced by a force *F* at *B* and a moment *Fd*.

2.7 COMPOSITION OF CONCURRENT COPLANAR FORCES

General Approach

In chapter 1, composition of concurrent forces by graphical method and the analytical method of composition of two force system has been discussed. In this article composition of concurrent coplanar forces is explained by a general analytic method.

Analytical method consists in finding the components of given forces in two mutually perpendicular directions and then combining them to get the resultant. Finding the component of a force is called *resolution of forces* and is exactly the opposite to the process of composition of forces. Finding the components of forces in two mutually perpendicularly directions is preferable. The following points associated with the analytical method of finding rectangular components may be noted:

- (*i*) Imagine that the arrow drawn to show force represents it to some scale
- (*ii*) Travel from tail to head of the arrow in the directions of coordinates

(*iii*) The direction of the travel gives the direction of component forces

(*iv*) From the triangle law of forces, the magnitude of the components can be calculated.

After finding the components of all the forces in the system in the two mutually perpendicular directions, the component in each direction are algebraically added to get the two components. These two components, which are mutually perpendicular, are combined to get the resultant.

Let F_1 , F_2 , F_3 and F_4 shown in Fig. 2.12(*a*) be the system of four forces the resultant of which is required.

The procedure to get the resultant is given below:

Step 1*:* Find the components of all the forces in *X* and *Y* directions. Thus, F_{1x} , F_{2x} , F_{3x} , F_{4x} , F_{1y} F_{2y} , F_{3y} , and F_{4y} , are obtained.

Step 2*:* Find the algebraic sum of the component forces in *X* and *Y* directions.

$$
\sum F_x = F_{1x} + F_{2x} + F_{3x} + F_{4x}
$$

$$
\sum F_y = F_{1y} + F_{2y} + F_{3y} + F_{4y}
$$

(**Note:** In the above case F_{2x} , F_{3x} , F_{3y} and F_{4y} have negative values.)

Step 3: Now the system of forces is equal to two mutually perpendicular forces, namely, ΣF_x and ΣF _y as shown in Fig. 2.12(*b*). Since these two forces are perpendicular, the parallelogram of forces becomes a rectangle. Hence the resultant *R* is given by:

$$
R = \sqrt{\left(\Sigma F_x\right)^2 + \left(\Sigma F_y\right)^2} \qquad \qquad \dots (2.7)
$$

and its inclination to *x* axis is given by:

$$
\alpha = \tan^{-1} \left(\frac{\sum F_y}{\sum F_x} \right) \tag{2.8}
$$

Note: $R \cos \alpha = \sum F_x = R_x$...(2.9)

$$
\sim
$$

and
$$
R \sin \alpha = \Sigma F_y = R_y
$$
 ...(2.10)

i.e., ΣF_x and ΣF_y are the *x* and *y* components of the resultant.

The procedure of finding the component of forces and then finding the resultant is illustrated with examples 2.3 to 2.9.

Example 2.3. Determine the resultant of the three forces acting on a hook as shown in Fig. 2.13(a).

Example 2.4. A system of four forces acting on a body is as shown in Fig. 2.14(a). Determine the resultant.

Solution: If θ_1 is the inclination of the 200 N force to *x* axis,

Fig. 2.14
Similarly for the force 120 N,

$$
\tan \theta_2 = \frac{4}{3}, \sin \theta_2 = \frac{4}{5}, \cos \theta_2 = \frac{3}{5}
$$

Now,
 $R_x = \Sigma F_x = 200 \times \frac{2}{\sqrt{5}} - 120 \times \frac{3}{5} - 50 \cos 60^\circ + 100 \sin 40^\circ$
 $= 146.16 \text{ N}.$
 $R_y = \Sigma F_y = 200 \times \frac{1}{\sqrt{5}} + 120 \times \frac{4}{5} - 50 \times \sin 60^\circ - 100 \cos 40^\circ$
 $= 65.54 \text{ N}$
 $R = \sqrt{146.16^2 + 65.54^2}$
 $R = 160.18 \text{ N}.$
 $\alpha = \tan^{-1} \frac{65.54}{146.16} = 24.15^\circ$ as shown in Fig. 2.14(b).

Example 2.5. A system of forces acting on a body resting on an inclined plane is as shown in
2.15. Determine the resultant force if $\theta = 60^\circ$ and if $W = 1000 \text{ N}$; $N = 500 \text{ N}$; $F = 100 \text{ N}$
 $T = 1200 \text{ N}.$

Solution: In this problem, note that selecting X and Y axes parallel to the plane and perpend
to the plane is convenient.

 $R_x = \Sigma F_x = T - F - W \sin \theta$
 $F_y = \sum F_y = N - W \cos 60^\circ = 500 - 1000 \cos 60^\circ = 0$
 $= 1200 - 1000 - 1000 \cos 60^\circ = 0$
 \therefore Resultant is force of 233.97 N directed up the plane.

Example 2.6. Two forces acting on a body are 500 N and 1000 N as shown in Fig. 2.

Determine the third force F such that the resultant of all the three forces is 1000 N directed a
to x axis.

Solution: Let the third force F make an angle θ with x axis.

Example 2.5. A system of forces acting on a body resting on an inclined plane is as shown in Fig. 2.15. Determine the resultant force if θ *= 60° and if W = 1000 N; N = 500 N; F = 100 N; and T = 1200 N.*

Solution: In this problem, note that selecting *X* and *Y* axes parallel to the plane and perpendicular to the plane is convenient.

$$
R_x = \Sigma F_x = T - F - W \sin \theta
$$

= 1200 - 100 - 1000 sin 60° = 233.97 N

$$
R_y = \Sigma F_y = N - W \cos 60° = 500 - 1000 \cos 60° = 0.
$$

∴ **Resultant is force of 233.97 N directed up the plane.**

Example 2.6. Two forces acting on a body are 500 N and 1000 N as shown in Fig. 2.16(a). Determine the third force F such that the resultant of all the three forces is 1000 N directed at 45° to x axis.

Solution: Let the third force *F* make an angle θ with *x* axis.

Example 2.7. Three forces acting at a point are shown in Fig. 2.17. The direction of the 300 N forces may vary, but the angle between them is always 40°. Determine the value of θ *for which the resultant of the three forces is directed parallel to b-b.*

Solution: Let the *x* and *y* axes be as shown in Fig. 2.17. If the resultant is directed along the *x* axis, its component in *y* direction is zero.

i.e., $0 = \Sigma F_y = 300 \sin \theta + 300 \sin (40 + \theta) - 500 \sin 30^\circ$

$$
\therefore \sin \theta + \sin (40 + \theta) = \frac{500 \sin 30^{\circ}}{300}
$$

= 0.8333

$$
\therefore \sin \theta + \sin(40 + \theta) = 0.8333
$$

$$
2 \sin \left(\frac{40 + \theta + \theta}{2} \right) \times \cos \left(\frac{40 + \theta - \theta}{2} \right) = 0.8333
$$

$$
2 \sin (20 + \theta) \times \cos (20) = 0.8333
$$

$$
\theta = 6.35^{\circ}
$$

Fig. 2.17

2.8 EQUILIBRIANT OF A FORCE SYSTEM

We have seen that the resultant of a system of forces is a single force which will have the same effect as the system of forces. According to Newton's second law of motion, the body starts moving with acceleration in the direction of the resultant force. If we apply a force equal and opposite to the resultant, the body should come to the equilibrium state. Such a force is called equilibriant. Thus an equilibriant of a system of forces may be defined as *the force which brings the body to the state of equilibrium and obviously, this forces is equal in magnitude, but opposite in the direction to the resultant.*

2.9 COMPOSITION OF COPLANAR NON-CONCURRENT FORCE SYSTEM

Let F_1 , F_2 and F_3 [Fig. 2.18(*a*)] constitute a system of forces acting on a body. Each force can be replaced by a force of the same magnitude and acting in the same direction at point *O* and a moment about *O*. Thus, the given system in Fig. 2.18(*a*) is equal to the system shown in Fig. 2.18(*b*) where ΣM_O is the algebraic sum of the moments of the given forces about *O*.

At *O*, the concurrent force F_1 , F_2 and F_3 can be combined as usual to get the resultant force *R*. Now the resultant of the given system is equal to force *R* at *O* and a moment ΣM_0 as shown in Fig. 2.18(*c*).

The force *R* and moment ΣM_O shown in Fig. 2.18(*c*) can be replaced by a single force *R* acting at a distance *d* from *O* such that the moment produced by this force *R* is equal to ΣM_0 [Ref. 2.18(*d*)].

Thus, we get a single force *R* acting at a distance *d* from the point *O* which gives the same effect as the constituent forces of the systems. Thus, the resultant of the given forces may be reduced to a single force.

Mathematically,

$$
R = \sqrt{(\Sigma F_x)^2 + (\Sigma F_y)^2}
$$

\n
$$
\tan \alpha = \frac{\Sigma F_y}{\Sigma F_x}
$$
...(2.11)
\n
$$
d = \frac{\Sigma M_o}{R}
$$

where, ΣF_r – algebraic sum of the components of all forces in *x* direction ΣF _{*y*} – algebraic sum of the components of all forces in *y* direction α – inclination of the resultant *R* to *x* direction ΣM_{\odot} – algebraic sum of the moments of all the forces about point *O*

d – is distance of the resultant *R* from the point *O*.

Note: *R* is marked at distance *d* such that it produces the same direction of moment about point *O* as $\sum M_O$.

Sometimes the values of ΣF_x and ΣF_y may come out to be zero, but ΣM_o may exist. This means that the resultant of the system gets reduced to a pure couple.

2.10 *x* **AND** *y* **INTERCEPTS OF RESULTANT**

In some situations we may be interested in finding only the distance of *R* along *x* or *y* axis, that is *x* and *y* intercepts.

Let *d* be the distance of the resultant from *O* and α be its inclination to *x* axis (Fig. 2.19). Then the intercepts are given by:

$$
\tan \alpha = \frac{2M_o}{2F_x}
$$
\nwhere, $\sum F_x$ – algebraic sum of the components of all forces in *x* direction
\n $\sum F_y$ – algebraic sum of the components of all forces in *y* direction
\n $\sum M_O$ – inclination of the resultant *R* to *x* direction
\n $\sum M_O$ – algebraic sum of the moments of all the forces about point *O*
\nNote: *R* is marked at distance *d* such that it produces the same direction of moment about point *O* as $\sum M_O$
\nSometimes the values of $\sum F_x$ and $\sum F_y$ may come out to be zero, but $\sum M_O$ may exist. This means
\nthat the resultant of the system gets reduced to a pure couple.
\n**2.10 x AND y INTERCEPTS OF RESULTANT**
\nIn some situations we may be interested in finding only the distance of *R* along *x* or *y* axis, that is
\n*x* and *y* intercepts.
\nLet *d* be the distance of the resultant from *O* and *α* be its inclination to *x* axis (Fig. 2.19). Then
\nthe intercepts are given by:
\n
$$
x = \frac{d}{\sin \alpha} \qquad y = \frac{d}{\cos \alpha} \qquad ...(2.12)
$$
\nAnother method of finding the intercepts is as follows:
\nLet *R_x* = $\sum F_x$ and *R_y* = $\sum F_y$ be the components of the resultant *R* in *x* and *y* directions. Considering
\nthe moment of *R* about *O* as the sum of moments of its components about *B* (Varignon's theorem)
\nwe get (ref. Fig. 2.20).

Another method of finding the intercepts is as follows:

Let $R_x = \Sigma F_x$ and $R_y = \Sigma F_y$ be the components of the resultant *R* in *x* and *y* directions. Considering the moment of R about O as the sum of moments of its components about B (Varignon's theorem) we get (ref. Fig. 2.20).

Similarly, resolving the resultant into its components at *A*, it can be shown that:

$$
y = \frac{\Sigma M_o}{R_x} = \frac{\Sigma M_o}{\Sigma F_x}
$$
...(2.14)

Example 2.8. A system of loads acting on a beam is shown in Fig. 2.21(a). Determine the resultant of the loads.

Solution: Taking horizontal direction towards left as *x* axis and the vertical downward direction as *y* axis.

> $\Sigma F_x = 20 \cos 60^\circ = 10 \text{ kN}$ $\Sigma F_y = 20 + 30 + 20 \sin 60^\circ = 67.3205$

$$
\tan \alpha = \frac{\Sigma F_y}{\Sigma F_x} = 6.7321
$$

$$
\alpha = 81.55^\circ.
$$

Now taking moment about *A,*

$$
\Sigma M_A = 20 \times 1.5 + 30 \times 3.0 + 20 \sin 60^\circ \times 6
$$

= 223.9231 kN-m

∴ The distance of the resultant from point *O* is given by:

$$
d = \frac{\Sigma M_A}{R} = \frac{223.9231}{68.0592} = 3.290 \text{ m}
$$

$$
x = \frac{d}{\sin \alpha} = \frac{3.290}{\sin 81.55^{\circ}}
$$

x **= 3.326 m.**

The value of *x* intercept may be obtained using Eqn. 2.13 also. Thus,

$$
x = \frac{\Sigma M_0}{R_y} = \frac{\Sigma M_0}{\Sigma F_y} = \frac{223.9231}{67.3205} = 3.326 \text{ m}.
$$

Example 2.9. Find the resultant of the force system shown in Fig. 2.22(a) acting on a lamina of equilateral triangular shape.

Solution: $\Sigma F_x = 80 - 100 \cos 60^\circ - 120 \cos 30^\circ = -73.92 \text{ N}$

$$
R_x = 73.92 \text{ N}
$$

\n
$$
\Sigma F_y = 80 + 120 \sin 30^\circ - 100 \sin 60^\circ
$$

\n
$$
R_y = 53.40 \text{ N}
$$

\n
$$
R = \sqrt{73.92^2 + 53.40^2}
$$

\n
$$
R = 91.19 \text{ N}
$$

\n
$$
\tan \alpha = \frac{\Sigma F_y}{\Sigma F_x} = \frac{53.40}{73.92}
$$

\n
$$
\alpha = 35.84^\circ
$$

Let *x* be the distance from A at which the resultant cuts *AC*. Then taking *A* as moment centre,

53.40 $x = 80 \times 100 \sin 60^{\circ} + 80 \times 50 + 120 \sin 30^{\circ} \times 100$

 $x = 317.008$ mm to the right of A as shown in Fig. 2.22(a).

Example 2.10. Four forces having magnitudes of 200 N, 400 N, 600 N and 800 N respectively, are acting along the four sides (1 m each) of a square ABCD taken in order, as shown in Fig. 2.23. Determine the magnitude and direction of the resultant force.

 $x = \frac{\sum M_A}{\sum A} = \frac{1000}{\sum A}$ *^y* 400 *M* $\frac{\Sigma M_A}{\Sigma F_v} = \frac{1000}{400} = 2.5 \text{ m}$ as shown in the Fig. 2.23(*a*).

Example 2.11. Forces 2, $\sqrt{3}$, $\sqrt{3}$ *and 2 kN respectively act at one of the angular points of a regular hexagon towards five other angular points. Determine the magnitude and direction of the resultant force.*

Solution: Let the system of forces be as shown in Fig. 2.24 shown below:

Let O be the centre of the encircling circle A , B , C , D , E and F . In regular hexagon each side is equal to the radius *AO*. Hence *OAB* is equilateral triangle.

∴
$$
\angle OAB = 60^{\circ}
$$

\nIn $\triangle ABC$, $BA = BC$
\n∴ $\angle CAB = \angle BCA$
\nBut $\angle CAB + \angle BCA = \text{interior angle at } B$
\n $= 180^{\circ} - 120^{\circ} = 60^{\circ}$

Example 2.12. Find the resultant of a set of coplanar forces acting on a lamina as shown in Fig. 2.25(a). Each square has side of 10 mm.

Solution: If θ_1 , θ_2 and θ_3 are the slopes of the forces 2 kN, 5 kN and 1.5 kN forces with respect to *x* axis, then

$$
\tan \theta_1 = \frac{10}{10} = 1 \quad \therefore \quad \theta_1 = 45^{\circ}
$$
\n
$$
\tan \theta_2 = \frac{30}{40} \quad \therefore \quad \theta_2 = 36.87^{\circ}
$$
\n
$$
\tan \theta_3 = \frac{10}{20} \quad \therefore \quad \theta_3 = 26.565
$$
\n
$$
R_x = \Sigma F_x = 2 \cos 45^{\circ} + 5 \cos 36.87^{\circ} - 1.5 \cos 26.565^{\circ} = 4.072 \text{ kN}
$$
\n
$$
R_y = \Sigma F_y = 2 \sin 45^{\circ} - 5 \sin 36.87^{\circ} - 1.5 \sin 26.565^{\circ} = 2.26 \text{ kN}
$$
\n
$$
R = \sqrt{(\Sigma F_x)^2 + (\Sigma F_y)^2} = 4.66 \text{ kN}.
$$
\n
$$
\tan \alpha = \frac{2.26}{4.072}
$$
\n
$$
\therefore \quad \alpha = 28.99^{\circ}.
$$
\nDistance *d* of the resultant from *O* is given by
\n
$$
Rd = \Sigma M_0
$$
\n
$$
4.66d = 2 \times \cos 45^{\circ} \times 30 + 5 \times \sin 36.87^{\circ} \times 50 + 1.5 \times \sin 26.565^{\circ} \times 10 = 199.13
$$
\nNote: To find moment of forces about *O*, 2 kN force is resolved at its intersection with *y* axis and *g* kN and 1.5 kN forces are resolved at their intersection with *x* axis, and then Varignon theorem is used.
\nExample 2.13 Determine the resultant of four parallel forces acting on the axle of a vehicle as shown in Fig. 2.26.
\nSolution: Let *x* and *y* axes be selected as shown in the Figure
\n
$$
R_x = \Sigma F_x = 0
$$
\n
$$
R_y = \Sigma F_y = 60 + 20 + 30 + 40 = 150 \text{ kN}
$$

Distance *d*

 $4.66d = 2 \times \cos 45^{\circ} \times 30 + 5 \times \sin 36.87^{\circ} \times 50 + 1.5 \times \sin 26.565^{\circ} \times 10$ $= 199.13$ *d* **= 42.77 mm** as shown in Fig. 2.25(*a*).

Note: To find moment of forces about *O*, 2 kN force is resolved at it's intersection with *y* axis and 5 kN and 1.5 kN forces are resolved at their intersection with *x* axis, and then Varignon theorem is used.

Example 2.13 Determine the resultant of four parallel forces acting on the axle of a vehicle as shown in Fig. 2.26.

Solution: Let *x* and *y* axes be selected as shown in the Figure

$$
R_x = \Sigma F_x = 0
$$

\n
$$
R_y = \Sigma F_y = 60 + 20 + 30 + 40 = 150 \text{ kN}
$$

$$
R = \sqrt{0^2 + 150^2} = 150 \text{ kN}
$$

Taking clockwise moment as +ve,

$$
\Sigma M_A = 60 \times 0 + 20 \times 1 + 30 \times 3 + 40 \times 4
$$

 = 270 kN-m

∴ Distance of resultant from *A*

$$
x = \frac{270}{150} = 1.8 \text{ m as shown in the figure.}
$$

Example 2.14. Determine the resultant of system of parallel forces acting on a beam as shown in Fig. 2.27.

Fig. 2.27

Solution: Selecting *x* and *y* axes as shown in Fig. 2.27

$$
R_x = \Sigma F_x = 0
$$

\n
$$
R_y = \Sigma F_y = 80 - 30 + 40 - 50 + 60 = 100 \text{ kN}
$$

\n
$$
\therefore R = \sqrt{0^2 + 100^2} = 100 \text{ kN, in } y \text{-direction}
$$

Taking clockwise moment as positive,

$$
\Sigma M_A = 80 \times 0 - 30 \times 2 + 40 \times 4 - 50 \times 8 + 60 \times 10 = 300
$$
 kN-m

∴ The distance of resultant from *A* is,

$$
x = \frac{\Sigma M_A}{R_y} = \frac{300}{100} = 3 \text{ m as shown in Fig. 2.27.}
$$

Example 2.15. The system of forces acting on a bell crank is shown in Fig. 2.28(a). Determine the magnitude, direction and the point of application of the resultant.

Solution:
$$
R_x = \Sigma F_x = 500 \text{ cos } 60^\circ - 700 = -450 \text{ N} = 450 \text{ N (from right to left)}
$$

\n $R_y = \Sigma F_y = -1200 - 1000 - 500 \sin 60^\circ = -2633.01 \text{ N}$
\n $R = \sqrt{450^2 + (2633.01)^2}$
\n $R = 2671.19 \text{ N.}$
\n $\tan \alpha = \frac{2633.01}{450}$
\nLet the point of application of the resultant be at a distance x from the point *O* along the horizontal arm. Then,
\n $x \times 2633.01 = 500 \sin 60^\circ \times 300 + 1000 \times 150$
\n $-1200 \times 150 \cos 60^\circ + 700 \times 300 \sin 60^\circ$
\n $x = 141.195 \text{ mm, as shown in Fig. 2.28(a).$
\nExample 2.16. Various forces to be considered for the stability analysis of a dam are shown in H
\nFig. 2.29. The dam is safe if the resultant force passes through middle third of the base. Verify,
\nwhether the dam is safe.
\nSolution:
\n $R_x = \Sigma F_x = 500 \text{ kN}$
\n $R_y = \Sigma F_y = +1120 - 120 + 420 = 1420 \text{ kN}$
\nLet x be the distance from *O* where the resulting cut downward
\n $= 1420 \text{ kN}$
\n $R_y = 2.29$
\nThen
\n $x = 4.126 \text{ m}$
\n $R_y = 2.50 \text{ kN}$
\n $R_y = 2.670 \text{ kN}$
\n $R_z = \Sigma F_z = 500 \text{ kN}$
\n $R_x = 2.500 \text{ kN}$
\n $R_y = 2.29$
\nThen
\n $x = 4.126 \text{ m}$
\n

Let the point of application of the resultant be at a distance x from the point O along the horizontal arm. Then,

> $x \times 2633.01 = 500 \sin 60^{\circ} \times 300 + 1000 \times 150$ $-1200 \times 150 \cos 60^{\circ} + 700 \times 300 \sin 60^{\circ}$ *x* **= 141.195 mm,** as shown in Fig. 2.28(*a*).

Example 2.16. Various forces to be considered for the stability analysis of a dam are shown in the Fig. 2.29. The dam is safe if the resultant force passes through middle third of the base. Verify whether the dam is safe.

 $Solution:$

$$
R_x = \Sigma F_x = 500 \text{ kN}
$$

\n
$$
R_y = \Sigma F_y = + 1120 - 120 + 420 = 1420 \text{ kN}
$$

\n= 1420 kN (downward)

Let x be the distance from O where the resultant cuts the base.

Then
\n
$$
x \sum F_y = \sum M_O
$$

\n $x \times 1420 = 500 \times 4 + 1120 \times 2 - 120 \times 4 + 420 \times 5$
\n $x = 4.126 \text{ m}$

The resultant passes through the middle third of the base i.e., between ⁷ $\frac{7}{3}$ m, and $2 \times \frac{7}{3}$ **3** $x \frac{7}{2}$ m. **Hence, the dam is safe.**

Solution: The roof is inclined at 45° to horizontal and loads are at 90° to the roof. Hence, the loads are also inclined at 45° to vertical/horizontal.

Now,

$$
R_x = \Sigma F_x = (5 + 10 + 10 + 5 + 5 + 10 + 10 + 5) \cos 45^\circ
$$

= 60 \times $\frac{1}{\sqrt{2}}$ = 42.426 kN

$$
\Sigma F_y = -(5 + 10 + 10 + 5) \sin 45^\circ + (5 + 10 + 10 + 5) \sin 45^\circ
$$

= 0

$$
\therefore R = \Sigma F_x = 42.426 kN
$$

and its direction is horizontal.

 $60 \times -$

Let *R* be at a distance *d* from the ridge *A*.

Then, $Rd = \sum M_A$

$$
\times \frac{1}{\sqrt{2}} \times d = 5 \times \frac{3}{\sqrt{2}} + 10 \times \frac{2}{\sqrt{2}} + 10 \times \frac{1}{\sqrt{2}} + 10 \times \frac{1}{\sqrt{2}} + 10 \times \frac{2}{\sqrt{2}} + 5 \times
$$

d **= 1.5 m**

∴ **Resultant is a horizontal force of magnitude 42.426 kN at 1.5 m below A.**

Example 2.18. Determine the magnitude, direction and line of action of the equilibriant of the given set of coplanar forces acting on a planar structure shown in Fig. 2.31.

Solution: The two 40 kN forces acting on the smooth pulley may be replaced by a pair of 40 kN forces acting at centre of pulley *C* and parallel to the given forces, since the sum of moments of the two given forces about *C* is zero.

Now,
$$
R_x = \Sigma F_x = 20 \cos 45^\circ - 30 \cos 60^\circ - 50 \cos 30^\circ + 40 \cos 20^\circ - 40 \sin 30^\circ
$$

 $\Sigma F_x = -26.5714 \text{ kN}$
= 26.5714 kN (from right to left)

 $5 \times \frac{3}{\sqrt{2}}$ 2

Fig. 2.31

$$
\Sigma F_y = -20 \sin 45^\circ - 20 + 20 - 30 \sin 60^\circ - 50 \sin 30^\circ
$$

\n
$$
-40 \sin 20^\circ - 40 \cos 30^\circ
$$

\n
$$
= -113.4447 \text{ kN} = 113.4447 \text{ kN downward.}
$$

\n
$$
R = \sqrt{(\Sigma F_x)^2 + (\Sigma F_y)^2} = 116.515 \text{ kN}
$$

\n
$$
\tan \alpha = \frac{113.4447}{26.5714}
$$

\n
$$
\alpha = 76.82^\circ \text{ as shown in Fig. 2.31}(b)
$$

\n
$$
\cot AB \text{ at a distance } x \text{ from } A. \text{ Then,}
$$

\n
$$
r \Sigma F_x = \Sigma M
$$

Let the resultant interse

$$
x \Sigma F_y = \Sigma M_A
$$

\n
$$
x \times 113.4447 = 20 \times 4 - 20 \times 4 + (30 \sin 60^\circ) \times 6 + (50 \sin 30^\circ)
$$

\n
$$
\times 2 - (50 \cos 30^\circ) \times 2 + (40 \cos 20^\circ) \times 3
$$

\n
$$
- (40 \sin 30^\circ) \times 3
$$

\n
$$
x = 1.516 \text{ m}
$$

The equilibriant is equal and opposite to the resultant. Hence, it is as shown in Fig. 2.31(*a*) in which **E** = 116.515 kN, α = 76.82° and x = 1.516 m.

2.11 TYPES OF FORCES ON A BODY

Before taking up equilibrium conditions of a body, it is necessary to identify the various forces acting on it. The various forces acting on a body may be grouped into:

(*a*) Applied Forces

(*b*) Non-applied Forces

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(a) Applied Forces

These are the forces applied externally to a body. Each of the forces is having contact with the body. Depending upon type of their contact with the body, the applied forces may be classified as:

- (*i*) Point Force
- (*ii*) Distributed Forces

(*i***)** *Point Force***:** It is the one which has got contact with the body at a point. Practically there is no force which will have contact with the body at a single point. However, when the contact area is small compared to the other dimensions in the problem, for simplicity of calculation the force may be considered as a point load. If a person stands on a ladder, his weight be taken as an applied point load [Fig. 1.2]. Characteristics of such forces are already explained in Art. 1.1. (i) Pout Force: This the neutrich has got contact with the body at a point. Practice (i) Point Force: It is the one which has got contact with the body at a point. However, when the cis small compared to the other dimensi

(*ii***)** *Distributed Forces***:** Distributed forces may act over a line, a surface or a volume. Correspondingly they are known as linear, surface and body forces.

Linear Force: A linear force is one that acts along a line on the body. It is usually represented with abscissa representing the position on the body and ordinate representing the magnitude of the load. An example of this force is a beam subjected to uniformly distributed/uniformly varying or general load. The force *'dF*' at any small length '*ds*' is given by

$$
dF = w \, ds
$$

Surface Force: A force acting on the surface of a body is known as surface force. The hydrostatic pressure acting on a Dam is an example of surface force. The force *dF* acting on any area *dA* is given by

$$
dF = p \, dA
$$

where 'p' is the intensity of force per unit area.

Body Force: A body force is the force exerted from each and every particle of the mass of the body. Example of this type of force is the weight of a block acting on the body under consideration. The body force *dF* of volume *dV* is given by

$$
dF = \gamma \, dV
$$

where γ is the force per unit volume.

(b) Non-applied Forces

There are two types of non-applied forces: (*a*) Self weight and (*b*) Reactions.

Self weight: Everybody subjected to gravitational acceleration and hence has got a selfweight. *W = mg*

where *m* is mass of the body and *g* is gravitational acceleration $(9.81 \text{ m/sec}^2 \text{ near the earth surface})$

Selfweight always acts in vertically downward direction. When analysing equilibrium conditions of a body, selfweight is treated as acting through the centre of gravity of the body. If selfweight is very small, it may be neglected.

Reactions: These are self-adjusting forces developed by the other bodies which come in contact with the body under consideration. According to Newton's third law of motion, the reactions are equal and opposite to the actions. The reactions adjust themselves to bring the body to equilibrium.

If the surface of contact is smooth, the direction of the reaction is normal to the surface of contact. If the surface of contact is not smooth, apart from normal reaction, there will be frictional reaction also. Hence the resultant reaction will not be normal to the surface of contact.

2.12 FREE BODY DIAGRAM

In many problems, it is essential to isolate the body under consideration from the other bodies in contact and draw all the forces acting on the body. For this, first the body is drawn and then applied forces, selfweight and the reactions at the points of contact with other bodies are drawn. Such *a diagram of the body in which the body under consideration is freed from all the contact surfaces and shows all the forces acting on it (including reactions at contact surfaces), is called a* **Free Body Diagram (FBD).** Free Body Diagrams (FBD) are shown for few typical cases in Table 2.2.

Table 2.2 Free Body Diagrams (FBD) for a Few Typical Cases

2.13 EQUILIBRIUM OF BODIES

A body is said to be in equilibrium when it is at rest or has uniform motion. According to Newton's law of motion, it means the resultant of all the forces acting on a body in equilibrium is zero. The resultant of coplanar system of forces acting on a body is zero when—

- (*a*) The algebraic sum of the component of forces along each of the two mutually perpendicular directions is zero (translatory motion is zero).
- (*b*) The algebraic sum of moment of all the forces about any point in the plane is zero (rotational moment is zero).

The above conditions for coplanar concurrent and non-concurrent forces is discussed and illustrated in this article.

2.14 EQUILIBRIUM OF CONCURRENT FORCE SYSTEMS

If the moment equilibrium condition is considered about the concurrent point of application of the forces, the equation results into zero equal to zero. Hence this is not at all a useful equation. Thus only the conditions to ensure translatory motion is zero gives useful equations. In planar problems the resultant R of a system of forces is zero only when the following conditions are satisfied.

$$
\Sigma F_x = 0 \qquad \qquad \dots (2.15a)
$$

$$
\Sigma F_y = 0 \qquad \qquad \dots (2.15b)
$$

It may be observed that only one of the above two conditions is not sufficient. For example, ΣF _x = 0 means that *R* cos α = 0. This will ensure that the resultant *R* cannot exist in any direction except in *y*-direction ($\alpha = 90^{\circ}$). Hence the condition $\Sigma F_y = 0$ also should be satisfied to ensure the resultant *R* does not exist, that is, the equilibrium condition exists. After drawing free body diagrams for each of the body under the action of concurrent force system equations 2.15a and 2.15b may be written and the problems may be solved.

If a body is in equilibrium under the action of only three concurrent forces, Lami's theorem also may be used.

Lami's theorem states : If a body is in equilibrium under the action of three forces, each force is proportional to the sine of the angle between the other two forces.

Thus, for the system of forces shown in Fig. 2.32(*a*).

$$
\frac{F_1}{\sin \alpha} = \frac{F_2}{\sin \beta} = \frac{F_3}{\sin \gamma}
$$
...(2.15c)

Proof: Draw the three forces F_1 , F_2 and F_3 one after the other in direction and magnitude starting from point *a*. Since the body is in equilibrium (resultant is zero), the last point must coincide with *a*. Thus, it results in triangle of forces *abc* as shown in Fig. 2.32(*b*). Now, the external angles at *a*, *b* and *c* are equal to β , γ and α .

Fig. 2.32

Since *ab* is parallel to F_1 ,

bc parallel to F_2 and

ca parallel to F_3 .

In the triangle of forces *abc*

$$
ab = F_1,
$$

$$
bc = F_2 \text{ and } ca = F_3.
$$

Applying sine rule for the triangle *abc,*

$$
\frac{ab}{\sin(180 - \alpha)} = \frac{bc}{\sin(180 - \beta)} = \frac{ca}{\sin(180 - \gamma)}
$$

i.e.,

$$
\frac{F_1}{\sin \alpha} = \frac{F_2}{\sin \beta} = \frac{F_3}{\sin \gamma}
$$

Note: While determining the direction of the reaction on a body note that if the body is in equilibrium under the action of only three coplanar forces, those three forces must be concurrent.

Example 2.19. A sphere of weight 100 N is tied to a smooth wall by a string as shown in Fig. 2.33(a). Find the tension T in the string and reaction R of the wall.

Solution: Free body diagram of the sphere is as shown in Fig. 2.33(*b*). Figure 2.33(*c*) shows all the forces moving away from the centre of the ball. Applying Lami's theorem to the system of forces.

$$
\frac{T}{\sin 90^\circ} = \frac{R}{\sin(180 - 15)} = \frac{100}{\sin(90 + 15)}
$$

$$
T = 103.53 \text{ N.}
$$

$$
R = 26.79 \text{ N.}
$$

The above problem may be solved using equations of equilibrium also.

Taking horizontal direction as *x* axis and vertical direction as *y* axis,

Fig. 2.33

T cos 15[°] - 100 = 0
\nT = 103.53 N.
\n
$$
\Sigma F_x = 0
$$
 gives
\nR - T sin 15[°] = 0
\nR = 26.79 N.

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Notes:

1. The string can have only tension in it (it can pull a body), but cannot have compression in it (cannot push a body).

2. The wall reaction is a push, but cannot be a pull on the body.

3. If the magnitude of reaction comes out to be negative, then assumed direction of reaction is wrong. It is acting exactly in the opposite to the assumed direction. However, the magnitude will be the same. Hence no further analysis is required. This advantage is not there in using Lami's equation. Hence, it is advisable for beginners to use equations of equilibrium, instead of Lami's theorem even if the body is in equilibrium under the action of only three forces.

Example 2.20. Determine the horizontal force P to be applied to a block of weight 1500 N to hold it in position on a smooth inclined plane AB which makes an angle of 30° with the horizontal [Fig. 2.34(a)].

Solution: The body is in equilibrium under the action of applied force *P*, self-weight 1500 *N* and normal reaction *R* from the plane. Since *R*, which is normal to the plane, makes 30° with the vertical (or 60° with the horizontal),

Note: Since the body is in equilibrium under the action of only three forces the above problem can be solved using Lami's theorem as given below:

$$
\frac{R}{\sin 90^\circ} = \frac{P}{\sin(180 - 30)} = \frac{1500}{\sin(90 + 30)}
$$

R = 1732.06 and P = 866.03.

Example 2.21. A roller of weight 10 kN rests on a smooth horizontal floor and is connected to the floor by the bar AC as shown in Fig. 2.35 (a). Determine the force in the bar AC and reaction from floor, if the roller is subjected to a horizontal force of 5 kN and an inclined force of 7 kN as shown in the figure.

Solution: A bar can develop a tensile force or a compressive force. Let the force developed be a compressive force S (push on the cylinder). Free Body Diagram of the cylinder is as shown in Fig. 2.35(*b*).

Since there are more than three forces in the system, Lami's equations cannot be applied. Consider the components in horizontal and vertical directions.

$$
\Sigma H = 0
$$

\n
$$
S \cos 30^{\circ} + 5 - 7 \cos 45^{\circ} = 0
$$

\n
$$
S = \frac{7 \cos 45^{\circ} - 5}{\cos 30^{\circ}} = -0.058 \text{ kN}
$$

Since the value of S is negative the force exerted by the bar is not a push, but it is pull (tensile force in bar) of magnitude 0.058 kN.

> $\Sigma V = 0$ *R* – 10 – 7 sin 45° + *S* sin 30° = 0 $R = 10 + 7 \sin 45^\circ - S \sin 30^\circ$ $= 10 + 7 \sin 45^\circ - (-0.058) \sin 30^\circ$ *R* **= 14.979 kN.**

Example 2.22. A cord ACB 5 m long is attached at points A and B to two vertical walls 3 m apart as shown in Fig. 2.36(a). A pully C of negligible radius carries a suspended load of 200 N and is free to roll without friction along the cord. Determine the position of equilibrium, as defined by the distance X, that the pulley will assume and also the tensile force in the cord.

Solution: The pulley *C* is in equilibrium under the action of tensile forces in *CA* and *CB* and vertical downward load 200 N. The tensile forces in segment *CA* and *CB* are the same since the pulley is frictionless. Now consider the equilibrium of pulley *C* shown in Fig. 2.36(*b*).

$$
\Sigma H = 0
$$

\n
$$
T \cos \theta_1 - T \cos \theta_2 = 0
$$

\n
$$
\theta_1 = \theta_2
$$
, say, equal to θ
\nNow, let *BC* be extended to *D*.
\n
$$
\Delta CFD = \Delta CFA
$$

\n
$$
CD = AC
$$

 $BD = BC + CD = BC + AC =$ length of chord = 5 m *DE* = 3 m ∴ *BE* = 4 m B (a) (b) 3 m D F ^A ^I G E 1 m <u>200 N</u> x - →
- - - - - - - H T ¦T $θ$ \sim θ θ θ T Θ_1 T 200 N C **Fig. 2.36** As ∆*BHI* is similar to ∆*BDE* $HI = \frac{BI}{BE} \times DE =$ 1 4 \times 3 = 0.75 ∴ *AH* = 3 – 0.75 = 2.25 $x = 1.125$ m. Since $AH = 2x$. At *C*, $\Sigma V = 0$ $2 \times T \sin \theta = 200$ $2 \times T \times$ $\frac{4}{5}$ = 200 **Example 2.3.** A roller of reaction of Reference of the wheel to just until the force is able to contain the contained force P applied to the end of a string variable the terms when the force is able to the solution. When

Example 2.23. A roller of radius r = 300 mm and weight 2000 N is to be pulled over a curb of height 150 mm [Fig. 2.37(a)] by a horizontal force P applied to the end of a string wound tightly around the circumference of the roller. Find the magnitude of P required to start the roller move over the curb. What is the least pull P through the centre of the wheel to just turn the roller over the curb? **Solution:** When the roller is about to turn over the curb, the contact with the floor is lost and hence there is no reaction from the floor. The reaction *R* from the curb must pass through the intersection of *P* and the line of action of self weight, since the body is in equilibrium under the action of only three forces (all the three forces must be concurrent).

Referring to Fig. 2.37(*b*),

∴ $T = 125 N$.

$$
\cos \alpha = \frac{OC}{AO} = \frac{(300 - 150)}{300} = \frac{1}{2}
$$
\n
$$
\alpha = 60^{\circ}
$$

Now the reaction from the curb must pass through the centre of the wheel since the other two forces pass through that point. Its inclination to vertical is $\theta = 60^{\circ}$.

If the triangle of forces *ABC* is constructed [Fig. 2.38(*b*)], representing selfweight by *AB*, reaction *R* by *BC* and pull *P* by *AC*, it may be observed that *AC* to be least, it should be perpendicular to *BC*. In other words, *P* makes 90° with the line of action of *R*.

From triangle of forces *ABC*, we get

 $P = AC = AB \sin \theta = 2000 \sin 60^\circ$ *P* **= 1732.05 N.**

2.15 EQUILIBRIUM OF CONNECTED BODIES

When two or more bodies are in contact with one another, the system of forces appears as though it is a non-concurrent forces system. However, when each body is considered separately, in many situations it turns out to be a set of concurrent force system. In such instances, first, the body subjected to only two unknown forces is to be analysed followed by the analysis of other connected body/bodies. Such examples are illustrated below.

Example 2.24. A system of connected flexible cables shown in Fig. 2.39(a) is supporting two vertical forces 200 N and 250 N at points B and D. Determine the forces in various segments of the cable.

Solution: Free body diagrams of points *B* and *D* are shown in Fig. 2.39(*b*). Let the forces in the members be as shown in the figure.

Applying Lami's theorem to the system of forces at point *D*,

situations it turns out to be a set of concurrent force system. In such instances, first, subject to only two unknown forces is to be analysed followed by the analysis of other body/bodies. Such examples are illustrated below. **Example 2.24.** A system of connected flexible cables shown in Fig. 2.39(a) is supporting to forces 200 N and 250 N at points B and D. Determine the forces in various segments of
$$
\frac{1}{2}
$$
 from 250 N. **Figure 2.24.** A system of connected flexible cables shown in Fig. 2.39(a) is supporting to the system of forces at point D, and 120° = $\frac{T_1}{\sin 120^\circ}$ = $\frac{T_2}{\sin 135^\circ}$ = $\frac{250}{\sin 105^\circ}$.\n\nSolution: Free body diagrams of points B and D are shown in Fig. 2.39(b). Let the form of the system of forces at point D, and 135° = $\frac{T_1}{\sin 105^\circ}$ = 24.14 N.\n\nT₂ = 183.01 N.\n\nConsider the system of forces at $EV = 0$ \nT₃ cos 30° = 200 - T₂ cos 60° = 0 \nT₃ = 336.60 N.\n\nT₄ = T₂ sin 105° = T₂ sin 105° = T₃ sin 30° = 0 \nT₄ = 183.01 x sin 60° + 336.60 sin 30° = T₄ = 326.79 N.\n\nExample 2.25. A rope AB, 4.5 m long is connected at two points A and B at the same level A load of 1500 N is suspended from a point C on the rope 1.5 m from A as shown in Fi. What bad connected at a point D on the rope, 1 m from B will be necessary to keep the

Example 2.25. A rope AB, 4.5 m long is connected at two points A and B at the same level 4 m apart. A load of 1500 N is suspended from a point C on the rope 1.5m from A as shown in Fig. 2.40(a). What load connected at a point D on the rope, 1 m from B will be necessary to keep the position CD level ?

48
\n
$$
\frac{1}{4} \sqrt[3]{\frac{1}{2}} \sqrt[3]{\frac{1}{2}} = \frac{1}{2} \sqrt[3]{\frac{1}{2}} = 2.40
$$
\nSolution: Drop perpendiculars *CF* and *AE* = x
\n
$$
AE = x
$$

\nFrom $\triangle ABC$,
\nNow,
\n
$$
AB = 4 \text{ m}
$$

\nand
\n
$$
AC + CD + BD = 4.5 \text{ m}
$$

\ni.e.,
\n
$$
ED = 4.5 - 1.5 = 1.0 = 2.0 \text{ m}
$$

\n
$$
\therefore
$$

\nFrom $\triangle BFD$,
\n
$$
BP = AB - (AE + EF)
$$

\n
$$
= 2.0 \text{ m}
$$

\n
$$
\therefore
$$

\n
$$
BF = AB - (AE + EF)
$$

\n
$$
= 2.5 \text{ m}
$$

\n
$$
x^2 - (2 - x)^2 = 1.25
$$

\ni.e.,
\n
$$
x^2 - 4 + 4x = x^2 = 1.25
$$

\n
$$
x = 1.3125 \text{ m}
$$

\n
$$
\therefore
$$

\n
$$
\alpha = \cos^{-1}(\frac{1.3125}{1.5}) = 28.955^{\circ}
$$

\n
$$
\beta = \cos^{-1}(\frac{2 - 1.3125}{1.5}) = 46.567^{\circ}
$$

\nApplying 1ami's theorem to the system of forces acting at point C [Fig. 2.40(b)], we get
\n

$$
T_2 = 2711.09 \text{ N}
$$

Applying Lami's theorem to the system of forces at B [Fig. 2.40(c)], we get

$$
\frac{T_3}{\sin 90^\circ} = \frac{W}{\sin(180 - 46.567)^\circ} = \frac{T_2}{\sin 136.567^\circ}
$$

$$
T_3 = 3993.28 \text{ N}
$$

and
$$
W = 2863.53 \text{ N}.
$$

Example 2.26. A wire rope is fixed at two points A and D as shown in Fig. 2.41(a). Two weights 20 kN and 30 kN are attached to it at B and C, respectively. The weights rest with portions AB and BC inclined at angles 30° and 50° respectively, to the vertical as shown in figure. Find the tension in the wire in segments AB, BC and CD and also the inclination of the segments CD to vertical. Solution: Applying Lami's theorem for the system of forces at *B* [Fig. 2.41(*b*)], we get

$$
\frac{T_1}{\sin 50^\circ} = \frac{T_2}{\sin 150^\circ} = \frac{20}{\sin 160^\circ}
$$

$$
T_1 = 44.79 \text{ kN.}
$$

$$
T_2 = 29.24 \text{ kN.}
$$

Writing equations of equilibrium for the system of forces at *C* [Fig. 2.41(*c*)], we get

*Example 2.27***.** *A wire is fixed at two points A and D as shown in Fig. 2.42(a). Two weights 20 kN and 25 kN are supported at B and C, respectively. When equilibrium is reached it is found that inclination of AB is 30° and that of CD is 60° to the vertical. Determine the tension in the segments AB, BC and CD of the rope and also the inclination of BC to the vertical.*

 $T_2 \sin \theta = T_1 \sin 30^\circ$...(*i*) $T_2 \cos \theta = T_1 \cos 30^\circ - 20$...(*ii*) $A_{\mu\nu}$ 60° T_3 T_{1} T_3 T_4 T_1 T_1 $_{\rm I_2}$ T_2 T_2 $\overset{\mathsf{T}_2}{\longleftarrow}$ C 20 kN 25 kN 20 kN 25 kN (b) (c) θ 30° $\theta \rightarrow 60^\circ$ (a) \overline{B} T₂ \overline{A} \overline{A} 30° θ **Fig. 2.42** T_2 sin $\theta = T_3$ sin 60°(*iii*) $T_2 \cos \theta = 25 - T_3 \cos 60^\circ$...(*iv*) From (1) and (3) *T*₁ sin 30° = *T*₃ sin 60° ∴ $T_1 = \sqrt{3}T_3$...(*v*) From (2) and (4), $T_1 \cos 30^\circ - 20 = 25 - T_3 \cos 60^\circ$ $\sqrt{3}T_3 \frac{\sqrt{3}}{2}$ 1 $T_3 \frac{V_3}{2} + T_3 \frac{1}{2} = 45$ *T***³ = 22.5 kN.** From (5), $T_1 = 38.97$ kN. From (1), $T_2 \sin \theta = 19.48$ From (2), $T_2 \cos \theta = 13.75$ ∴ tan $\theta = 1.4167$ θ = **54.78.** *T***² = 23.84 kN.** Fig. 2.42

To T_2 is not T_3 in T_4
 T_5
 T_6
 T_7
 T_8
 T_8
 T_9
 T_1 is not T_2
 T_3 is not T_3
 T_4 is not T_5

Fig. 2.42

To T_2 is not T_3
 T_1 is not T_2
 T_3 is not T_3
 T

Solution: Writing equations of equilibrium for the system of forces at joints *B* and *C* [Figs. 2.42(*b*) and (*c*)], we have

Example 2.28. A 600 N cylinder is supported by the frame BCD as shown in Fig. 2.43(a). The frame is hinged at D. Determine the reactions at A, B, C and D.

Solution: Free body diagram of sphere and frame are shown in Figs. 2.43(*b*) and 2.43(*c*), respectively.

Considering the sphere,

\n
$$
\Sigma V = 0, \text{ gives}
$$
\n
$$
R_B = 600 \text{ N.}
$$
\n
$$
\Sigma H = 0, \text{ gives}
$$
\n
$$
R_A = R_C
$$
\n... (i)

As the frame is in equilibrium under the action of three forces only, they must be concurrent forces. In other words, reaction at *D* has line of action alone *OD*. Hence, its inclination to horizontal is given by:

$$
\tan \alpha = \frac{450}{150} = 3
$$
\n
$$
\alpha = 71.5650^{\circ}
$$
\n
$$
\Sigma V = 0, \text{ gives}
$$
\n
$$
R_D \sin \alpha = R_B = 600
$$
\n
$$
\therefore \quad R_D = 632.456 \text{ N.}
$$
\n
$$
\Sigma H = 0, \text{ gives}
$$
\n
$$
R_C = R_D \cos \alpha
$$
\n
$$
\therefore \quad R_C = 200 \text{ N.}
$$
\nFrom (2),
$$
\therefore \quad R_A = 200 \text{ N.}
$$

Example 2.29. Two smooth spheres each of radius 100 mm and weight 100 N, rest in a horizontal channel having vertical walls, the distance between which is 360 mm. Find the reactions at the points of contacts A, B, C and D shown in Fig. 2.44(a).

Solution: Let O_1 and O_2 be the centres of the first and second spheres. Drop perpendicular O_1P to the horizontal line through O_2 . Figures 2.44(*b*) and 2.44(*c*) show free body diagram of the sphere 1 and 2, respectively. Since the surface of contact are smooth, reaction of *B* is in the radial direction, *i.e.*, in the direction O_1O_2 . Let it make angle *a* with the horizontal. Then,

$$
\cos \alpha = \frac{O_2 P}{O_1 O_2} = \frac{360 - O_1 A - O_2 D}{O_1 B + B O_2} = \frac{360 - 100 - 100}{100 + 100} = 0.8
$$

$$
\sin \alpha = 0.6.
$$

Consider sphere No. 1.

$$
\Sigma V = 0, \text{ gives}
$$
\n
$$
R_B \times 0.6 = 100
$$
\n
$$
R_B = 166.67 \text{ N.}
$$
\n
$$
\Sigma H = 0, \text{ gives}
$$
\n
$$
R_A = R_B \times 0.8
$$
\n
$$
\therefore R_A = 133.33 \text{ N.}
$$

Example 2.30. Two cylinders, A of weight 4000 N and B of weight 2000 N rest on smooth inclines as shown in Fig. 2.45(a). They are connected by a bar of negligible weight hinged to each cylinder at its geometric centre by smooth pins. Find the force P to be applied as shown in the figure such that it will hold the system in the given position.

Solution: Figures 2.45(*b*) and 2.45(*c*) show the free body diagram of the two cylinders. Applying Lami's theorem to the system of forces on cylinder *A*, we get

$$
\frac{C}{\sin 60^{\circ}} = \frac{4000}{\sin(60 + 90 - 15)}
$$

$$
C = 4898.98 \text{ N}
$$

Consider cylinder *B.*

Summation of the forces parallel to the inclined plane $(45^{\circ}$ to horizontal) = 0, gives: *P* cos 15° + 2000 cos 45° – *C* cos 60° = 0

∴ $P = \frac{4898.98 \cos 60^\circ - 2000 \cos 45}{150}$

P **= 1071.8 N.**

2.16 EQUILIBRIUM OF NON-CONCURRENT FORCE SYSTEMS

To satisfy the condition that there should not be translatory motion of the body subjected to coplanar forces, in this case also the following two conditions are to be satisfied:

$$
\Sigma F_x = 0; \quad \Sigma F_y = 0
$$

To satisfy the condition that there should not be rotational motion the following condition should be satisfied.

 $\Sigma M_A = 0$, where *A* is any point in the plane.

cos 15

°

 \degree – 2000 cos 45 \degree

Hence, the necessary and sufficient conditions for equilibrium of coplanar force system are:

and
$$
\Sigma F_x = 0
$$

$$
\Sigma F_y = 0
$$

$$
\Sigma M_A = 0
$$
...(2.16)

Referring to Fig. 2.46, *A, B* and *C* are three points in the *x*-*y* plane which are not collinear. Any system of forces acting on a body may be replaced by a single force and a couple. Let *R* be such resultant force.

Fig. 2.46

If the condition $\Sigma M_A = 0$ is satisfied, we can conclude that the resultant *R* should be a force passing through *A*. If $\Sigma M_B = 0$ is satisfied, the line of action of the resultant passes through *B*, *i.e.*, *AB* is the line of action. Now, if $\Sigma M_C = 0$ is also satisfied, the resultant should be zero since the line of action of the resultant (line *AB*) cannot pass through *C* also. Thus if *A, B, C* are not collinear points, the following conditions are necessary and sufficient conditions of equilibrium.

$$
\Sigma M_A = 0
$$

\n
$$
\Sigma M_B = 0
$$

\n
$$
\Sigma M_C = 0
$$
...(2.17)

The equilibrium Equations 2.16 are not independent of equilibrium Equations 2.15. Two of them are common to the two sets. Referring to Fig. 2.46, *R* can be resolved into its components perpendicular to and parallel to *AC*.

Thus, $\Sigma M_C = 0$ is identical to $\Sigma F_x = 0$ and $\Sigma M_B = 0$ is identical to $\Sigma F_y = 0$. Hence any one of the following sets may be used as equations of equilibrium:

> $\overline{1}$ $\overline{1}$

> $\overline{\mathcal{L}}$ ì

> \mathbf{I}

- (1) $\Sigma F_x = 0; \Sigma F_y = 0; \Sigma M_A = 0$
- (2) If line AB is not in y direction,

$$
\Sigma F_y = 0; \Sigma M_A = 0; \Sigma M_B = 0
$$

(3) If line AB is not in x direction

 $\Sigma F_x = 0; \Sigma M_A = 0; \Sigma M_B = 0$

(4) If A, B and C are non-collinear

 $\sum M_A = 0; \sum M_B = 0; \sum M_C = 0$

It can be proved that *if a system is in equilibrium under the action of three forces, those three forces must be concurrent.*

Let P_1 , P_2 and P_3 be the forces acting on a body as shown in Fig. 2.47 and let P_1 and P_2 intersect at *A*. Then applying moment equilibrium condition about *A*, we get

$$
\Sigma M_A = 0 \qquad i.e. \ P_3 d = 0
$$

where *d* is distance of line of action of P_3 from *A*.

Since P_3 is not zero,

 $d = 0$, *i.e.* P_3 also must pass through *A*.

Hence, the proposition is proved.

and

...(2.18)

Three problems are solved in this article. The problem of finding the reactions at supports of beams and trusses also fall under this categories. But these problems are solved in the chapters analysis of beams and trusses.

Example 2.31. The 12 m boom AB weighs 1 kN, the distance of the centre of gravity G being 6 m from A. For the position shown, determine the tension T in the cable and the reaction at B [Ref. Fig. 2.48(a)].

Solution: The free body diagram of the boom is shown in Fig. 2.48(*b*).

$$
\alpha = \tan^{-1} \frac{6.0981}{9.6962}
$$

 $= 32.17^{\circ}$ as shown in Fig. 2.48(*c*).

Example 2.32. A cable car used for carrying materials in a hydroelectric project is at rest on a track formed at an angle of 30° with the vertical. The gross weight of the car and its load is 60 kN and its centroid is at a point 800 mm from the track half way between the axles. The car is held by a cable as shown in Fig. 2.49. The axles of the car are at a distance 1.2 m. Find the tension in the cables and reaction at each of the axles neglecting friction of the track.

Solution: Let *T* be the tension in the cable and the reaction at the pair of wheels be R_1 and R_2 as shown in Fig. 2.49.

Now, Σ of forces parallel to the track = 0, gives

T – 60 sin 60° = 0

$$
T = 51.9615
$$
 kN.

Taking moment equilibrium condition about upper axle point on track, we get $R_1 \times 1200 + T \times 600 - 60 \sin 60^\circ \times 800 - 60 \cos 60^\circ \times 600 = 0$

R1 = **23.6603 kN.**

 \sum of forces normal to the plane = 0, gives

$$
R_1 + R_2 - 60 \cos 60^\circ = 0
$$

 $R_2 = 30 - 23.6603$

*R***² = 6.3397 kN.**

*Example 2.33***.** *A hollow right circular cylinder of radius 800 mm is open at both ends and rests on a smooth horizontal plane as shown in Fig. 2.50(a). Inside the cylinder there are two spheres having weights 1 kN and 3 kN and radii 400 mm and 600 mm, respectively. The lower sphere also rests on the horizontal plane. Neglecting friction find the minimum weight W of the cylinder for which it will not tip over.* Example 2.32. Acadie car used for carrying materials in a hydrodectric project is at a point of the verical. The gross weight of the car and its load be its centroid is at a point 800 mm from the track old p wy between th

Solution: Join the centres of spheres, O_1 and O_2 and drop O_1D perpendicular to horizontal through O_2 .

Now,
\n
$$
O_1O_2 = 400 + 600 = 1000 \text{ mm}
$$
\n
$$
O_2 = 1600 - 400 - 600 = 600 \text{ mm}
$$

If α is the inclination of O_2O_1 to horizontal,

$$
\cos \alpha = \frac{600}{1000} = 0.6
$$

$$
\sin \alpha = 0.8
$$

Fig. 2.50

Free body diagrams of cylinder and spheres are shown in Fig. 2.50(*b*) and (*c*). Considering the equilibrium of the spheres,

$$
\Sigma M_{O_2} = 0, \text{ gives}
$$
\n
$$
R_1 \times O_1O_2 \sin \alpha - 1 \times O_2D = 0
$$
\n
$$
R_1 \times 1000 \times 0.8 = 1 \times 600
$$
\n
$$
R_1 = 0.75 \text{ kN}
$$
\n
$$
\Sigma H = 0, \text{ gives}
$$
\n
$$
R_2 = R_1 = 0.75 \text{ kN}
$$
\n
$$
\Sigma V = 0, \text{ gives}
$$
\n
$$
R_3 = 1 + 3 = 4 \text{ kN}
$$

Now consider the equilibrium of cylinder. When it is about to tip over *A*, there is no reaction from ground at *B*. The reaction will be only at *A*. $\sum M_A = 0$ gives,

 $R_1 h_1 - R_2 h_2 - W \times 800 = 0$ $R_1(h_1 - h_2) - W \times 800 = 0$ Since $R_1 = R_2$ $0.75 \times O_1D = W \times 800$ $0.75 \times 1000 \times 0.8 = W \times 800$ ∴ $W = 0.75$ kN.

IMPORTANT FORMULAE

- 1. Newton's third law, $F = R$.
	- **2.** Newton's gravitational law, *F = G* m_1m *d* $rac{1}{a^2}$.
- **3.** Weight of a body, $W = mg$.
- **4.** Moment of a force about a point $= F \times d$.

5. Varignon's theorem

$$
Ra = P_1d_1 + P_2d_2 + P_3d_3 + \dots
$$

6. Magnitude of a couple

 $M = Pd$.

- **7.** *P* at any point *A* may be resolved into force *P* at *B* plus a moment $P \times d$, where '*d*' is the perpendicular distance of *B* from the line of action of *P* through *A*. 7. P at any proint A may be resulted into force P at B plus a moment P x d, where

reprendicular riskines of B from the line of action of P through A.

8. $R_x = R \cos \alpha = \Sigma F_x$
 $R_y = R \sin \alpha = \Sigma F_y$
 $\tan \alpha = \frac{\Sigma F_x}{F_x}$
 $\tan \alpha = \frac{\Sigma F_x}{$
	- **8.** $R_x = R \cos \alpha = \sum F_x$ $R_y = R \sin \alpha = \sum F_y$

$$
\tan \alpha = \frac{\Sigma F_y}{\Sigma F_x}
$$

and $R = \sqrt{(\Sigma F_x)^2 + (\Sigma F_y)^2}$.

9. *x* and *y* intercepts of the resultant are given by

$$
x = \frac{d}{\sin \alpha} = \frac{\sum M_o}{\sum F_y}
$$

and
$$
y = \frac{d}{\cos \alpha} = \frac{\sum M_o}{\sum F_x}
$$
.

10. Lami's theorem:

$$
\frac{F_1}{\sin \alpha} = \frac{F_2}{\sin \beta} = \frac{F_3}{\sin \gamma} .
$$

THEORY QUESTIONS

- **1.** State and explain Newton's laws of mechanics.
- **2.** Explain the following laws of mechanics:
	- (*i*) Law of transmissibility of forces
	- (*ii*) Parallelogram law of forces.
- **3.** Write short notes on
	- (*i*) Principles of physical independence
	- (*ii*) Principles of superposition.
- **4.** Explain with examples the following system of forces
	- (*i*) Coplanar concurrent forces
	- (*ii*) Coplanar non-concurrent forces
	- (*iii*) Non-coplanar non-concurrent forces.
- **5.** State and explain Varignon's theorem of moments.
- **6.** Define a couple-moment and list its characteristics.

PROBLEMS FOR EXERCISE

- **1.** A body is subjected to the three forces as shown in Fig. 2.51. If possible determine the direction of the force *F* so that the resultant is in *x*-direction, when
	- (*a*) $F = 5000 \text{ N}$

(*b*) $F = 3000 \text{ N}$

[**Ans.** (*a*) 36.87°; (*b*) Not possible]

2. The frictionless pulley *A* shown in Fig. 2.52 is supported by two bars *AB* and *AC* which are hinged at *B* and *C* to a vertical wall. The flexible cable *DG* hinged at *D*, goes over the pulley and supports a load of 20 kN at *G*. The angles between the various members are shown in the figure. Determine the forces in the bars *AB* and *AC*. Neglect the size of the pulley.

Fig. 2.52

[*Hint*: Force in $AD = 20$ kN, tensile] [Ans. $F_{AB} = 0$; $F_{AC} = 34.641$ kN]

3. A cord supported at *A* and *B* carries a load of 10 kN at *D* and a load of *W* at *C* a shown in Fig. 2.53. Find the value of *W* so that *CD* remains horizontal.[**Ans.** *W* = 30 kN]

4. Three bars, hinged at *A* and *D* and pinned at *B* and *C* as shown in Fig. 2.54 form a four-linked mechanism. Determine the value of *P* that will prevent movement of bars.

 $[Ans. $P = 3047.2 \text{ N}]$$

5. Two identical rollers, each of weights 100 N are supported by an inclined plane and a vertical wall as shown in Fig. 2.55. Assuming smooth surfaces, find the reactions induced at the points of supports *A, B, C* and *D.*

7. Two smooth spheres each of weight *W* and radius '*r*' are in equilibrium in a horizontal channel of width '*b*' $(b < 4r)$ and vertical sides as shown in Fig. 2.57. Find the three reactions from the sides of the channel which are all smooth. Also find the force exerted by each spheres on the other.

[Ans. $R = 10$ kN, downward at $x = 4$ m from A]

9. Determine the resultant of the parallel coplanar force system shown in Fig. 2.59.

[Ans. $R = 800$ N (towards left); $d = 627.50$ mm]

10. An equilateral triangular plate of sides 200 mm is acted upon by four forces as shown in Fig. 2.60. Determine the magnitude and direction of the resultant of this system of forces and its position. **[Ans.** $R = 57.3541$ kN; $\alpha = 6.70^{\circ}$; $d = 11.51$ mm]

11. Determine the magnitude, direction and the line of application of the equilibriant of the set of forces shown in Fig. 2.61, which will keep the plane body ABCDEFGH in equilibrium.

[Ans.
$$
E = 23.6517
$$
 kN; $\alpha = 24.37^{\circ}$; x from $A = 1.041$ m]

12. Determine the resultant of the four forces acting on a body as shown in Fig. 2.62.

[Ans. $R = 200$; $\alpha = 60^{\circ}$; at a distance $y = 8.7679$ m below *O*]

13. A bracket is subjected to the system of forces and couples as shown in Fig. 2.63. Find the resultant of the system and the point of intersection of its line of action with (*a*) line *AB*, (*b*) line *BC* and (*c*) line *CD*.

[Ans. $R = 485.4$ N; $\alpha = 34.50^{\circ}$; $Y_{BA} = 112.5$ mm; $X_{BC} = 163.6$ mm; $Y_{CD} = 93.75$ mm]

Fig. 2.63

14. Determine the resultant of the three forces acting on the dam section shown in Fig. 2.64 and locate its intersection with the base *AB*. For a safe design this intersection should occur within the middle third. Is it a safe design?

[**Ans***.* Resultant intersects *AB* 3.333 m from *A.* It is a safe design]

15. A 1000 N cylinder is supported by a horizontal rod *AB* and a smooth uniform rod *CD* which weighs 500 N as shown in Fig. 2.65. Assuming the pins at *A*, *B*, *C* and *D*, to be frictionless and weight of *AB* negligible, find the reactions at *C* and *D*.

[Ans. $V_C = 937.1$ N (upward); $H_C = 577.3$ N (towards left) $R_D = 562.9$ N (upward)]

16. The frame shown in Fig. 2.66 is supported by a hinge at *E* and by a roller at *D*. Determine the horizontal and vertical components of the reactions at hinge *C* as it acts upon member *BD.* [Ans. $H_C = 140 \text{ N}$ (towards right); $V_C = 35 \text{ N}$ (upward)]

Fig. 2.66

3 Trusses

A truss is a structure made up of slender members pin-connected at ends and is capable of taking loads at joints. They are used as roof trusses to support sloping roofs and as bridge trusses to support deck. In many machines steel trusses are used. Transmission towers are also the examples of trusses. In the case of wooden trusses, the ends are connected by making suitable joints or by nailing and bolting whereas in steel trusses ends are connected by bolting or welding. The trusses are also known as '*pinjointed frames*'.

A truss in which all the members lie in a single plane is called as a *plane truss.* In such trusses loads act in the plane of the truss only. Roof trusses and bridge trusses can be considered as plane trusses. If all the members of a truss do not lie in a single plane, then it is called a space truss. Tripod and transmission towers are the examples of space trusses. In this chapter, the analysis of only plane trusses is considered.

3.1 PERFECT, DEFICIENT AND REDUNDANT TRUSSES

A pinjointed truss which has got just sufficient number of members to resist the loads without undergoing appreciable deformation in shape is called a *perfect truss*. Triangular truss is the simplest perfect truss and it has three joints and three members (Fig. 3.1). Perfect trusses with four and five joints are shown in Figs. 3.2 and 3.3 respectively.

It may be observed that to increase one joint in a perfect truss, two more members are required. Hence the following expression may be written down as the relationship between number of joints *j*, and the number of members *m*, in a perfect truss.

$$
m = 2j - 3 \tag{3.1}
$$

However, the above equation gives only a necessary, but not a sufficient condition of a perfect truss. For example, the two trusses shown in Fig. 3.4(*a*) and (*b*) have the same number of members and joints. The truss shown in Fig. 3.4(*a*) is perfect whereas the one shown in Fig. 3.4(*b*) is not capable of retaining its shape if loaded at the joint marked 6. Therefore, the only necessary and sufficient condition of a perfect truss is that it should retain its shape when load is applied at any joint in any direction. **EPUA:**

A truss is a structure made up of shender members pin-summected at ends and is capable

heads at ionis. They are used as roof trasses to support doping roofs and as bright under

the case of wooden trasses, the e

A truss is said to be *deficient* if the number of members in it are less than that required for a perfect truss. Such trusses cannot retain their shape when loaded. A deficient truss is shown in Fig. 3.5.

A truss is said to be *redundant* if the number of members in it are more than that required in a perfect truss. Such trusses cannot be analysed by making use of the equations of equilibrium alone. Thus, a redundant truss is statically indeterminate. Each extra member adds one degree of indeterminancy. For the analysis of such members the consistency of deformations is to be considered. The truss shown in the Fig. 3.6 is a typical redundant truss. In this truss one diagonal member in each panel is extra. Hence it is a two-degree redundant truss.

In this chapter, only the analysis of perfect frames is considered.

3.2 ASSUMPTIONS

In the theory that is going to be developed in this chapter, the following assumptions are made:

- (1) The ends of the members are pin-connected (hinged);
- (2) The loads act only at the joints;
- (3) Self-weights of the members are negligible;
- (4) Cross-section of the members is uniform.

If at all the cross-section varies, the centre of gravity of the section is assumed to be located along the same longitudinal line.

In reality the members are connected by bolting, riveting or by welding. No special care is taken to ensure perfect pin-connections. However, experiments have shown that assuming pin-connected ends is quite satisfactory since the members used are slender.

Fig. 3.6

Fig. 3.5

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In most of the frames the loads act at the joints. Even if a load is not acting at a joint, it can be replaced by its reaction at the joint and a local bending effect on the member. The frame may be analysed for the joint loads and the local bending effect on the member superposed in the design of that member.

In most of the trusses, the self-weight is really small compared to the loads they carry. Hence self-weight of the members may be neglected.

It is the duty of construction engineer to see that the centroid of all cross-sections lie along a single axis so that the member is held in equilibrium by the two forces acting at its ends.

Because of the assumption of pin-connected ends, it is more appropriate to call the theory that is going to be developed in this chapter as *analysis of pin-connected plane trusses.* Analysis of rigid frames is not covered in this book.

3.3 NATURE OF FORCES IN MEMBERS

The members of a truss are subjected to either tensile or compressive forces. A typical truss *ABCDE* loaded at joint *E* is shown in Fig. 3.7(*a*). The member *BC* is subjected to compressive force *C* as shown in Fig. 3.7(*b*). Effect of this force on the joint *B* (or *C*) is equal and opposite to the force *C* as shown in Fig. 3.7(*b*).

The member *AE* is subjected to tensile force *T*. Its effect on the joints *A* and *E* are as shown in Fig. 3.7(*b*). In the analysis of frame we mark the forces on the joints, instead of the forces in the members as shown in Fig. 3.7(*c*). It may be noted that compressive force in a member is represented in a figure by two arrows going away from each other and a tensile force by two arrows coming towards each other. This is quite logical considering the fact that the markings on the members represent the internal reactive forces developed which are opposite in direction to the applied forces.

3.4 METHODS OF ANALYSIS

The following three methods are available for the analysis of pin-connected frames:

- (1) Method of joints
- (2) Method of section
- (3) Graphical method.

The first two are analytical methods and they are dealt in this chapter.

3.5 METHOD OF JOINTS

At each joint the forces in the members meeting and the loads acting, if any, constitute a system of concurrent forces. Hence, two independent equations of equilibrium can be formed at each joint. First, a joint is selected where there are only two unknown forces. Many times such a joint can be identified only after finding the reaction at the support by considering the equilibrium of the entire frame. Then making use of the two equations of equilibrium at that joint the two unknown forces are found. Then, the next joint is selected for analysis where there are now only two unknown forces. Thus, the analysis proceeds from joints to joint to find the forces in all the members. **EXECUTE ASSEMATE ASSEM AND THE SIGN CONSUMITION**
 C (1) Method of junits, and a windlether the analysis of pin-connected frames:

(1) Method of section

(2) Method of section

(3) Graphical methods and they are dealt i

It may be noted that if there are *j* number of joints, 2*j* number of the equations can be formed. There will be three reactions in a general determinate truss. The force in each member is unknown. Hence, if there are *m* number of members, the total number of unknowns will be *m*+3. A problem can be analysed if there are as many equations as there are unknowns. Hence, a frame analysis problem is determinate if:

$$
2j = m + 3
$$
...(3.2)

This equation is the same as Equation 3.1 which was derived on the consideration of a perfect frame. Hence, a perfect frame is determinate. If $m > 2j - 3$, then the number of unknowns is more than the number of equations. Hence, a redundant frame is indeterminate. If $m < 2j - 3$, then the number of equations is more than the number of unknowns. Since a set of solutions can satisfy such equations, it shows instability of the structure. Hence, a deficient frame is not stable.

EXAMPLES

The method of joints is illustrated with the examples 3.1 to 3.7.

Example 3.1. Find the forces in all the members of the truss shown in Fig. 3.8(a). Tabulate the results.

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Solution: Step 1: Determine the inclinations of all inclined members. In this case,

$$
\tan \theta = \frac{3}{3} = 1
$$
\n
$$
\therefore \qquad \theta = 45^{\circ}
$$

Step 2: Look for a joint at which there are only two unknowns. If such a joint is not available, determine the reactions at the supports, and then at the supports these unknowns may reduce to only two. Step 2: Look for a joint at which there only two unknowns. If such a joint is not
determine the reactions at the supports, and then at the supports these unknowns range real
two.
Wow at joints C, there are only two unknow

Now at joints *C*, there are only two unknowns, *i.e.*, forces in members *CB* and *CD*, say F_{CB} and F_{CD} .

Note: Usually in cantilever type frames, we find such joints without the need to find reactions.

Step 3: Now there are two equations of equilibrium for the forces meeting at the joint and two unknown forces. Hence, the unknown forces can be determined.

At joint *C* [Ref. Fig. 3.8(*b*)] $\Sigma V = 0$ condition shows that the force F_{CB} should act away from the joint *C* so that its vertical component balances the vertical downward load at *C*.

$$
\overline{a}
$$

 $F_{CR} = 40\sqrt{2}$ kN

Now $\Sigma H = 0$ indicates that F_{CD} should act towards C.

$$
F_{CD} - F_{CB} \cos 45^\circ = 0
$$

 F_{CR} sin 45° = 40

$$
F_{CD} = F_{CB} \cos 45^\circ = 40\sqrt{2} \times \frac{1}{\sqrt{2}} = 40 \text{ kN}
$$

Fig. 3.8(b)

Note: If the assumed direction of unknown force is opposite, the value will be negative. Then reverse the direction and proceed.

Step 4: On the diagram of the truss, mark arrows on the members near the joint analysed to indicate the forces on the joint. At the other end, mark the arrows in the reverse direction.

In the present case, near the joint *C*, the arrows are marked on the members *CB* and *CD* to indicate forces F_{CB} and F_{CD} directions as found in the analysis of joint *C*. Then reversed directions are marked in the members *CB* and *CD* near joints *B* and *D,* respectively. $F_{\rm DR}$

Step 5: Look for the next joint where there are only two unknown forces and analyse that joint.

In this case, there are only two unknown forces at the joint *D* as shown in Fig. 3.8(*c*).

$$
\Sigma V = 0
$$

\n
$$
F_{DB} = 40 \text{ kN}
$$

\n
$$
\Sigma H = 0
$$

\n
$$
F_{DE} = 40 \text{ kN}
$$

40 kN

 F_{DE}

D

F_{cD}

45°

 F_{CB}

C

Fig. 3.8(d)

Step 6: Repeat steps 4 and 5 till forces in all the members are found.

The directions of these forces are marked on the diagram. Now the analysis is complete since the forces in all the members are determined.

Step 7: Determine the nature of forces in each member and tabulate the results. Note that if the arrow marks on a member are towards each other, then the member is in tension and if the arrow marks are away from each other, the member is in compression [Ref. Fig. 3.8(*e*)]. In this case,

Example 3.2. Determine the forces in all the members of the truss shown in Fig. 3.9(a) and indicate the magnitude and nature of forces on the diagram of the truss. All inclined members are at 60° to horizontal and length of each member is 2 m.

Solution: Now, we cannot find a joint with only two unknown forces without finding reactions. Consider the equilibrium of the entire frame.

$$
\sum M_A = 0, \text{ gives}
$$

R_D × 4 – 40 × 1 – 60 × 2 – 50 × 3 = 0
\n∴ R_D = 77.5 kN
\n
$$
\Sigma H = 0
$$
, gives
\n∴ R_A = 70.5 kN
\n∴ R_A = 72.5 kN
\nJoint A:
\n $R_A = 72.5$ kN
\nJoint A:
\n $\Sigma V = 0$, gives
\n F_{AB} sin 60° = R_A = 72.5
\n $F_{AB} = 33.7158$ kN (Comp.)
\n $\Sigma H = 0$, gives
\n $F_{AB} = 33.7158$ kN (Comp.)
\n $F_{AE} = 83.7158$ kN (Comp.)
\n $F_{AE} = 41.8579$ kN (Tension)
\nJoint D:
\n $\Sigma V = 0$, gives
\n F_{DC} sin 60° = R_D = 77.5
\n∴
\n $F_{DC} = 89.4893$ kN (Comp.)
\n $F_{DE} = 44.7446$ kN (Tension)
\nJoint B:
\n $\Sigma V = 0$, gives
\n $F_{DE} = 44.7446$ kN (Tension)
\nJoint B:
\n $\Sigma V = 0$, gives
\n $F_{BE} = \frac{72.5 - 40}{\sin 60°} = 37.5278$ (Tension)
\n $F_{BE} = \frac{72.5 - 40}{\sin 60°} = 37.5278$ (Tension)
\n $F_{BC} = (83.7158 + 37.5274) × 0.5$
\n $F_{BC} = (60.6218$ kN (Comp.)
\n $F_{BC} = 60.6218$ kN (Comp.)

Now the forces in all the members are known. If joint *E* is analysed it will give the check for the analysis. The results are shown on the diagram of the truss in Fig. 3.9(*f*).

Example 3.3. Determine the forces in all the members of the truss shown in Fig. 3.10 (a).

Example 3.4. Determine the forces in the members of truss shown in Fig. 3.11(a).

Example 3.5. Analyse the truss shown in Fig. 3.12(a). Solution: All inclined members have the same inclination to horizontal. Now, length of an inclined member is

$$
= \sqrt{3^2 + 4^2} = 5 \text{ m}
$$

$$
\sin \theta = \frac{4}{5}
$$

$$
\cos \theta = \frac{3}{5}
$$

As soon as a joint is analysed the forces on the joint are marked on members [Fig. 3.12(*b*)]

At this stage as no other joint is having only two unknowns, no further progress is possible. Let us find the reactions at the supports considering the whole structure. Let the reaction be as shown in Fig. 3.12(*b*).

$$
\Sigma M_A = 0, \text{ gives}
$$
\n
$$
R_C \times 8 - 20 \times 6 = 0
$$
\n
$$
R_C = 15 \text{ kN}
$$
\n
$$
\Sigma V = 0, \text{ gives}
$$
\n
$$
V_A = 20 \text{ kN}
$$
\n
$$
\Sigma H = 0, \text{ gives}
$$
\n
$$
H_A = R_C = 15 \text{ kN}
$$
\n
$$
\Sigma V = 0, \text{ gives}
$$
\n
$$
F_{AB} - V_A = 0
$$
\n
$$
F_{AB} = 20 \text{ kN (Comp.)}
$$
\n
$$
\Sigma H = 0, \text{ gives}
$$
\n
$$
F_{AF} - H_A = 0
$$
\n
$$
F_{AF} = 15 \text{ kN (Comp.)}
$$

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Note: When three members are meeting at an unloaded joint and out of them two are collinear, then the force in third member will be zero. Such situations are illustrated in Fig. 3.12(*h*) and (*i*).

Example 3.6. Find the forces in all the members of the truss shown in Fig. 3.13(a).

Joint-by-joint analysis is carried out as given below and the joint forces are marked in Fig. 3.13(*b*). Then nature of the force in the members is determined.

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Joint *F*: $\Sigma V = 0$, gives $F_{FG} - 12 = 0$

 $F_{FG} = 12$ kN (Tension)

Example 3.7. Analyse the truss shown in Fig. 3.14(a). All the members are of 3 m length.

Solution: Since all members are 3 m long, all triangles are equilateral and hence all inclined members are at 60° to horizontal. Joint-by-joint analysis is carried out and the forces are represented in Fig. 3.14(*b*). Then nature of the force is determined.

Solution: Since all members are 3 m long, all triangles are equilateral and hence all inclined are at 60° to horizontal. Joint-by-joint analysis is carried out and the forces are repre-
Fig. 3.14(b). Then nature of the force is determined.
Joint G:
$$
\sum V = 0
$$
, gives
 $F_{GF} = 23.0940$ kN (Tension)
 $\sum H = 0$, gives
 $F_{GF} = 23.0940$ kN (Tension)
 $F_{GF} = F_{GF} \cos 60^\circ = 0$
 $F_{FG} = 11.5470$ kN (Comp.)
Joint F: $\sum V = 0$, gives
 $F_{FG} = 13.0940$ kN (Comp.)
 $F_{FG} = 13.0940$ kN (Comp.)
 $\sum H = 0$, gives
 $F_{FD} = 13.0940$ kN (Comp.)
 $F_{FD} = 13.0940$ kN (Comp.)
 $\sum H = 0$, gives
 $F_{FD} = 13.0940$ kN (Comp.)
Now, without finding reaction we cancel. Hence, consider equilibrium of the er-
down that is 30 kN
out $\sum_{i=1}^{40} \sum_{i=1}^{40} \sum_{i=1}^{40}$

Now, without finding reaction we cannot proceed. Hence, consider equilibrium of the entire truss.

 $\sum M_A = 0$ $R_E \times 6 + 10 \times 3 \times \sin 60^\circ - 40 \times 1.5 - 30 \times 4.5 - 20 \times 9 = 0$ ∴ $R_E = 58.1699$ kN

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 $\Sigma V = 0$, gives $V_A = 40 + 30 + 20 - R_E = 31.8301$ kN $\Sigma H = 0$, gives $H_A = 10 \text{ kN}$ Joint *A*: $\Sigma V = 0$, gives F_{AB} sin 60° – 31.8301 = 0 *FAB* **= 36.7543 kN (Comp.)** $\Sigma H = 0$, gives $F_{AC} - F_{AB} \cos 60^{\circ} + 10 = 0$ $F_{AC} = 8.3771 \text{ kN (Tension)}$ Joint *B*: $\Sigma V = 0$, gives F_{BC} sin 60° + F_{AB} sin 60° – 40 = 0 $F_{BC} = 9.4337 \text{ kN (Comp.)}$ $\sum H = 0$, gives $F_{BD} + F_{BC} \cos 60^\circ - F_{BA} \cos 60^\circ = 0$ *FBD* **= 13.6603 kN (Comp.)** Joint *C*: $\Sigma V = 0$, gives F_{CD} sin 60° – F_{BC} sin 60° = 0 $F_{CD} = F_{BC} = 9.4337 \text{ kN (Tension)}$ $\Sigma H = 0$, gives $F_{CE} + F_{AC} - F_{CD} \cos 60^{\circ} - F_{BC} \cos 60^{\circ} = 0$ $F_{CE} = 2 \times 9.4337 \times \frac{1}{2} - 8.3771 = 1.0566$ kN (Comp.) Joint *D*: $\Sigma V = 0$, gives F_{DE} sin 60° – F_{CD} sin 60° – 30 = 0 *FDE* **= 44.0747 kN (Comp.)** I oint A:
 $H_A = 10 \text{ kN}$
 F_{AB} sin 60° - 31.8301 = 0
 $F_{AB} = 36.7543 \text{ kN (Comp.)}$
 $E_{AC} = F_{AB} \cos 60° + 10 = 0$
 $F_{AC} = 8.3771 \text{ kN (Tension)}$

Joint B:
 F_{BC} in 60° - F_{BC} is gives.
 F_{BC} in 60° - F_{BC} is gives.
 $F_{BC} = 9$

3.6 METHOD OF SECTION

In the method of section, after determining the reactions, a section line is drawn passing through not more than three members in which forces are not known such that the frame is cut into two separate parts. Each part should be in equilibrium under the action of loads, reactions and forces in the members that are cut by the section line. Equilibrium of any one of these two parts is considered and the unknown forces in the members cut by the section line are determined. The system of forces acting on either part of truss constitutes a non-concurrent force system. Since there are only three independent equation of equilibrium, there should be only three unknown forces. Hence, in this method it is an essential condition that the section line should pass through not more than three members in which forces are not known and it should separate the frame into two parts.

Thus, the method of section is the application of nonconcurrent force system analysis whereas the method of joints, described in previous article was the application of analysis of concurrent force system.

Under the following two conditions the method of section is preferred over the method of joints:

(1) In a large truss in which forces in only few members are required;

(2) In the situation where the method of joints fails to start/proceed with analysis.

The method of section is illustrated with the examples 3.8 to 3.11. Examples 3.8 and 3.9 are the cases in which method of section is advantageous since forces in only few members are required. Examples 3.10 and 3.11 are the cases in which method of joints fails to start/proceed to get the solution. In practice the frames may be analysed partly by method of section and partly by method of joints as illustrated in example 3.11.

Example 3.8. Determine the forces in the members FH, HG and GI in the truss shown in Fig. 3.15(a). Each load is 10 kN and all triangles are equilateral with sides 4 m.

Solution: Due to symmetry,

$$
R_A = R_0 = \frac{1}{2} \times 10 \times 7 = 35
$$
 kN

Take section (*A*)–(*A*), which cuts the members *FH, GH* and *GI* and separates the truss into two parts. Consider the equilibrium of left hand side part as shown in Fig. 3.15(*b*) (Prefer part in which number of forces are less).

Σ*MG* = 0, gives *FFH* × 4 sin 60° – 35 × 12 + 10 × 10 + 10 × 6 + 10 × 2 = 0 *FFH* **= 69.2820 kN (Comp.)** ∑*V* = 0, gives *FGH* sin 60° + 10 + 10 + 10 – 35 = 0 *FGH* **= 5.7735 kN (Comp.)** ∑ *H* = 0, gives Civil[datas.blogspo](http://Civildatas.blogspot.in)t.in

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$$
F_{GI} - F_{FH} - F_{GH} \cos 60^\circ = 0
$$

$$
F_{GI} = 69.2820 + 5.7735 \cos 60^\circ
$$

= 72.1688 kN (Tension)

*Example 3.9.**Find the magnitude and nature of the forces in the members U3U4, L3L4 and U4L3 of the loaded truss shown in Fig. 3.16(a).*

Solution: To determine reactions, consider equilibrium equations.

Fig. 3.16(a)

Now, $\sum M_{LO} = 0$, gives

 $R_2 \times 36 - 200 \times 6 - 200 \times 12 - 150 \times 18 - 100 \times 24 - 100 \times 30 = 0$

$$
R_2 = 325 \text{ kN}
$$

\n
$$
\Sigma V = 0, \text{ gives}
$$

$$
R_1 = 200 + 200 + 150 + 100 + 100 - 325 = 425
$$
 kN

Take the section (I) – (I) and consider the right hand side part.

Fig. 3.16(b)

Now,
\n
$$
U_3 U_4 = \sqrt{1^2 + 6^2} = 6.0828
$$
\n
$$
\sin \theta_1 = \frac{1}{6.0828} = 0.1644
$$
\n
$$
\cos \theta_1 = \frac{6}{6.0828} = 0.9864
$$

$$
L_3U_4 = \sqrt{6^2 + 8^2} = 10
$$

\n $\sin \theta_2 = 0.6 \cos \theta_2 = 0.8$
\n $\sum M_{U_1} = 0$, gives
\n $F_{L_2L_4} \times 8 - 325 \times 12 + 100 \times 6 = 0$
\n $F_{L_3L_4} = 412.5 \text{ kN (Tension)}$
\n $\sum M_{L_3} = 0$, gives
\n $F_{U_4U_3} \times \cos \theta_1 \times 9 + 100 \times 6 + 100 \times 12 - 325 \times 18 = 0$
\n $F_{U_4U_3} = 456.2072 \text{ kN (Comp.)}$
\n $\sum H = 0$
\n $F_{U_4U_3} \sin \theta_2 - F_{U_4U_3} \cos \theta_1 + F_{L_4L_3} = 0$
\n $F_{U_4U_3} = \frac{456.2072 \times 0.9864 - 412.5}{0.6}$
\n $= 62.5 \text{ kN (Tension)}$
\nExample 3.10. Find the forces in the members (1), (2) and (3) of French truss shown in Fig.
\n $\frac{20 \text{ kN}}{6}$
\n $\$

= 62.5 kN (Tension)

Example 3.10. Find the forces in the members (1), (2) and (3) of French truss shown in Fig. 3.17(a).

*Solution***:** Due to symmetry

$$
R_A = R_B = \frac{1}{2} \times 20 \times 7 = 70
$$
 kN

Now

$$
AC = 4 \times a = \frac{9}{\cos 30^{\circ}}
$$

$$
a = 2.5981 \text{ m.}
$$

Take Section (*A*)–(*A*) and consider the equilibrium of left hand side part of the French Truss shown in Fig. 3.17(*b*).

Drop perpendicular *CE* on *AB.*

Now *CE* = 9 tan 30° and *DE* = 3 m [∴] tan ^θ ⁼ 9 30 3 ³ ¹ 3 ³ tan [×] ° = = ∴ θ **= 60°** ∑ *M ^A* = 0, gives *F*2 sin 60° × 6 – 20 × 2.5981 cos 30° – 20 × 2 × 2.5981 cos 30° – 20 × 3 × 2.5981 cos 30° = 0 *^F*² = 20 × 2.5981 () ¹²³ 6 + + since sin 60° = cos 30° *F***² = 51.9615 kN (Tension)** ∑*V* = 0, gives *F*1 sin 30° – 70 + 20 + 20 + 20 – 51.9615 sin 60° = 0 *F1* **= 110 kN (Comp.)** ∑ *H* = 0, gives *F*3 + *F*2 cos 60° – *F*1 cos 30° = 0 *F***³ = 69.2820 (Tension)** Civil[datas.blogspo](http://Civildatas.blogspot.in)t.in

Note: In this problem, the method of joints cannot give complete solution.

Example 3.11. Find the forces in all the members of the truss shown in Fig. 3.18(a).

*Solution***:** Due to symmetry, the reactions are equal.

$$
R_A = R_E = \frac{1}{2} \times \text{Total load}
$$

$$
= \frac{1}{2} (15 + 30 + 30 + 30 + 15) = 60 \text{ kN}
$$

Drop perpendicular *CH* on *AF.*

In $\triangle ACH$, ∠*ACH* = 45°

∴ *FC* is inclined at 30° to vertical *i.e.*, 60° to horizontal and *CH* = 5 m \bigcirc

It is not possible to find a joint where there are only two unknowns. Hence, consider section $(1)–(1)$.

For left hand side part of the frame:

$$
= \frac{1}{2} (15 + 30 + 30 + 30 + 15) = 60 \text{ kN}
$$

\nDrop perpendicular *CH* on *AF*.
\nIn *AACH*, *AACH* = 45°
\n \therefore *FC* is inclined at 30° to vertical *i.e.*, 60° to horizontal and *CH* = 5 m
\nIt is not possible to find a joint where there are only two unknowns. Hence, consider section
\n(1)-(1).
\nFor left hand side part of the frame:
\n
$$
\sum M_C = 0
$$
, gives
\n $F_{AE} \times 5 - 60 \times 5 + 15 \times 5 + 30 \times 2.5 = 0$
\n \therefore
\nAssuming the directions for F_{FC} and F_{RC} as shown in Fig. 3.18(b)
\n
$$
\sum V = 0
$$
, gives
\n $F_{FC} \sin 60^\circ - F_{BC} \sin 45^\circ + 60 = 15 - 30 = 0$...(i)
\n $F_{FC} \cos 60^\circ - F_{BC} \cos 45^\circ + F_{AE} = 0$
\n $F_{FC} \times \frac{1}{2} - F_{BC} \frac{1}{\sqrt{2}} + 30 = 0$
\n \therefore
\n $F_{FC} = F_{BC} \sqrt{2} - 60$...(ii)
\n $F_{FC} = F_{BC} \sqrt{2} - 60$...(iii)
\n $F_{FC} = F_{BC} \sqrt{2} - 60$...(iv)
\n $F_{BC} = F_{AC} \frac{30 \text{ kN}}{45^\circ}$
\n $F_{FC} = F_{BC} \sqrt{2} - 60$...(iv)
\n $F_{BC} = 30 \text{ kN}$
\nFig. 3.18(b)
\nSubstituting this value of F_{FC} in (1), we get
\n $F_{BC}(\sqrt{2} \sin 60^\circ - \frac{1}{\sqrt{2}}) = 60 \sin 60^\circ - 15$
\n $F_{BC}(\sqrt{2} \sin 60^\circ - \frac{1}{\sqrt{2}}) = 60 \sin 60^\circ - 1$

Substituting this value of F_{FC} in (1), we get

$$
(F_{BC} \sqrt{2} - 60) \sin 60^\circ - F_{BC} \frac{1}{\sqrt{2}} + 15 = 0
$$

$$
F_{BC} \left(\sqrt{2} \sin 60^\circ - \frac{1}{\sqrt{2}} \right) = 60 \sin 60^\circ - 15
$$

$$
F_{BC} = 71.40 \text{ kN (Comp.)}
$$

$$
T \text{RUSSES} \hspace{1.5cm} \textbf{87}
$$

∴ $F_{FC} = 71.4042 \times \sqrt{2} - 60$

= 40.98 kN (Tension)

Assumed directions of F_{BC} and F_{FC} are correct.

Therefore, F_{BC} is in compression and F_{FC} is in tension.

Now we can proceed with method of joints to find the forces in other members. Since it is a symmetric truss, analysis of half the truss is sufficient. Other values may be written down by making use of symmetry.

Joint *B*:

 Σ forces normal to $AC = 0$, gives F_{BF} – 30 cos 45° = 0 $F_{BF} = 21.21 \text{ kN (Comp.)}$ \sum forces parallel to *AC* = 0, gives $F_{AB} - F_{BC} - 30 \sin 45^\circ = 0$ $F_{AB} = 71.40 + 21.21$ **= 92.61 kN (Comp)** Joint *A*: $\Sigma V = 0$, gives F_{AF} sin 30° – F_{AB} sin 45° – 15 + 60 = 0

FAF = **40.98 kN (Tension)**

The results are tabulated below:

+ means tension and – means compression

IMPORTANT FORMULA

1. In a perfect truss

 $m = 2j - 3$

where $m =$ number of members

 j = number of joints.

THEORY QUESTIONS

- **1.** Bring out the differences among perfect, deficient and redundant trusses.
- **2.** State the assumptions made in the analysis of pin jointed trusses.
- **3.** How method of joint differs from the method of section in the analysis of pin jointed trusses?

PROBLEMS FOR EXERCISE

1 to 17: Determine the forces in all the members of the trusses shown in Fig. 3.19 to 3.35. Indicate the nature of forces using the convention tension as +ve and compression as –ve. **1.**

[Ans. F_{AB} = +67.5 kN; F_{BC} = +15 kN; F_{CD} = -25 kN; F_{DE} = -30 kN; F_{EF} = –105 kN; F_{AE} = +62.5 kN; F_{BE} = –62.5 kN; F_{BD} = +25 kN]

2.

3.

Fig. 3.20 (Prob. 2)

 $[Ans. F_{AB} = +82.0738$ kN; $F_{BC} = +73.866$ kN; $F_{CD} = 49.2443$ kN; $F_{DE} = -45$ kN; F_{EF} = -45 kN; F_{FG} = -67.5 kN; F_{BG} = -10.0 kN; F_{FC} = +24.622 kN; F_{CE} = 0; F_{BF} = 10 kN]

4.

are zero members]

[Ans. $F_{AC} = F_{CE} = F_{EG} = +193.1852 \text{ kN}; F_{BD} = F_{DE} = F_{FG} = -193.1852 \text{ kN}; \text{ all others}$

E \overline{C} \overline{A} 200 kN B 4 m 4 m 2 m 2 m D **Fig. 3.22 (Prob. 4)**

[Ans. F_{EC} = +447.2136 kN; F_{CA} = +400 kN; F_{AB} = -447.2136 kN; F_{BD} = -400 kN; $F_{CD} = 0$; $F_{CB} = -200$ kN] F

- **5.** [Ans. $F_{DB} = F_{BA} = +5.7735 \text{ kN}; F_{BC} = F_{DE} = -5.7738 \text{ kN};$ F_{AC} = -2.8868 kN; F_{CE} = -14.4338 kN; F_{DC} = +17.3205 kN; F_{DF} = +20.0 kN]
- **6.** [Ans. $F_{AB} = -30$ kN; $F_{AC} = -160$ kN; $F_{BC} = +50$ kN; F_{BD} = -200 kN; F_{CD} = -50 kN; F_{CE} = -120 kN; F_{DF} = -266.67 kN; F_{DE} = +83.33 kN]

7. [Ans. $F_{AB} = -200 \text{ kN}; F_{AD} = -100 \text{ kN}; F_{BC} = F_{CE} = 0;$ $F_{BD} = 100\sqrt{2}$ kN; $F_{BF} = -100\sqrt{2}$ kN; $F_{DE} = -100$ kN; F_{DG} = 0; F_{EF} = +100 kN; F_{EH} = -100 $\sqrt{2}$ kN; F_{EG} = +100 $\sqrt{2}$ kN; F_{GH} = +100 kN]

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- **11.** [Ans. $F_{AB} = -15$ kN; $F_{AC} = +12\sqrt{2}$ kN; $F_{BD} = -27.5\sqrt{2}$. kN; $F_{BC} = -12.5\sqrt{2}$ kN; $F_{CE} = 0$; $F_{CD} = +25$ kN; $F_{ED} = -27.5\sqrt{2}$ kN]
- **12.** [Ans. $F_{AB} = -17.32 \text{ kN}; F_{AC} = +5 \text{ kN}; F_{BC} = -20 \text{ kN};$ $F_{BD} = -17.32 \text{ kN}; F_{CD} = +20 \text{ kN}; F_{CE} = -15 \text{ kN};$ $F_{DE} = -30$ kN]

13. [Ans. $F_{AB} = 60 \text{ kN}; F_{AC} = +51.96 \text{ kN}; F_{BC} = -20 \text{ kN}; F_{BD} = -40 \text{ kN}; F_{CD} = +40 \text{ kN}; \text{ symmetry}$]

14. [Ans. $F_{AC} = -4.5\sqrt{13}$ kN; $F_{AB} = +13.5$ kN; $F_{BC} = +6$ kN; $F_{BD} = +13.5$ kN; F_{CD} $= -0.5\sqrt{13}$ kN; $F_{CE} = -4\sqrt{10}$ kN; $F_{DE} = 8$ kN]

Fig. 3.32 (Prob. 14)

15. [Ans. $F_{AB} = +10\sqrt{13}$ kN; $F_{AC} = -20$ kN; $F_{CB} = -48.75$ kN; $F_{CE} = -20$ kN; $F_{CD} = -7.5$ kN; F_{BE} = +6.25 $\sqrt{13}$ kN; F_{DE} = 18.75 kN; F_{DF} = -3.75 $\sqrt{13}$ kN; F_{FE} = -7.5 kN]

Fig. 3.33 (Prob. 15)

16. [Ans. $F_{AB} = 16.91 \text{ kN}; F_{AF} = +31.55 \text{ kN}; F_{BF} = +23.91 \text{ kN}; F_{BD} = -23.91 \text{ kN}; F_{BC} = +40 \text{ kN};$ F_{CD} = -40 kN; F_{DE} = -63.1 kN; F_{DF} = +23.91 kN; F_{EF} = +31.55 kN]

17. [Ans. F_{AC} = –67.48 kN; F_{AB} = +53.99 kN; F_{BC} = +10 kN; F_{CD} = –8.33 kN; F_{CE} = –59.15 kN; F_{EF} = -24.5 kN; F_{ED} = +52.81 kN; F_{FD} = +47.21 kN; F_{FG} = -34.64 kN; F_{DG} = +47.32 kN]

18. Find the force in the member *FG* of the triangular Howe truss shown in Fig. 3.36.

(*Hint:* Take section (1)–(1) and find force in *FD*. Then analyse joint *F*) [**Ans.** *+*28 kN]

19. Determine the forces in the members *AB, AC, DF* and *CE* of the scissors truss shown in Fig. 3.37.

(*Hint:* Find reaction *RA* and analyse joint *A*. Take section (1)–(1) and find force in *DF* and *CE*) [Ans. $F_{AB} = -6.25$ W; $F_{AC} = 4.51$ W; $F_{DF} = -3.75$ W; $F_{CE} = +2.75$ W] **20.** Find the force in member *KL* of the French truss shown in Fig. 3.38.

(*Hint:* Take section (*A*)–(*A*) and find F_{LE} and F_{DE} . From joint *D* find F_{DL} . Then analyse joint *L* to get F_{KL} [Ans. +41.96 kN]

Distributed Forces, Centre of Gravity and Moment of Inertia

The term 'distributed force' has been explained in Art 2.11 and this applied force has been classified into linear surface and the body force. The number of such forces acting on a body is infinite. However, these forces can be replaced by their resultant which acts through a point, known as the centre of gravity of the body. In this chapter the method of finding areas of given figures and volumes is explained. Then the terms centroid and centre of gravity and second moment of area (moment of inertia of areas) are explained and method of finding them is illustrated with examples. Theorem of Pappus-Guldinus is introduced which is very useful for finding surface areas and volumes of solids then the method of finding centre of gravity and moment of inertia (mass moment of inertia) of solids is illustrated. **Distributed Forces, Central

Cravity and Moment of Income

The term 'distributed force' has been explained in Art 2.11 and this applied force has been

thow there a strike and the body force. The numeric q start because**

4.1 DETERMINATION OF AREAS AND VOLUMES

In the school education methods of finding areas and volumes of simple cases are taught by many methods. Here we will see the general approach which is common to all cases *i.e.* by the method of integration. In this method the expression for an elemental area will be written then suitable integrations are carried out so as to take care of entire surface/volume. This method is illustrated with standard cases below, first for finding the areas and latter for finding the volumes:

A: Area of Standard Figures

(*i*) Area of a rectangle

Let the size of rectangle be $b \times d$ as shown in Fig. 4.1. *dA* is an elemental area of side $dx \times dy$.

Area of rectangle,
$$
A = \oint dA = \int_{-b/2}^{b/2} \int_{-d/2}^{d/2} dx dy
$$

= $[x]_{-b/2}^{b/2} [y]_{-d/2}^{d/2}$
= bd.

If we take element as shown in Fig. 4.2,

Fig. 4.1 $d/2$ d/2 $-h/2$ $-h/2$ O y dx
⊠⊿dy x

$$
^{94}
$$

$$
A = \int_{-d/2}^{d/2} dA = \int_{-d/2}^{d/2} b \cdot dy
$$

= $b [y]_{-d/2}^{d/2}$
= bd

(*ii*) Area of a triangle of base width '*b*' height '*h*': Referring to Fig. 4.3, let the element be selected as shown by hatched lines

Then
$$
dA = b'dy = b\frac{y}{h} dy
$$

$$
A = \int_0^h dA = \int_0^h b\frac{y}{h} dy
$$

$$
= \frac{b}{h} \left[\frac{y^2}{2} \right]_0^h = \frac{bh}{2}
$$

dy y d/2 d/2 b **Fig. 4.2** y h b dy b **Fig. 4.3**

(*iii*) Area of a circle

Consider the elemental area $dA = rd\theta dr$ as shown in Fig. 4.4. Now,

$$
= b \text{ [y]}^{2/2} = bd
$$
\n(ii) Area of a triangle of base width 'b' height 'h': Referring to
\nFig. 4.3, let the element be selected as shown by hatched
\nlines
\nThen $dA = b'dy = b\frac{y}{h} dy$
\n
$$
A = \int_{0}^{h} dA = \int_{0}^{h} b \frac{y}{h} dy
$$
\n
$$
= \frac{h}{h} \left[\frac{y^{2}}{2} \right]_{0}^{h} = \frac{bh}{2}
$$
\n(iii) Area of a circle
\nConsider the elemental area $dA = r d\theta dr$ as shown in Fig. 4.4. Now,
\n
$$
dA = r d\theta dr
$$
\n
$$
r \text{ varies from 0 to } R \text{ and } \theta \text{ varies from 0 to } 2\pi
$$
\n
$$
= \int_{0}^{2\pi} \left[\frac{r^{2}}{2} \right]_{0}^{R} d\theta
$$
\n
$$
= \int_{0}^{2\pi} \left[\frac{r^{2}}{2} \right]_{0}^{R} d\theta
$$
\n
$$
= \frac{R^{2}}{2} \cdot 2\pi = \pi R^{2}
$$
\nIn the above derivation, if we take variation of θ from 0 to π , we get the area of $\frac{\pi R^{2}}{4}$.
\n
$$
\frac{\pi R^{2}}{2}
$$
 and if the limit is from 0 to $\pi/2$ the area of quarter of a circle is obtained as $\frac{\pi R^{2}}{4}$.

2

O dr $rd\theta$ $r_{\rm d0}$ θ _R \rightarrow \rightarrow x y

In the above derivation, if we take variation of θ from 0 to π , we get the area of semicircle as πR^2 2 R^2 and if the limit is from 0 to $\pi/2$ the area of quarter of a circle is obtained as πR^2 4 $\frac{R^2}{\cdots}$.

(*iv*) Area of a sector of a circle

Area of a sector of a circle with included angle 2α shown in Fig. 4.5 is to be determined. The elemental area is as shown in the figure

 $dA = rd\theta \cdot dr$

θ varies from –α to α and *r* varies from 0 to *R*

$$
dA = r dθ \cdot dr
$$
\n0 varies from –α to α and r varies from 0 to R
\n∴ $A = \oint dA = \int_{-a0}^{aR} r d\theta dr$
\n
$$
= \int_{-a}^{a} \left[\frac{r^2}{2} \right]_0^a d\theta = \int_{-a}^{a} \frac{R^2}{a^2} d\theta
$$

\n
$$
= \left[\frac{R^2}{2} \theta \right]_{-a}^{a} = \frac{R^2}{2} (2\alpha) = R^2 \alpha.
$$
\n(v) Area of a parabolic spandrel
\nTwo types of parabolic curves are possible
\n(a) $y = kx^2$
\n(b) $y^2 = kx$
\nCase a: This curve is shown in Fig. 4.6.
\nThe area of the element
\n $dA = y dx$
\n
$$
= k \left[\frac{x^3}{3} \right]_0^a = \frac{k a^3}{3}
$$

\n
$$
= k \left[\frac{x^3}{3} \right]_0^a = \frac{k a^3}{3}
$$

\nWe know, when $x = a$, $y = h$
\ni.e., $h = ka^2$ or $k = \frac{h}{a^2}$
\n∴ $A = \frac{ka^3}{3} = \frac{h}{a^2} \frac{a^3}{3} = \frac{1}{3} ha = \frac{1}{3} \text{rd the area of rectangle of size } a \times h$
\nCase b: In this case $y^2 = kx$
\nReferring to Fig. 4.7
\n $dA = y dx = \sqrt{kx} dx$
\n $A = \int_0^a y dx = \int_0^a kx dx$

(*v*) Area of a parabolic spandrel Two types of parabolic curves are possible (*a*) $y = kx^2$

$$
(b) \ \ y^2 = kx
$$

Case a: This curve is shown in Fig. 4.6. The area of the element

$$
dA = y dx
$$

= $kx^2 dx$

$$
\therefore A = \int_0^a dA = \int_0^a kx^2 dx
$$

= $k \left[\frac{x^3}{3} \right]_0^a = \frac{ka^3}{3}$

Fig. 4.5 δ^{α} đθ $\overline{\theta}$ dr \Box \Box rd θ \rightarrow R

y

We know, when $x = a$, $y = h$

i.e.,
$$
h = ka^2 \text{ or } k = \frac{h}{a^2}
$$

∴ $A =$ $\frac{3}{2}$ $\frac{1}{2}$ $\frac{3}{2}$ $\frac{a^3}{2}$ = $\frac{1}{3}$ ha = $\frac{1}{3}$ rd $3 \t a^2 \t 3 \t 3 \t 3$ $\frac{ka^3}{3} = \frac{h}{a^2} \frac{a^3}{3} = \frac{1}{3} ha = \frac{1}{3}$ rd the area of rectangle of size $a \times h$ *Case b:* In this case $y^2 = kx$

Referring to Fig. 4.7

$$
dA = y \, dx = \sqrt{kx} \, dx
$$

$$
A = \int_0^a y \, dx = \int_0^a \sqrt{kx} \, dx
$$

a

x=a y = kx² x dx **Fig. 4.6**

h

x

y

DISTRIBUTED FORCES, CENTRE OF GRAVITY AND MOMENT OF INERTIA **97**

 $=$ \sqrt{k} $x^{3/2}$ 0 2 3 *a* $\left[x^{3/2} \frac{2}{3}\right]_0^a = \sqrt{k} \frac{2}{3} a^{3/2}$ We know that, when $x = a$, $y = h$ \therefore $h^2 = ka$ or $k = \frac{h^2}{h}$ *a* Hence $A = \frac{h}{\sqrt{2}} \cdot \frac{2}{2}$ 3 *h a* $\frac{2}{9} \cdot a^{3/2}$ *i.e.*, $A = \frac{2}{3}$ *ha* = $\frac{2}{3}$ rd the area of rectangle of size *a* × *h*. (*vi*) Surface area of a cone Consider the cone shown in Fig. 4.8. Now, $y = \frac{x}{7}$ *h R* .. $h^2 = ka$ or $k = \frac{h^2}{a}$

Hence $A = \frac{h}{\sqrt{a}}, \frac{2}{3}$ or a^{12}

Hence $A = \frac{2}{\sqrt{a}}, \frac{2}{3}$ or a^{12}

Le., $A = \frac{2}{3}$ hat $a = \frac{2}{3}$ rule area of rectungle of size $a \times h$,

(v) Surface area of a cone

Consider th

Surface area of the element,

$$
dA = 2\pi y \, dl = 2\pi \frac{x}{h} R \, dl
$$

$$
= 2\pi \frac{x}{h} R \, \frac{dx}{\sin \alpha}
$$

$$
A = \frac{2\pi R}{h \sin \alpha} \left[\frac{x^2}{2} \right]_0^h
$$

$$
= \frac{\pi R h}{\sin \alpha} = \pi R l
$$

y $v^2 = kx$ h x dx x=a **Fig. 4.7**

(*vii*) Surface area of a sphere

Consider the sphere of radius *R* shown in Fig. 4.9. The element considered is the parallel circle at distance *y* from the diametral axis of sphere.

$$
dS = 2\pi x \ R d\theta
$$

= 2\pi R cos \theta R d\theta, since $x = R \cos \theta$

$$
S = 2\pi R^2 \int_{-\pi/2}^{\pi/2} \cos \theta \ d\theta
$$

= 2\pi R² [sin \theta]^{\pi/2}_{\pi/2}
= 4\pi R²

B: Volume of Standard Solids

(i) Volume of a parallelopiped

\nLet the size of the parallelopiped be a × b × c. The volume of the element is

\n
$$
dV = dV \cdot dV = dx \, dy \, dz
$$
\n
$$
V = \int_{0}^{a} \int_{0}^{b} dx \, dy \, dz
$$
\n
$$
= [x]_{0}^{a} \int_{0}^{b} \int_{0}^{b} [z]_{0}^{c} = abc
$$
\n(ii) Volume of a cone

\nReferring to Fig. 4.8

\n
$$
dV = \pi y^{2} \cdot dx = \pi \frac{x^{2}}{h^{2}} R^{2} dx, \text{ since } y = \frac{x}{h} R^{2}
$$
\n
$$
V = \frac{\pi}{h^{2}} R^{2} \int_{0}^{h} x^{2} dx = \frac{\pi}{h^{2}} R^{2} \left[\frac{x^{3}}{3} \right]_{0}^{h}
$$
\n
$$
= \frac{\pi}{h^{2}} R^{2} \frac{h^{3}}{3} = \frac{\pi R^{2} h^{3}}{3}
$$
\n(iii) Volume of a sphere

\nReferring to Fig. 4

\n
$$
dV = \pi (R^{2} - y^{2}) dy
$$
\nBut

\n
$$
x^{2} + y^{2} = R^{2}
$$
\ni.e.,

\n
$$
dV = \pi (R^{2} - y^{2}) dy
$$
\n
$$
V = \int_{R}^{\infty} \pi (R^{2} - y^{2}) dy
$$
\n
$$
V = \pi \left[R^{2} \cdot R - \frac{R^{3}}{3} - \left\{ -R^{3} - \frac{(-R^{3})}{3} \right\} \right]
$$
\n
$$
= \pi R \left[1 - \frac{1}{3} + 1 - \frac{1}{3} \right] = \frac{4}{3} \pi R^{3}
$$
\nThe surface areas and volumes of solids of revolutions like cone, spheres may be easily found using theorems of Pappus and Guditmus. This will be taken up latter in this chapter, since it needs the term centroid of generating lines.

The surface areas and volumes of solids of revolutions like cone, spheres may be easily found using theorems of Pappus and Guldinus. This will be taken up latter in this chapter, since it needs the term centroid of generating lines.

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4.2 CENTRE OF GRAVITY AND CENTROIDS

Consider the suspended body shown in Fig. $4.10(a)$. The self weight of various parts of this body are acting vertically downward. The only upward force is the force *T* in the string. To satisfy the equilibrium condition the resultant weight of the body. *W* must act along the line of string 1–1. Now, if the position is changed and the body is suspended again (Fig. 4.10(*b*)), it will reach equilibrium condition in a particular position. Let the line of action of the resultant weight be 2–2 intersecting 1–1 at *G*. It is obvious that if the body is suspended in any other position, the line of action of resultant weight *W* passes through *G*. This point is called the centre of gravity of the body. Thus *centre of gravity can be defined as the point through which the resultant of force of gravity of the body acts*.

The above method of locating centre of gravity is the practical method. If one desires to locating centre of gravity of a body analytically, it is to be noted that the resultant of weight of various portions of the body is to be determined. For this Varignon's theorem, which states the moment of resultant force is equal to the sum of moments of component forces, can be used.

Referring to Fig. 4.11, let W_i be the weight of an element in the given body. *W* be the total weight of the body. Let the coordinates of the element be x_i , y_i , z_i and that of centroid *G* be x_c , y_c , z_c . Since *W* is the resultant of W_i forces,

$$
W = W_1 + W_2 + W_3 + \cdots
$$

\n
$$
= \Sigma W_i
$$

\nand
\n
$$
Wx_c = W_1x_1 + W_2x_2 + W_3x_3 + \cdots
$$

\n
$$
\therefore \qquad Wx_c = \Sigma W_i x_i = \oint x dw
$$

\nSimilarly,
\n
$$
W_y_c = \Sigma W_i y_i = \oint y dw
$$

\nand
\n
$$
W_z = \Sigma W_i z_i = \oint z dw
$$

Fig. 4.11

xi xc

z

Wi ^G ^W

 \bigcirc y_i y_c \longrightarrow x

 $z_i^{\prime} \mid z_{\rm c}^{\prime}$

...(4.1)

If M is the mass of the body and m_i that of the element, then

$$
M = \frac{W}{g} \text{ and } m_i = \frac{W_i}{g}, \text{ hence we get}
$$

\n
$$
Mx_c = \sum m_i x_i = \oint x_i dm
$$

\n
$$
My_c = \sum m_i y_i = \oint y_i dm
$$

\nand
\n
$$
Mz_c = \sum m_i z_i = \oint z_i dm
$$

\n...(4.2)

If the body is made up of uniform material of unit weight χ then we know $W_i = V_i \chi$, where V represents volume, then equation 4.1 reduces to

$$
Vx_c = \Sigma V_i x_i = \oint x dV
$$

\n
$$
Vy_c = \Sigma V_i y_i = \oint y dV
$$

\n
$$
Vz_c = \Sigma V_i z_i = \oint z dV
$$
...(4.3)

If the body is a flat plate of uniform thickness, in *x*-*y* plane, $W_i = \gamma A_i t$ (Ref. Fig. 4.12). Hence equation 4.1 reduces to

If the body is a wire of uniform cross-section in plane *x*, *y* (Ref. Fig. 4.13) the equation 4.1 reduces to

$$
Lx_c = \sum L_i x_i = \oint x dL
$$

\n
$$
Ly_c = \sum L_i y_i = \oint y dL
$$
...(4.5)

The term centre of gravity is used only when the gravitational forces (weights) are considered. This term is applicable to solids. Equations 4.2 in which only masses are used the point obtained is termed as *centre of mass*. The central points obtained for volumes, surfaces and line segments (obtained by eqn. 4.3, 4.4 and 4.5) are termed as *centroids.*

4.3 CENTROID OF A LINE

Centroid of a line can be determined using equation 4.5. Method of finding the centroid of a line for some standard cases is illustrated below:

(*i*) *Centroid of a straight line*

Selecting the *x*-coordinate along the line (Fig. 4.14)

$$
Lx_c = \int_0^L x \, dx = \left[\frac{x^2}{2}\right]_0^L = \frac{L^2}{2}
$$

$$
\therefore \qquad x_c = \frac{L}{2}
$$

Thus the centroid lies at midpoint of a straight line, whatever be the orientation of line (Ref. Fig. 4.15).

From equation (*i*) and (*ii*) we can get the centroid of semicircle shown in Fig. 4.17 by putting $\alpha = \pi/2$ and for quarter of a circle shown in Fig. 4.18 by putting α varying from zero to $\pi/2$.

4.5 in the form

$$
x_c L = \Sigma L_i x_i
$$

$$
y_c L = \Sigma L_i y_i
$$

may be used to get centroid x_c and y_c . If the line segments is in space the expression $z_c L = \sum L_i z_i$ may also be used. The method is illustrated with few examples below:

Example 4.1. Determine the centroid of the wire shown in Fig. 4.19.

Solution: The wire is divided into three segments *AB, BC* and *CD*. Taking *A* as origin the coordinates of the centroids of *AB, BC* and *CD* are

*G*₁(300, 0); *G*₂(600, 100) and *G*₃ (600 – 150 cos 45°, 200 + 150 sin 45°) *i.e.,* G_3 (493.93, 306.07) $L_1 = 600$ mm, $L_2 = 200$ mm, $L_3 = 300$ mm ∴ Total length $L = 600 + 200 + 300 = 1100$ mm ∴ From the eqn. $Lx_c = \sum L_i x_i$, we get $1100 x_c = L_1 x_1 + L_2 x_2 + L_3 x_3$ $= 600 \times 300 + 200 \times 600 + 300 \times 493.93$ ∴ $x_c = 407.44$ mm Now, $Ly_c = \Sigma L_i y_i$ 1100 $y_c = 600 \times 0 + 200 \times 100 + 300 \times 306.07$ ∴ $y_c = 101.66$ mm i.e., $G_2(493,306.07)$

Content L = 600 + 200 + 300 = 1100 mm

Content L = 600 + 200 + 300 = 1100 mm

Content L = 600 + 200 + 300 = 1100 mm

Contex 30 = 14, we get
 $x_1 = 407.44$ mm

Now, $L_y = 2L_y$
 $y = 600 \times 100 + 300 \times$

Example 4.2. Locate the centroid of the uniform wire bent as shown in Fig. 4.20.

Solution: The composite figure is divided into three simple figures and taking *A* as origin coordinates of their centroids noted down as shown below:

AB—a straight line

$$
L_1 = 400 \text{ mm}, \qquad G_1 \text{ (200, 0)}
$$

BC—a semicircle

$$
L_2 = 150 \pi = 471.24, \qquad G_2 \left(475, \frac{2 \times 150}{\pi} \right)
$$

i.e. G_2 (475, 95.49)

CD—a straight line

 $L_3 = 250$; $x_3 = 400 + 300 + \frac{250}{2}$ cos 30° = 808.25 mm $y_3 = 125 \sin 30^\circ = 62.5 \text{ mm}$ ∴ Total length $L = L_1 + L_2 + L_3 = 1121.24$ mm $Lx_c = \sum L_i x_i$ *xi* gives 1121.24 $x_c = 400 \times 200 + 471.24 \times 475 + 250 \times 808.25$ $x_c = 451.20$ mm $Ly_c = \Sigma L_i y_i$ *yi* gives

1121.24 $y_c = 400 \times 0 + 471.24 \times 95.49 + 250 \times 62.5$

yc **= 54.07 mm**

Example 4.3. Locate the centroid of uniform wire shown in Fig. 4.21. Note: portion AB is in x-z plane, BC in y-z plane and CD in x-y plane. AB and BC are semi circular in shape.

Table 4.1

	A100	$r = 140$	0 $\overline{45^\circ}$ y	
		Fig. 4.21		
			Solution: The length and the centroid of portions AB, BC and CD are as shown in table below:	
		Table 4.1		
Portion	L_i	\boldsymbol{x}_i	y_i	Z_i
AB	100π	100	0	2×100
				π
BC	140π	$\mathbf 0$	140	2×140 π
CD	300	300 sin 45°	280 + 300 cos 45° $= 492.13$	0
$\ddot{\cdot}$	From eqn. $Lx_c = \sum L_i x_i$, we get $x_c = 90.19$ mm	$L = 100\pi + 140\pi + 300 = 1053.98$ mm 1053.98 $x_c = 100\pi \times 100 + 140\pi \times 0 + 300 \times 300 \sin 45^\circ$		
		Similarly, 1053.98 $y_c = 100\pi \times 0 + 140\pi \times 140 + 300 \times 492.13$		
	$y_c = 198.50$ mm			
		1053.98 $z_c = 100\pi \times \frac{200}{\pi} + 140\pi \times \frac{2 \times 140}{\pi} + 300 \times 0$		
and	$z_c = 56.17$ mm			
		4.4 FIRST MOMENT OF AREA AND CENTROID		
From equation 4.1, we have		$x_c = \frac{\sum W_i x_i}{W}$, $y_c = \frac{\sum W_i y_i}{W}$ and $z_c = \frac{\sum W_i z_i}{W}$		

$$
x_c = 90.19 \text{ mm}
$$

$$
y_c = 198.50 \text{ mm}
$$

4.4 FIRST MOMENT OF AREA AND CENTROID

$$
x_c = \frac{\Sigma W_i x_i}{W}, \quad y_c = \frac{\Sigma W_i y_i}{W} \quad \text{and} \quad z_c = \frac{\Sigma W_i z_i}{W}
$$

From the above equation we can make the statement that distance of centre of gravity of a body from an axis is obtained by dividing moment of the gravitational forces acting on the body, about the axis, by the total weight of the body. Similarly from equation 4.4, we have,

$$
x_c = \frac{\sum A_i x_i}{A}, \quad y_c = \frac{\sum A_i y_i}{A}
$$

By terming $\Sigma A_i x$: as the moment of area about the axis, we can say centroid of plane area from any axis is equal to moment of area about the axis divided by the total area. The moment of area $\Sigma A_i x$: is termed as first moment of area also just to differentiate this from the term $\Sigma A_i x_i^2$, which will be dealt latter. It may be noted that since the moment of area about an axis divided by total area gives the distance of the centroid from that axis, the moment of area is zero about any centroidal axis. $x_c = \frac{\sum k_x}{A}$, $y_c = \frac{\sum k_x}{A}$, $y_c = \frac{\sum k_y}{A}$

askis is equal to moment of area about the axis, we can say centroid of plane axis is equal to moment of area about the six divided by the total area. The moment of area as

Difference between Centre of Gravity and Centroid

From the above discussion we can draw the following differences between centre of gravity and centroid:

- (1) The term centre of gravity applies to bodies with weight, and centroid applies to lines, plane areas and volumes.
- (2) Centre of gravity of a body is a point through which the resultant gravitational force (weight) acts for any orientation of the body whereas centroid is a point in a line plane area volume such that the moment of area about any axis through that point is zero.

Use of Axis of Symmetry

Centroid of an area lies on the axis of symmetry if it exits. This is useful theorem to locate the centroid of an area.

This theorem can be proved as follows:

Consider the area shown in Fig. 4.22. In this figure *y*-*y* is the axis of symmetry. From eqn. 4.4, the distance of centroid from this axis is given by:

$$
\frac{\sum A_i x_i}{A}
$$

Consider the two elemental areas shown in Fig. 4.22, which are equal in size and are equidistant from the axis, but on either side. Now the sum of moments of these areas cancel each other since the areas and distances are the same, but signs of distances are opposite. Similarly, we can go on considering an area on one side of symmetric axis and

corresponding image area on the other side, and prove that total moments of area $(\Sigma A_i x_i)$ about the symmetric axis is zero. Hence the distance of centroid from the symmetric axis is zero, *i.e.,* centroid always lies on symmetric axis.

Making use of the symmetry we can conclude that:

- (1) Centroid of a circle is its centre (Fig. 4.23);
- (2) Centroid of a rectangle of sides *b* and *d* is at distance 2 $\frac{b}{a}$ and 2 $\frac{d}{dx}$ from the corner as shown in Fig. 4.24.

Determination of Centroid of Simple Figures From First Principle

For simple figures like triangle and semicircle, we can write general expression for the elemental area and its distance from an axis. Then equation 4.4 reduces to:

$$
\overline{y} = \frac{\int y dA}{A}
$$

$$
\overline{x} = \frac{\int x dA}{A}
$$

The location of the centroid using the above equations may be considered as finding centroid from first principle. Now, let us find centroid of some standard figures from first principle.

Centroid of a Triangle

Consider the triangle *ABC* of base width *b* and height *h* as shown in Fig. 4.25. Let us locate the distance of centroid from the base. Let b_1 be the width of elemental strip of thickness dy at a distance *y* from the base. Since Δ*AEF* and Δ*ABC* are similar triangles, we can write:

$$
\frac{b_1}{b} = \frac{h - y}{h}
$$

$$
b_1 = \left(\frac{h - y}{h}\right)b = \left(1 - \frac{y}{h}\right)b
$$

∴ Area of the element

$$
= \left(1 - \frac{y}{h}\right) b \ dy
$$

 $= dA = b_1 dv$

Area of the triangle

 $\frac{1}{2}$ bh

∴ From eqn. 4.4

 \overline{y} = $\frac{\text{Moment of area}}{\text{Total area}}$ $=\frac{\int y dA}{A}$ Now, $\int y dA = \int_0^b$ $\int_{0}^{h} y \left(1 - \frac{y}{h}\right) b \, dy$

$$
= \int_{0}^{h} \left(y - \frac{y^{2}}{h} \right) b \, dy
$$

$$
= b \left[\frac{y^{2}}{2} - \frac{y^{3}}{3h} \right]_{0}^{h}
$$

$$
= \frac{bh^{2}}{6}
$$

$$
\therefore \qquad \overline{y} = \frac{\int y dA}{A} = \frac{bh^{2}}{6} \times \frac{1}{\frac{1}{2}bh}
$$

$$
\therefore \qquad \overline{y} = \frac{h}{3}
$$

Thus the centroid of a triangle is at a distance $\frac{h}{3}$ from the base (or $\frac{2}{3}$ 3 $\frac{h}{\sqrt{2}}$ from the apex) of the triangle, where *h* is the height of the triangle.

Centroid of a Semicircle

Consider the semicircle of radius *R* as shown in Fig. 4.26. Due to symmetry centroid must lie on *y* axis. Let its distance from diametral axis be \bar{y} . To find \bar{y} , consider an element at a distance *r* from the centre O of the semicircle, radial width being dr and bound by radii at θ and $\theta + d\theta$.

Area of element = $r d\theta dr$.

Its moment about diametral axis x is given by:

 $r d\theta \times dr \times r \sin \theta = r^2 \sin \theta dr d\theta$

∴ Total moment of area about diametral axis,

$$
= b\left[\frac{y^{2}}{2} - \frac{y^{2}}{3h}\right]_{0}
$$
\n
$$
= \frac{bh^{2}}{6}
$$
\n
$$
\therefore \qquad \bar{y} = \frac{h}{3}bh^{2} \qquad \bar{y} = \frac{bh^{2}}{6} \qquad \frac{1}{2}bh
$$
\n
$$
\therefore \qquad \bar{y} = \frac{h}{3}
$$
\nThus the centroid of a triangle is at a distance $\frac{h}{3}$ from the base (or $\frac{2h}{3}$ from the a triangle, where *h* is the height of the triangle.\n\n**Centre of a Semicircle**\nConsider the semicircle of radius *R* as shown in Fig. 4.26. Due to symmetry centroid m axis. Let its distance from diameter a axis be \bar{y} . To find \bar{y} , consider an element at a distate the centre *O* of the semicircle, radial width being *dr* and bound by radii at θ and $\theta + d\theta$. Area of element = *r d\theta dr*.\n\nIt is moment about diametral axis, \bar{x} is given by:\n
$$
r d\theta \times dr \times r \sin \theta = r^{2} \sin \theta dr d\theta
$$
\n
$$
\therefore \text{Total moment of area about diameral axis,}
$$
\n
$$
\int_{0}^{\pi} \int_{0}^{a} r^{2} \sin \theta dr d\theta = \int_{0}^{\pi} \left[\frac{r^{3}}{3}\right]_{0}^{\pi} \sin \theta d\theta
$$
\n
$$
= \frac{R^{3}}{3} \left[-\cos \theta\right]_{0}^{\pi}
$$
\n
$$
= \frac{R^{3}}{3} \left[1 + 1\right] = \frac{2R^{3}}{3}
$$
\nArea of semicircle\n
$$
A = \frac{1}{2} \pi R^{2}
$$
\n
$$
\bar{y} = \frac{\text{Moment of area}}{\text{Total area}} = \frac{\frac{2R^{3}}{3}}{\frac{1}{2} \pi R^{2}}
$$
\n
$$
= \frac{4R}{3\pi}
$$

 $d\theta$ r θ dr O \leftarrow $R \rightarrow$ X Y R **Fig. 4.26**

Area of semicircle

2 πR^2

$$
\overline{y} = \frac{\text{Moment of area}}{\text{Total area}} = \frac{\frac{2R^3}{3}}{\frac{1}{2}\pi R^2}
$$

$$
= \frac{4R}{3\pi}
$$

Thus, the centroid of the circle is at a distance $\frac{4}{3}$ 3π *^R* from the diametral axis.

Centroid of Sector of a Circle

Consider the sector of a circle of angle 2α as shown in Fig. 4.27. Due to symmetry, centroid lies on *x* axis. To find its distance from the centre *O*, consider the elemental area shown.

Area of the element =*rd*θ *dr*

Its moment about *y* axis

The distance of centroid from centre *O*

$$
= \frac{\text{Moment of area about y axis}}{\text{Area of the figure}}
$$

$$
= \frac{2R^3}{3} \sin \alpha = \frac{2R}{3\alpha} \sin \alpha
$$

Centroid of Parabolic Spandrel

Consider the parabolic spandrel shown in Fig. 4.28. Height of the element at a distance *x* from *O* is $y = kx^2$

Centroids of some common figures are shown in Table 4.2.

Table 4.2 Centroid of Some Common Figures

Centroid of Composite Sections

So far, the discussion was confined to locating the centroid of simple figures like rectangle, triangle, circle, semicircle, etc. In engineering practice, use of sections which are built up of many simple sections is very common. Such sections may be called as built-up sections or composite sections. To locate the centroid of composite sections, one need not go for the first principle (method of integration). The given composite section can be split into suitable simple figures and then the centroid of each simple figure can be found by inspection or using the standard formulae listed in Table 4.2. Assuming the area of the simple figure as concentrated at its centroid, its moment about an axis can be

DISTRIBUTED FORCES, CENTRE OF GRAVITY AND MOMENT OF INERTIA **111**

found by multiplying the area with distance of its centroid from the reference axis. After determining moment of each area about reference axis, the distance of centroid from the axis is obtained by dividing total moment of area by total area of the composite section.

Example 4.4. Locate the centroid of the T-section shown in the Fig. 4.29.

Solution: Selecting the axis as shown in Fig. 4.29, we can say due to symmetry centroid lies on *y* axis, *i.e.* $\bar{x} = 0$. Now the given *T*-section may be divided into two rectangles A_1 and A_2 each of size 100 \times 20 and 20 \times 100. The centroid of A_1 and A_2 are $g_1(0, 10)$ and $g_2(0, 70)$ respectively.

∴ The distance of centroid from top is given by:

$$
\overline{y} = \frac{100 \times 20 \times 10 + 20 \times 100 \times 70}{100 \times 20 + 20 \times 100}
$$

 $= 40$ mm

Hence, **centroid of T-section is on the symmetric axis at a distance 40 mm from the top.** *Example 4.5. Find the centroid of the unequal angle 200* \times *150* \times *12 mm, shown in Fig. 4.30. Solution:* The given composite figure can be divided into two rectangles:

 $A_1 = 150 \times 12 = 1800$ mm² $A_2 = (200 - 12) \times 12 = 2256$ mm² Total area $A = A_1 + A_2 = 4056$ mm²

Selecting the reference axis x and y as shown in Fig. 4.30. The centroid of A_1 is g_1 (75, 6) and that of A_2 is:

> g_2 6,12 + $\frac{1}{2}$ (200 – 12) $\left[\frac{0,12+2(200-12)}{2}\right]$

i.e., $g_2 (6, 106)$

$$
\therefore \quad \overline{x} = \frac{\text{Moment about y axis}}{\text{Total area}}
$$
\n
$$
= \frac{A_1 x_1 + A_2 x_2}{A}
$$
\n
$$
= \frac{1800 \times 75 + 2256 \times 6}{4056} = 36.62 \text{ mm}
$$
\n
$$
\overline{y} = \frac{\text{Moment about x axis}}{\text{Total area}}
$$
\n
$$
= \frac{A_1 y_1 + A_2 y_2}{A}
$$
\n
$$
= \frac{1800 \times 6 + 2256 \times 106}{4056} = 61.62 \text{ mm}
$$

Thus, the centroid is at \bar{x} = 36.62 mm and \bar{y} = 61.62 mm as shown in the figure.

Example 4.6. Locate the centroid of the I-section shown in Fig. 4.31.

Solution: Selecting the coordinate system as shown in Fig. 4.31, due to symmetry centroid must lie on *y* axis,

i.e.,
$$
\overline{x} = 0
$$

Now, the composite section may be split into three rectangles

$$
A_1 = 100 \times 20 = 2000
$$
 mm²

Centroid of A_1 from the origin is:

$$
y_1 = 30 + 100 + \frac{20}{2} = 140
$$
 mm
 $A_2 = 100 \times 20 = 2000$ mm²

Similarly

Similarly
\n
$$
A_2 = 100 \times 20 = 2000 \text{ mm}^2
$$

\n $y_2 = 30 + \frac{100}{2} = 80 \text{ mm}$
\n $A_3 = 150 \times 30 = 4500 \text{ mm}^2$,
\nand
\n $y_3 = \frac{30}{2} = 15 \text{ mm}$
\n $\overline{y} = \frac{A_1 y_1 + A_2 y_2 + A_3 y_3}{A}$
\n $= \frac{2000 \times 140 + 2000 \times 80 + 4500 \times 15}{2000 + 2000 + 4500}$
\n= 59.71 mm

Thus, **the centroid is on the symmetric axis at a distance 59.71 mm from the bottom as shown in Fig. 4.31.**

Example 4.7. Determine the centroid of the section of the concrete dam shown in Fig. 4.32.

Fig. 4.32

Solution: Let the axis be selected as shown in Fig. 4.32. Note that it is convenient to take axis in such a way that the centroids of all simple figures are having positive coordinates. If coordinate of any simple figure comes out to be negative, one should be careful in assigning the sign of moment of area of that figure.

The composite figure can be conveniently divided into two triangles and two rectangles, as shown in Fig. 4.32.

Now,
\n
$$
A_1 = \frac{1}{2} \times 2 \times 6 = 6 \text{ m}^2
$$
\n
$$
A_2 = 2 \times 7.5 = 15 \text{ m}^2
$$
\n
$$
A_3 = \frac{1}{2} \times 3 \times 5 = 7.5 \text{ m}^2
$$
\n
$$
A_4 = 1 \times 4 = 4 \text{ m}^2
$$
\n
$$
A = \text{total area} = 32.5 \text{ m}^2
$$

Centroids of simple figures are:

$$
x_1 = \frac{2}{3} \times 2 = \frac{4}{3} \text{ m}
$$

\n
$$
y_1 = \frac{1}{3} \times 6 = 2 \text{ m}
$$

\n
$$
x_2 = 2 + 1 = 3 \text{ m}
$$

\n
$$
y_2 = \frac{7.5}{2} = 3.75 \text{ m}
$$

\n
$$
x_3 = 2 + 2 + \frac{1}{3} \times 3 = 5 \text{ m}
$$

$$
y_3 = 1 + \frac{1}{3} \times 5 = \frac{8}{3} \text{ m}
$$

\n $x_4 = 4 + \frac{4}{2} = 6 \text{ m}$
\n $y_4 = 0.5 \text{ m}$
\n $\overline{x} = \frac{A_1x_1 + A_2x_2 + A_3x_3 + A_4x_4}{A}$
\n $= \frac{6 \times \frac{4}{3} + 15 \times 3 + 7.5 \times 5 + 4 \times 6}{32.5}$
\n $= 3.523 \text{ m}$
\n $\overline{y} = \frac{A_1y_1 + A_2y_2 + A_3y_3 + A_4y_4}{A}$
\n $= \frac{6 \times 2 + 15 \times 3.75 + 7.5 \times \frac{5}{3} + 4 \times 0.5}{32.5}$
\nand
\n $\overline{y} = 2.777 \text{ m}$
\nThe centroid is at
\n $\overline{x} = 3.523 \text{ m}$
\nand
\nExample 4.8. Determine the centroid of the area shown in Fig. 4.33 with respect to the axis shown.
\n \overrightarrow{y}
\n \overrightarrow{y}
\n \overrightarrow{z}
\n \overrightarrow{z} <

Solution: The composite section is divided into three simple figures, a triangle, a rectangle and a semicircle

Now, area of triangle 2 \times 3 \times 4 = 6 m² Area of rectangle $A_2 = 6 \times 4 = 24$ m² Area of semicircle 2 $\times \pi \times 2^2 = 6.2832$ m² ∴ Total area *A* = 36.2832 m²

The coordinates of centroids of these three simple figures are:

=

 6×4

 \times

$$
x_1 = 6 + \frac{1}{3} \times 3 = 7 \text{ m}
$$

\n
$$
y_1 = \frac{4}{3} \text{ m}
$$

\n
$$
x_2 = 3 \text{ m}
$$

\n
$$
x_3 = \frac{-4R}{3\pi} = -\frac{4 \times 2}{3\pi} = -0.8488 \text{ m}
$$

\n
$$
y_3 = 2 \text{ m}
$$

\n
$$
x_3 = \frac{4R}{3\pi} = -\frac{4 \times 2}{3\pi} = -0.8488 \text{ m}
$$

\n
$$
\overline{x} = \frac{A_1x_1 + A_2x_2 + A_3x_3}{A}
$$

\n
$$
= \frac{6 \times 7 + 24 \times 3 + 6.2832 \times (-0.8488)}{36.2832}
$$

\ni.e.,
\n
$$
\overline{x} = 2.995 \text{ m}
$$

\n
$$
\overline{y} = \frac{A_1y_1 + A_2y_2 + A_3y_3}{A}
$$

\n
$$
= \frac{\frac{6 \times 4}{3} + 24 \times 2 + 6.2832 \times 2}{36.2832}
$$

\ni.e.,
\nExample 4.9. In a guest plate, there are six river holes of 21.5 mm diameter as shown in
\nFind the position of the centroid of the gusset plate.
\n
$$
\overline{y} = \frac{4 \times 2}{\sqrt{130}} = \frac{200}{\sqrt{130}} = \frac{
$$

i.e., $\bar{y} = 1.890 \text{ m}$

Example 4.9. In a gusset plate, there are six rivet holes of 21.5 mm diameter as shown in Fig. 4.34. Find the position of the centroid of the gusset plate.

 $\frac{x}{3}$ + 24 \times 2 + 6.2832 \times 2 36.2832

 $\frac{x}{3}$ + 24 \times 2 + 6.2832 \times

Solution: The composite area is equal to a rectangle of size 160 × 280 mm *plus* a triangle of size 280 mm base width and 40 mm height and *minus* areas of six holes. In this case also the Eqn. 4.4 can be used for locating centroid by treating area of holes as negative. The area of simple figures and their centroids are as shown in Table 4.3.

∴ $A = \sum A_i = 48221.70$

$$
\therefore \qquad \Sigma A_i x_i = 44800 \times 140 + 5600 \times \frac{560}{3} - 363.05 (70 + 140 + 210 + 70 + 140 + 210)
$$

$$
= 7012371.3 \text{ mm}^3
$$

$$
\overline{x} = \frac{\sum A_i x_i}{A} = 145.42 \text{ mm}
$$

\n
$$
\sum A_i y_i = 44800 \times 80 + 5600 \times 173.33 - 363.05 (50 \times 3 + 120 + 130 + 140)
$$

\n= 4358601 mm³
\n
$$
\overline{y} = \frac{\sum A_i y_i}{A} = \frac{4358601}{48221.70}
$$

\n= 90.39 mm

Thus, the coordinates of centroid of composite figure is given by:

 \bar{x} = 145.42 mm *y* **= 90.39 mm**

Example 4.10. Determine the coordinates x_c *and* y_c *of the centre of a 100 mm diameter circular hole cut in a thin plate so that this point will be the centroid of the remaining shaded area shown in Fig. 4.35 (All dimensions are in mm).*

Solution: If x_c and y_c are the coordinates of the centre of the circle, centroid also must have the coordinates x_c and *yc* as per the condition laid down in the problem. The shaded area may be considered as a rectangle of size 200 mm \times 150 mm *minus* a triangle of sides 100 mm \times 75 mm and a circle of diameter 100 mm.

$$
\therefore \quad \text{Total area} \\
 = 200 \times 150 - \frac{1}{2} \times 100 \times 75 - \left(\frac{\pi}{4}\right) 100^2 \\
 = 18396 \text{ mm}^2
$$

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$$
x_c = \bar{x} = \frac{200 \times 150 \times 100 - \frac{1}{2} \times 100 \times 75 \times \left[200 - \left(\frac{100}{3}\right)\right] - \frac{\pi}{4} \times 100^2 \times x_c}{18396}
$$

\n
$$
\therefore \qquad x_c(18396) = 200 \times 150 \times 100 - \frac{1}{2} \times 100 \times 75 \times 166.67 - \frac{\pi}{4} \times 100^2 x_c
$$

\n
$$
x_c = \frac{2375000}{26250} = 90.48 \text{ mm}
$$

\nSimilarly,

18396
$$
y_c = 200 \times 150 \times 75 - \frac{1}{2} \times 100 \times 75 \times (150 - 25) - \frac{\pi}{4} \times 100^2 y_c
$$

$$
y_c = \frac{1781250.0}{26250} = 67.86 \text{ mm}
$$

Centre of the circle should be located at (90.48, 67.86) so that this point will be the centroid of the remaining shaded area as shown in Fig. 4.35.

Note: The centroid of the given figure will coincide with the centroid of the figure without circular hole. Hence, the centroid of the given figure may be obtained by determining the centroid of the figure without the circular hole also.

Example 4.11. Determine the coordinates of the centroid of the plane area shown in Fig. 4.36 with reference to the axis shown. Take x = 40 mm.

Solution: The composite figure is divided into the following simple figures:

(1) A rectangle
$$
A_1 = (14x) \times (12x) = 168x^2
$$

$$
x_1 = 7x; \ y_1 = 6x
$$

(2) A triangle
$$
A_2 = \frac{1}{2} (6x) \times (4x) = 12x^2
$$

 $x_2 = 14x + 2x = 16x$
 $y_2 = \frac{4x}{3}$

(3) A rectangle to be subtracted

$$
A_3 = (-4x) \times (4x) = -16x^2
$$

$$
x_3 = 2x; y_3 = 8x + 2x = 10x
$$

(4) A semicircle to be subtracted

$$
A_4 = -\frac{1}{2}\pi (4x)^2 = -8\pi x^2
$$

$$
x_4 = 6x
$$

$$
y_4 = \frac{4R}{3\pi} = 4 \times \frac{4(x)}{3\pi} = \frac{16x}{3\pi}
$$

(5) A quarter of a circle to be subtracted

$$
y_2 = \frac{4x}{3}
$$

\n(3) A rectangle to be subtracted
\n $A_3 = (-4x) \times (4x) = -16x^2$
\n $x_3 = 2x; y_3 = 8x + 2x = 10x$
\n(4) A semicircle to be subtracted
\n $A_4 = -\frac{1}{2}\pi (4x)^2 = -8\pi x^2$
\n $x_4 = 6x$
\n $y_4 = \frac{4R}{3\pi} = 4 \times \frac{4(x)}{3\pi} = \frac{16x}{3\pi}$
\n(5) A quarter of a circle to be subtracted
\n $A_5 = -\frac{1}{4} \times \pi (4x)^2 = -4\pi x^2$
\n $x_5 = 14x - \frac{4R}{3\pi} = 14x - (4)(\frac{4x}{3\pi}) = 12.3023x$
\nTotal area
\n $A = 168x^2 + 12x^2 - 16x^2 - 8\pi x^2 - 4\pi x^2$
\n $= 126.3009x^2$
\n $\bar{x} = \frac{\Sigma 4x}{A}$
\n $\Sigma A_{\bar{x}_1} = 168x^2 \times 7x + 12x^2 \times 16x - 16x^2 \times 2x - 8\pi x^2 \times 6x - 4\pi x^2 \times 12.3023x$
\n $= 1030.6083x^3$
\n $\bar{x} = \frac{1030.6083x^3}{126.3009x^2}$
\n= 8.1599x = 8.1599x 40 (since x = 40 mm)
\n $\bar{y} = \frac{\Sigma A_1 y_1}{A}$

$$
\Sigma A_i y_i = 168x^2 \times 6x + 12x^2 \times \frac{4x}{3} - 16x^2 \times 10x
$$

\n
$$
-8\pi x^2 \times \frac{16x}{3\pi} - 4\pi x^2 \times 10.3023x
$$

\n
$$
= 691.8708x^3
$$

\n
$$
\vec{y} = \frac{691.8708x^3}{126.3009x^2}
$$

\n
$$
= 5.4780x
$$

\n
$$
= 219.12 \text{ mm}
$$
 (since $x = 40 \text{ mm}$)
\n
$$
\text{(since } x = 40 \text{ mm})
$$

Centroid

4.5 SECOND MOMENTS OF PLANE AREA

Consider the area shown in Fig. 4.37(*a*). *dA* is an elemental area with coordinates as *x* and *y*. The term $\Sigma y_i^2 dA_i$ is called *moment of inertia* of the area about *x* axis and is denoted as I_{xx} . Similarly, the moment of inertia about *y* axis is

 $I_{yy} = \sum x_i^2 dA_i$

In general, if *r* is the distance of elemental area *dA* from the axis *AB* [Fig. 4.37(*b*)], the sum of the terms Σ*r* ² *dA* to cover the entire area is called moment of inertia of the area about the axis *AB*. If *r* and *dA* can be expressed in general term, for any element, then the sum becomes an integral. Thus,

The term *rdA* may be called as moment of area, similar to moment of a force, and hence *r* ² *dA* may be called as moment of area or the second moment of area. Thus, the moment of inertia of area is nothing but second moment of area. In fact, the term *'second moment of area'* appears to correctly signify the meaning of the expression Σr^2 *dA*. The term 'moment of inertia' is rather a misnomer. However, the term moment of inertia has come to stay for long time and hence it will be used in this book also.

Though moment of inertia of plane area is a purely mathematical term, it is one of the important properties of areas. The strength of members subject to bending depends on the moment of inertia of its cross-sectional area. Students will find this property of area very useful when they study subjects like strength of materials, structural design and machine design.

The moment of inertia is a fourth dimensional term since it is a term obtained by multiplying area by the square of the distance. Hence, in SI units, if metre (m) is the unit for linear measurements used then $m⁴$ is the unit of moment of inertia. If millimetre (mm) is the unit used for linear measurements, then mm⁴ is the unit of moment of inertia. In MKS system m^4 or cm^4 and in FPS system ft⁴ or in⁴ are commonly used as units for moment of inertia.

Polar Moment of Inertia

Moment of inertia about an axis perpendicular to the plane of an area is known as *polar moment of inertia*. It may be denoted as J or I_{zz} . Thus, the moment of inertia about an axis perpendicular to the plane of the area at *O* in Fig. 4.38 is called polar moment of inertia at point *O*, and is given by

$$
I_{zz} = \Sigma r^2 dA \tag{4.7}
$$

Radius of Gyration

Radius of gyration is a mathematical term defined by the relation

$$
k = \sqrt{\frac{I}{A}}
$$

where $k =$ radius of gyration,

 $I =$ moment of inertia,

and $A =$ the cross-sectional area

Suffixes with moment of inertia *I* also accompany the term radius of gyration *k*. Thus, we can have,

and so on.

The relation between radius of gyration and moment of inertia can be put in the form:

 $k_{xx} = \sqrt{\frac{I_{xx}}{A}}$

 $k_{yy} = \sqrt{\frac{I_{yy}}{A}}$

 $k_{AB} = \sqrt{\frac{I_{AB}}{A}}$

$$
I = Ak^2 \qquad \qquad \dots (4.9)
$$

From the above relation a geometric meaning can be assigned to the term 'radius of gyration.' We can consider *k* as the distance at which the complete area is squeezed and kept as a strip of negligible width (Fig. 4.39) such that there is no change in the moment of inertia. **Polar Moment of Inertia**

Moment of intertia about an axis perpendicular to the plane of an area, is

Morown as *polar moment of inertia*. It may be denoted as *I* or I_x . Thus, the

moment of inertia about an axis perp

...(4.8)

Theorems of Moments of Inertia

There are two theorems of moment of inertia:

- (1) Perpendicular axis theorem, and
- (2) Parallel axis theorem.

These are explained and proved below.

Perpendicular Axis Theorem

The moment of inertia of an area about an axis pependicular to its plane (polar moment of inertia) at any point O is equal to the sum of moments of inertia about any two mutually perpendicular axis through the same point O and lying in the plane of the area. y

Referring to Fig. 4.40, if *z-z* is the axis normal to the plane of paper passing through point *O*, as per this theorem,

$$
I_{zz} = I_{xx} + I_{yy}
$$
 ... (4.10)
The above theorem can be easily proved. Let us consider an elemental
area dA at a distance r from O. Let the coordinates of dA be x and y.
Then from definition:

$$
I_{zz} = \Sigma r^2 dA
$$

= $\Sigma (x^2 + y^2) dA$
= $\Sigma x^2 dA + \Sigma y^2 dA$
 $I_{zz} = I_{xx} + I_{yy}$

Parallel Axis Theorem

Moment of inertia about any axis in the plane of an area is equal to the sum of moment of inertia about a parallel centroidal axis and the product of area and square of the distance between the two parallel axis. Referring to Fig. 4.41 the above theorem means:

$$
I_{AB} = I_{GG} + A y_c^2
$$
...(4.11)

where

 I_{AB} = moment of inertia about axis *AB*

- I_{GG} = moment of inertia about centroidal axis *GG* parallel to *AB*.
	- $A =$ the area of the plane figure given and
- y_c = the distance between the axis *AB* and the parallel centroidal axis *GG*.
- *Proof:* Consider an elemental parallel strip *dA* at a distance *y* from the centroidal axis (Fig. 4.41).

(2) Parallel axis theorem.
\nThese are explained and proved below.
\n**Perpendicular Axis Theorem**
\nThe moment of inertia of an area about an axis perpendicular to its plane (polar moment of
\nany point O is equal to the sum of moments of inertia about any two mutually perpendicular
\nthrough the same point O and lying in the plane of the area.
\nRefering to Fig. 4.40, if
$$
z-z
$$
 is the axis normal to the plane of
\npaper passing through point O, as per this theorem,
\n
$$
I_{zz} = I_{xx} + I_{yy}
$$
\n
$$
I_{zz} = I_{xx} + I_{yy}
$$
\nThe above theorem can be easily proved. Let us consider an elemental
\narea dA at a distance r from O. Let the coordinates of dA be x and y.
\nThen from definition:
\n
$$
I_{zz} = \Sigma r^2 dA
$$
\n
$$
= \Sigma (x^2 + y^2) dA
$$
\n
$$
= \Sigma x^2 dA + \Sigma y^2 dA
$$
\n**Parallel Axis Theorem**
\n*Moment of inertia about any axis in the plane of an area
\nis equal to the sum of moment of inertia about a parallel
\ncentroid at axis and the product of area and square of the
\ndistance between the two parallel axis. Referring to
\nFig. 4.41 the above theorem means:
\n
$$
I_{AB} = \text{moment of inertia about a axis AB}
$$
\n
$$
I_{GG} = \text{moment of inertia about the axis AB}
$$
\n
$$
I_{GG} = \text{moment of inertia about the axis AB}
$$
\n
$$
A = \text{the area of the plane figure given and}
$$
\n
$$
y_c = \text{the distance between the axis AB and the parallel centroid}
$$
\n
$$
\text{axis GG.}
$$
\n
$$
Pref. 4.41
$$
\n
$$
Pref. 4.41
$$
\n
$$
Pref. 4.42
$$
\n
$$
Pref. 4.43
$$
\n
$$
I_{AB} = \Sigma (y^2 + y)^2 dA
$$
\n
$$
= \Sigma y^2 dA + \Sigma y y_c dA + \Sigma y^2 dA
$$
\n
$$
= \Sigma y^2 dA = \text{Moment of inertia about the axis GG}
$$*

 $= I_{GG}$

$$
\Sigma 2yy_c\,dA = 2y_c\,\Sigma y\,dA
$$

$$
= 2y_c A \frac{\Sigma y dA}{A}
$$

In the above term $2y_c A$ is constant and $\frac{\Sigma y dA}{A}$ $\frac{\Sigma y dA}{\Sigma y}$ is the distance of centroid from the reference axis

GG. Since *GG* is passing through the centroid itself *ydA* $\frac{dA}{A}$ is zero and hence the term $\Sigma 2yy_c dA$ is zero.

Now, the third term,

$$
\Sigma y_c^2 dA = y_c^2 \Sigma dA
$$

= Ay_c^2

$$
\therefore \qquad I_{AB} = I_{GG} + A y_c^2
$$

Note: The above equation cannot be applied to any two parallel axis. One of the axis (*GG*) must be centroidal axis only.

4.6 MOMENT OF INERTIA FROM FIRST PRINCIPLES

For simple figures, moment of inertia can be obtained by writing the general expression for an element and then carrying out integration so as to cover the entire area. This procedure is illustrated with the following three cases:

- (1) Moment of inertia of a rectangle about the centroidal axis
- (2) Moment of inertia of a triangle about the base
- (3) Moment of inertia of a circle about a diametral axis

(1) *Moment of Inertia of a Rectangle about the Centroidal Axis*: Consider a rectangle of width *b* and depth *d* (Fig. 4.42). Moment of inertia about the centroidal axis *x-x* parallel to the short side is required.

Consider an elemental strip of width *dy* at a distance *y* from the axis. Moment of inertia of the elemental strip about the centroidal axis *xx* is:

(2) *Moment of Inertia of a Triangle about its Base*: Moment of inertia of a triangle with base width *b* and height *h* is to be determined about the base *AB* (Fig. 4.43).

Consider an elemental strip at a distance *y* from the base *AB*. Let *dy* be the thickness of the strip and *dA* its area. Width of this strip is given by:

$$
b_1 = \frac{(h - y)}{h} \times b
$$

Moment of inertia of this strip about *AB*

$$
= y2 dA
$$

= $y2 b1 dy$
= $y2 \frac{(h - y)}{h} \times b \times dy$

∴ Moment of inertia of the triangle about *AB*,

Consider an elemental strip at a distance y from
\nthe base AB. Let dy be the thickness of the strip and
\n
$$
dA
$$
 its area. Width of this strip is given by:
\n
$$
b_1 = \frac{(h - y)}{h} \times b
$$
\n
$$
b_1 = \frac{(h - y)}{h} \times b
$$
\n
$$
= y^2 dA
$$
\n
$$
= y^2 h_1 dy
$$
\n
$$
= y^2 \frac{(h - y)}{h} \times b \times dy
$$
\n
$$
\therefore \text{ Moment of inertia of the triangle about AB,}
$$
\n
$$
I_{AB} = \int_0^h \frac{y^2 (h - y) b dy}{h}
$$
\n
$$
= \int_0^h \left(y^2 - \frac{y^3}{h} \right) b dy
$$
\n
$$
= \int_0^h \left(y^2 - \frac{y^3}{h} \right) b dy
$$
\n
$$
= \int_0^h \left(y^2 - \frac{y^4}{h} \right) b dy
$$
\n
$$
= \int_0^h \left(\frac{y^3}{3} - \frac{y^4}{4h} \right) \Big|_0^h
$$
\n
$$
= \int_0^h \left(\frac{h^3}{3} - \frac{h^4}{4h} \right) \Big|_0^h
$$
\n
$$
= \int_0^h \left(\frac{h^3}{3} - \frac{h^4}{4h} \right) \Big|_0^h
$$
\n
$$
= \int_0^h \left(\frac{h^3}{3} - \frac{h^4}{4h} \right) \Big|_0^h
$$
\n(3) *Moment of* inertia of a circle of radius R is required about its Diamert al
\nAxius: When end of inertia of a circle of radius R is required about R is required
\nabout t's diameter in this about the diameteral axis x-x:
\n
$$
= y^2 dA
$$
\n
$$
= (r \sin \theta) \sin \theta \, d\theta \, dr
$$
\n
$$
\therefore \text{ Moment of inertia of the circle about x-x is given by}
$$
\n
$$
I_{xx} = \int_0^h \left(\frac{h}{h} \right)^2 r^3 \sin^2 \theta \, d\theta \, dr
$$

(3) *Moment of Inertia of a Circle about its Diametral Axis*: Moment of inertia of a circle of radius *R* is required about it's diametral axis as shown in Fig. 4.44

Consider an element of sides *rd*θ and *dr* as shown in the figure. It's moment of inertia about the diametral axis *x-x*:

$$
= y2 dA
$$

= $(r \sin \theta)^2 r d\theta dr$
= $r^3 \sin^2 \theta d\theta dr$

∴ Moment of inertia of the circle about *x*-*x* is given by

$$
I_{xx} = \int_{0}^{R} \int_{0}^{2\pi} r^3 \sin^2 \theta \ d\theta \ dr
$$

Fig. 4.44

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$$
= \int_{0}^{R} \int_{0}^{2\pi} \frac{(1 - \cos 2\theta)}{2} d\theta dr
$$

$$
= \int_{0}^{R} \frac{r^{3}}{2} \left[\theta - \frac{\sin 2\theta}{2} \right]_{0}^{2\pi} dr
$$

$$
= \left[\frac{r^{4}}{8} \right]_{0}^{R} [2\pi - 0 + 0 - 0] = \frac{2\pi}{8} R^{4}
$$

$$
I_{xx} = \frac{\pi R^{4}}{4}
$$

If *d* is the diameter of the circle, then

$$
R = \frac{d}{2}
$$

\n
$$
I_{xx} = \frac{\pi}{4} \left(\frac{d}{2}\right)^4
$$

\n
$$
I_{xx} = \frac{\pi d^4}{64}
$$

d

Moment of Inertia of Standard Sections

Rectangle: Referring to Fig. 4.45.

(a)
$$
I_{xx} = \frac{bd^3}{12}
$$
 as derived from first principle.

(b)
$$
I_{yy} = \frac{db^3}{12}
$$
 can be derived on the same lines.

(*c*) About the base *AB*, from parallel axis theorem, $I_{AB} = I_{xx} + A y_c^2$

= 3 12 *bd* + *bd* 2 2 *d* , since *yc* = 2 *d* = 3 12 *bd* + 3 4 *bd* = 3 3 *bd* Civil[datas.blogspo](http://Civildatas.blogspot.in)t.in

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Now, y_c , the distance between the non-centroidal axis *AB* and centroidal axis *x*-*x*, is equal to $\frac{h}{3}$.

$$
\therefore \frac{bh^3}{12} = I_{xx} + \frac{1}{2}bh\left(\frac{h}{3}\right)
$$

$$
= I_{xx} + \frac{bh^3}{18}
$$

$$
\therefore \qquad I_{xx} = \frac{bh^3}{12} - \frac{bh^3}{18}
$$

$$
= \frac{bh^3}{36}
$$

Moment of Inertia of a Circle about any diametral axis

$$
= \frac{\pi d^4}{64} \tag{as found}
$$

2

d from first principle)

Moment of Inertia of a Hollow Circle: Referring to Fig. 4.48. I_{AB} = Moment of inertia of solid circle of diameter *D* about *AB* – Moment of inertia of circle of diameter *d* about *AB*. That is,

$$
= \frac{\pi D^4}{64} - \frac{\pi d^4}{64}
$$

$$
= \frac{\pi}{64} (D^4 - d^4)
$$

Moment of Inertia of a Semicircle: (*a*) *About Diametral Axis*: If the limit of integration is put as 0 to π instead of 0 to 2π in the derivation for the moment of inertia of a circle about diametral axis the moment of inertia of a semicircle is obtained. It can be observed that the moment of inertia of a semicircle (Fig. 4.49) about the diametral axis *AB*: $=\frac{\pi D^4}{64} - \frac{\pi d^4}{64}$
 $=\frac{\pi}{64} (D^4 - a^4)$
 $=\frac{\pi}{64} (D^4 - a^4)$

Moment of there in a denotical equivalent is to the solution of the minimal of a Smitteria of a simulated Attack (a) About Diameteral Attack and the env

$$
= \frac{1}{2} \times \frac{\pi d^4}{64} = \frac{\pi d^4}{128}
$$

(*b*) *About Centroidal Axis x-x*:

Now, the distance of centroidal axis y_c from the diametral axis is given by:

$$
y_c = \frac{4R}{3\pi} = \frac{2d}{3\pi}
$$

and,

$$
\text{Area } A = \frac{1}{2} \times \frac{\pi d^2}{4} = \frac{\pi d^2}{8}
$$

From parallel axis theorem,

$$
I_{AB} \stackrel{.}{=} I_{xx} + Ay_c^2
$$

$$
\frac{\pi d^4}{128} = I_{xx} + \frac{\pi d^2}{8} \times \left(\frac{2d}{3\pi}\right)^2
$$

$$
I_{xx} = \frac{\pi d^4}{128} - \frac{d^4}{18\pi}
$$

$$
= 0.0068598 \ d^4
$$

Moment of Inertia of a Quarter of a Circle: (*a*) *About the Base*: If the limit of integration is put as 0

to $\frac{\pi}{4}$ 2 instead of 0 to 2π in the derivation for moment of inertia of a circle the moment of inertia of a

quarter of a circle is obtained. It can be observed that moment of inertia of the quarter of a circle about the base *AB*.

 $X \mid X$ G

 $4R$ $\overline{3\pi}$

A B R

Fig. 4.50

(*b*) *About Centroidal Axis x-x*:

Now, the distance of centroidal axis y_c from the base is given by:

$$
y_c = \frac{4R}{3\pi} = \frac{2d}{3\pi}
$$

and the area $1 \times \pi d^2$ πd^2 4 4 16 $\times \frac{\pi d^2}{4} = \frac{\pi d}{4}$

From parallel axis theorem,

$$
I_{AB} = I_{xx} + Ay_c^2
$$

\n
$$
\frac{\pi d^4}{256} = I_{xx} + \frac{\pi d^2}{16} \left(\frac{2d}{3\pi}\right)^2
$$

\n
$$
I_{xx} = \frac{\pi d^4}{256} - \frac{d^4}{36\pi}
$$

\n= 0.00343 d⁴

The moment of inertia of common standard sections are presented in Table 4.4.

Table 4.4 Moment of Inertia of Standard Sections

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4.7 MOMENT OF INERTIA OF COMPOSITE SECTIONS

Beams and columns having composite sections are commonly used in structures. Moment of inertia of these sections about an axis can be found by the following steps:

- (1) Divide the given figure into a number of simple figures.
- (2) Locate the centroid of each simple figure by inspection or using standard expressions.
- (3) Find the moment of inertia of each simple figure about its centroidal axis. Add the term *Ay*² where *A* is the area of the simple figure and y is the distance of the centroid of the simple figure from the reference axis. This gives moment of inertia of the simple figure about the reference axis.
- (4) Sum up moments of inertia of all simple figures to get the moment of inertia of the composite section.

The procedure given above is illustrated below. Referring to the Fig. 4.51, it is required to find out the moment of inertia of the section about axis *A-B.*

(1) The section in the figure is divided into a rectangle, a triangle and a semicircle. The areas of the simple figures A_1 , A_2 and A_3 are calculated.

(2) The centroids of the rectangle (g_1) , triangle (g_2) and semicircle (g_3) are located. The distances y_1 , y_2 and y_3 are found from the axis *AB*.

(3) The moment of inertia of the rectangle about it's centroid (I_{g_1}) is calculated using standard expression. To this, the term A_1 y_1^2 is added to get the moment of inertia about the axis *AB* as: $I_1 = I_{g_1} + A_1 y_1^2$

Similarly, the moment of inertia of the triangle $(I_2 = I_{g_2} + A_2 y_2^2)$ and of semicircle $(I_3 = I_{g_3}$ $+A_3 y_3^2$ about axis *AB* are calculated.

(4) Moment of inertia of the composite section about *AB* is given by:

$$
I_{AB} = I_1 + I_2 + I_3
$$

= $I_{g_1} + A_1 y_1^2 + I_{g_2} + A_2 y_2^2 + I_{g_3} + A_3 y_3^2$...(4.12)

In most engineering problems, moment of inertia about the centroidal axis is required. In such cases, first locate the centroidal axis as discussed in 4.4 and then find the moment of inertia about this axis.

Referring to Fig. 4.52, first the moment of area about any reference axis, say *AB* is taken and is divided by the total area of section to locate centroidal axis *x-x*. Then the distances of centroid of individual figures y_{c1} , y_{c2} and y_{c3} from the axis $x-x$ are determined. The moment of inertia of the composite section about the centroidal axis $x-x$ is calculated using the expression:

$$
I_{xx} = I_{g_1} + A_1^2 y_{c1} + I_{g_2} + A_2^2 y_{c2} + I_{g_3} + A_3^2 y_{c3}
$$
...(4.13)
\n
$$
I_{xx} = I_{g_1} + A_1^2 y_{c1} + I_{g_2} + A_2^2 y_{c2} + I_{g_3} + A_3^2 y_{c3}
$$
\n
$$
I_{g_3}
$$
\n
$$
I_{g_4}
$$
\n
$$
I_{g_5}
$$
\n
$$
I_{g_6}
$$
\n
$$
I_{g_7}
$$
\n
$$
I_{g_8}
$$
\n
$$
I_{g_8}
$$
\n
$$
I_{g_9}
$$
\n
$$
I_{g_1} = I_{g_1} + A_3^2
$$
\n
$$
I_{g_1} = I_{g_1} + A_3^2
$$
\n
$$
I_{g_2} = I_{g_4} + A_3^2
$$
\nwhere A is the moment of inertia *i* at a constant, *ii* at a constant, *iii* at a constant, *iv* at a constant, *iv*

Sometimes the moment of inertia is found about a conveninet axis and then using parallel axis theorem, the moment of inertia about centroidal axis is found.

In the above example, the moment of inertia I_{AB} is found and y_c , the distance of *CG* from axis *AB* is calculated. Then from parallel axis theorem,

$$
I_{AB} = I_{xx} + Ay_c^2
$$

$$
I_{xx} = I_{AB} - Ay_c^2
$$

where *A* is the area of composite section.

Example 4.12. Determine the moment of inertia of the section shown in Fig. 4.53 about an axis passing through the centroid and parallel to the top most fibre of the section. Also determine moment of inertia about the axis of symmetry. Hence find radii of gyration.

Solution: The given composite section can be divided into two rectangles as follows:

Area $A_1 = 150 \times 10 = 1500$ mm² Area $A_2 = 140 \times 10 = 1400$ mm² Total Area $A = A_1 + A_2 = 2900$ mm² Due to symmetry, centroid lies on the symmetric axis *y-y*.

The distance of the centroid from the top most fibre is

given by:

Sum of moment of the areas about the top most fibre Total area $=\frac{1500 \times 5 + 1400(10 + 70)}{0000}$ 2900 $= 41.21$ mm

Referring to the centroidal axis *x*-*x* and *y*-*y*, the centroid of A_1 is g_1 (0.0, 36.21) and that of A_2 is *g*2 (0.0, 38.79).

Moment of inertia of the section about *x*-*x* axis

 I_{xx} = moment of inertia of A_1 about *x*-*x* axis + moment of inertia of A_2 about *x*-*x* axis.

$$
\therefore I_{xx} = \frac{150 \times 10^3}{12} + 1500 (36.21)^2 + \frac{10 \times 140^3}{12} + 1400 (38.79)^2
$$

i.e.,
$$
I_{xx} = 6372442.5 \text{ mm}^4
$$

Similarly,

$$
I_{yy} = \frac{10 \times 150^3}{12} + \frac{140 \times 10^3}{12} = 2824,166.7
$$
 mm⁴

Hence, the moment of inertia of the section about an axis passing through the centroid and parallel to the top most fibre is 6372442.5 mm⁴ and moment of inertia of the section about the axis of symmetry is 2824166.66 mm⁴.

The radius of gyration is given by:

Example 4.13. Determine the moment of inertia of the L-section shown in the Fig. 4.54 about its centroidal axis parallel to the legs. Also find out the polar moment of inertia.

Solution: The given section is divided into two rectangles *A*¹ and A_2 .

Area
$$
A_1 = 125 \times 10 = 1250
$$
 mm²

Area $A_2 = 75 \times 10 = 750$ mm²

Total Area = 2000 mm²

First, the centroid of the given section is to be located.

Two reference axis (1) – (1) and (2) – (2) are chosen as shown in Fig. 4.54.

The distance of centroid from the axis (1) – (1)

$$
= \frac{\text{sum of moment of areas } A_1 \text{ and } A_2 \text{ about}(1) - (1)}{\text{T}_{11}}
$$

$$
\overline{x} = \frac{120 \times 5 + 750 \left(10 + \frac{75}{2} \right)}{2000}
$$

 $= 20.94$ mm

Similarly,

the distance of the centroid from the axis (2) – (2)

$$
= \overline{y} = \frac{1250 \times \frac{125}{2} + 750 \times 5}{2000} = 40.94 \text{ mm}
$$

With respect to the centroidal axis *x*-*x* and *y*-*y*, the centroid of A_1 is g_1 (15.94, 21.56) and that of *A*2 is *g*2 (26.56, 35.94).

 I_{xx} = Moment of inertia of A_1 about *x-x* axis + Moment of inertia of A_2 about *x-x* axis

$$
\therefore I_{xx} = \frac{10 \times 125^3}{12} + 1250 \times 21.56^2 + \frac{75 \times 10^3}{12} + 750 \times 39.94^2
$$

i.e., $I_{rr} = 3411298.9$ mm⁴

$$
\frac{1}{2}
$$

Similarly,

$$
I_{yy} = \frac{125 \times 10^3}{12} + 1250 \times 15.94^2 + \frac{10 \times 75^3}{12} + 750 \times 26.56^2
$$

i.e., $I_{yy} = 1208658.9$ mm⁴ Polar moment of inertia $= I_{xx} + I_{yy}$

 $= 3411298.9 + 12,08658.9$

Izz **= 4619957.8 mm⁴**

Example 14. Determine the moment of inertia of the symmertic I-section shown in Fig. 4.55 about its centroidal axis x-x and y-y.

Also, determine moment of inertia of the section about a centroidal axis perpendicular to x-x axis and y-y axis. Solution: The section is divided into three rectangles A_1 , A_2 and A_3 .

Area $A_1 = 200 \times 9 = 1800$ mm²

Area $A_2 = (250 - 9 \times 2) \times 6.7 = 1554.4$ mm² Area $A_3 = 200 \times 9 = 1800$ mm²

Total Area $A = 5154.4$ mm²

The section is symmetrical about both *x-x* and *y-y* axis. Therefore, its centroid will coincide with the centroid of rectangle A_2 .

With respect to the centroidal axis *x*-*x* and *y*-*y*, the centroid of rectangle A_1 is g_1 (0.0, 120.5), that of A_2 is g_2 $(0.0, 0.0)$ and that of A_3 is g_3 (0.0, 120.5).

i.e.,

 I_{xx} = Moment of inertia of A_1 + Moment of inertia of A_2 + Moment of inertia of *A*3 about *x-x* axis

$$
I_{xx} = \frac{200 \times 9^3}{12} + 1800 \times 120.5^2 + \frac{6.7 \times 232^3}{12} + 0
$$

$$
+ \frac{200 \times 9^3}{12} + 1800(120.5)^2
$$

$$
I_{xx} = 5,92,69,202 \text{ mm}^4
$$

Similarly,

$$
I_{xy} = \frac{9 \times 200^3}{12} + \frac{232 \times 6.7^3}{12} + \frac{9 \times 200^3}{12}
$$

$$
I_{yy} = 1,20,05,815 \text{ mm}^4
$$

Moment of inertia of the section about a centroidal axis perpendicular to *x*-*x* and *y*-*y* axis is nothing but polar moment of inertia, and is given by:

$$
I_{xx} = I_{xx} + I_{yy}
$$

= 59269202 + 12005815

$$
I_{yy} = 7,12,75,017
$$
mm⁴

Example 4.15. Compute the second moment of area of the channel section shown in Fig. 4.56 about centroidal axis
$$
x
$$
- x and y - y .

Solution: The section is divided into three rectangles A_1 , A_2 and A_3 . Area $A_1 = 100 \times 13.5 = 1350$ mm² Area $A_2 = (400 - 27) \times 8.1 = 3021.3$ mm²

Area $A_3 = 100 \times 13.5 = 1350.00$ mm² Total Area $A = 5721.3$ mm²

The given section is symmetric about horizontal axis passing through the centroid g_2 of the rectangle A_2 . A reference axis (1)–(1) is chosen as shown in Fig. 4.56.

The distance of the centroid of the section from (1)*–*(1)

$$
=\frac{1350 \times 50 + 3021.3 \times \frac{8.1}{2} + 1350 \times 50}{5721.3}
$$

 $= 25.73$ mm

With reference to the centroidal axis *x-x* and *y-y*, the centroid of the rectangle A_1 is $g_1(24.27, 193.25)$ that of A_2 is g_2 (21.68, 0.0) and that of A_3 is g_3 (24.27, 193.25).

Similarly,

$$
I_{xx} = \text{Moment of inertia of } A_1, A_2 \text{ and } A_3 \text{ about } x \text{-}x
$$
\n
$$
= \frac{100 \times 13.5^3}{12} + 1350 \times 193.25^2
$$
\n
$$
+ \frac{8.1 \times 373^3}{12} + \frac{100 \times 13.5^3}{12} + 1350 \times 193.25^2
$$
\n
$$
I_{xx} = 1.359 \times 10^8 \text{ mm}^4
$$
\nSimilarly,\n
$$
I_{yy} = \frac{13.5 \times 100^3}{12} + 1350 \times 24.27^2 + \frac{273 \times 8.1^3}{12} + 3021.3
$$
\n
$$
\times 21.68^2 + \frac{13.5 \times 100^3}{12} + 1350 \times 24.27^2
$$

Iyy **= 52,72557.6 mm⁴**

Example 4.16. Determine the polar moment of inertia of the I-section shown in the Fig. 4.57. Also determine the radii of gyration with respect to x-x axis and y-y axis.

Solution: The section is divided into three rectangles as shown in Fig. 4.57

Area
\nArea
\nArea
\nArea
\nArea
\n
$$
A_1 = 80 \times 12 = 960 \text{ mm}^2
$$

\nArea
\n $A_2 = (150 - 22) \times 12 = 1536 \text{ mm}^2$
\nTotal area
\n $A = 3696 \text{ mm}^2$
\nDue to symmetry, centroid lies on axis by by The bottom fibre

Due to symmetry, centroid lies on axis *y-y*. The bottom fibre (1)–(1) is chosen as reference axis to locate the centroid.

The distance of the centroid from (1) – (1)

$$
= \frac{\text{Sum of moments of the areas of the rectangles about } (1) - (1)}{\text{Total area of the rectangle}} = \frac{\text{Sum of moments of the area of the rectangle}}{\text{Total area of the triangle}} = \frac{1}{\frac{1}{2} \cdot \frac{1}{2} \cdot \frac
$$

Total area of section

$$
\frac{960 \times (150 - 6) + 1536 \times \left(\frac{128}{2} + 10\right) + 1200 \times 5}{3696}
$$

= 69.78 mm

=

 \overline{I}

With reference to the centroidal axis *x*-*x* and *y*-*y*, the centroid of the rectangles A_1 is g_1 (0.0, 74.22), that of A_2 is g_2 (0.0, 4.22) and that of A_3 is g_3 (0.0, 64.78).

$$
I_{xx} = \frac{80 \times 12^3}{12} + 960 \times 74.22^2 + \frac{12 \times 128^3}{12} + 1536 \times 4.22^2 + \frac{120 \times 10^3}{12} + 1200 \times 64.78^2
$$

$$
I_{xx} = 1,24,70,028 \text{ mm}^4
$$

$$
V_{yy} = \frac{12 \times 80^3}{12} + \frac{128 \times 12^3}{12} + \frac{10 \times 120^3}{12}
$$

$$
= 19,70,432
$$
 mm⁴

Polar moment of inertia $= I_{xx}$

$$
= 1,24,70,027 + 19,70,432
$$

$$
= 1,44,40459 \text{ mm}^4
$$

$$
k_{xx} = \sqrt{\frac{I_{xx}}{A}} = \sqrt{\frac{1,24,70,027}{3696}}
$$

$$
= 58.09 \text{ mm}
$$

$$
k_{yy} = \sqrt{\frac{I_{yy}}{A}} = \sqrt{\frac{19,70,432}{3696}}
$$

$$
= 23.09 \text{ mm}.
$$

Example 4.17. Determine the moment of inertia of the built-up section shown in Fig. 4.58 about its centroidal axis x-x and y-y.

Solution: The given composite section may be divided into simple rectangles and triangles as shown in the Fig. 4.58

Due to symmetry, centroid lies on the axis *y-y*. A reference axis (1) – (1) is choosen as shown in the figure. The distance of the centroidal axis from (1) – (1)

 \overline{y} = $\frac{\text{sum of moment of areas about (1)–(1)}}{T}$ Total area = $3000 \times 135 + 2500 \times 70 + 4000 \times 10 + 875 \left(\frac{1}{3} \times 20 + 20 \right) \times 2$ 11250

 $= 59.26$ mm

With reference to the centroidal axis *x-x* and *y-y*, the centroid of the rectangle A_1 is g_1 (0.0, 75.74), that of A_2 is $g_2(0.0, 10.74)$, that of A_3 is g_3 (0.0, 49.26), the centroid of triangle A_4 is g_4 (41.66, 32.59) and that of A_5 is g_5 (41.66, 32.59).

$$
\frac{1}{25}
$$

\n
$$
\frac{1}{250}
$$

Example 4.18. Determine the moment of inertia of the built-up section shown in the Fig. 4.59 about an axis AB passing through the top most fibre of the section as shown.

Solution: In this problem, it is required to find out the moment of inertia of the section about an axis *AB.* So there is no need to find out the position of the centroid.

The given section is split up into simple rectangles as shown in Fig. 4.59. Now,

Moment of inertia about *AB* = Sum of moments of inertia of the rectangle about *AB*

$$
= \frac{400 \times 20^3}{12} + 400 \times 20 \times 10^2 + \left[\frac{100 \times 10^3}{12} + 100 \times 10 \times (20 + 5)^2 \right] \times 2 + \left[\frac{100 \times 380^3}{12} + 10 \times 380 \times (30 + 190)^2 \right] \times 2 + \left[\frac{100 \times 10^3}{12} + 100 \times 10 \times (20 + 10 + 380 + 5)^2 \right] \times 2
$$

$$
I_{AB} = 8.06093 \times 10^8 \text{ mm}^4.
$$

Example 4.19. Calculate the moment of inertia of the built-up section shown in Fig. 4.60 about a centroidal axis parallel to AB. All members are 10 mm thick.

Solution: The built-up section is divided into six simple rectangles as shown in the figure.

The distance of centroidal axis from *AB*

$$
= \frac{\text{Sum of the moment of areas about } AB}{\text{Total area}}
$$

$$
= \frac{\sum A_i y_i}{A}
$$

Fig. 4.60

Now,
\n
$$
\Sigma A_i y_i = 250 \times 10 \times 5 + 2 \times 40 \times 10 \times (10 + 20) + 40 \times 10 \times (10 + 5)
$$
\n
$$
+ 40 \times 10 \times 255 + 250 \times 10 \times (10 + 125)
$$
\n
$$
= 4,82,000 \text{ mm}^3
$$
\n
$$
A = 2 \times 250 \times 10 + 40 \times 10 \times 4
$$
\n
$$
= 6600 \text{ mm}^2
$$
\n
$$
\therefore \qquad \overline{y} = \frac{\sum A_i y_i}{A} = \frac{482000}{6600}
$$
\n
$$
= 73.03 \text{ mm}
$$

Now,

 $Now,$

Moment of inertia about the centroidal axis $\begin{cases} \text{Sum of the moment of } \\ \text{of the individual rectangles} \end{cases}$ \int $\begin{cases} \end{cases}$

$$
= \frac{250 \times 10^3}{12} + 250 \times 10 \times (73.03 - 5)^2
$$

+
$$
\left[\frac{10 \times 40^3}{12} + 40 \times 10 (73.03 - 30)^2 \right] \times 2
$$

+
$$
\frac{40 \times 10^3}{12} + 40 \times 10 (73.03 - 15)^2 + \frac{10 \times 250^3}{12} + 250
$$

$$
\times 10 (73.03 - 135)^2 + \frac{40 \times 10^3}{12} + 40 \times 10 (73.03 - 255)^2
$$

$$
I_{xx} = 5,03,99,395 \text{ mm}^4.
$$

Example 4.20. A built-up section of structural steel consists of a flange plate 400 mm × 20 mm, a web plate 600 mm × 15 mm and two angles 150 mm × 150 mm × 10 mm assembled to form a section as shown in Fig. 4.61. Determine the moment of inertia of the section about the horizontal centroidal axis.

$$
A = 600 \times 15 + 140 \times 10 \times 2 + 150 \times 10 \times 2 + 400 \times 20
$$

= 22,800 mm²

$$
\therefore \qquad \overline{y} = \frac{\sum A_i y_i}{A} = \frac{3315000}{22800}
$$

$$
= 145.39 \text{ mm}
$$

Moment of inertia of the section about centroidal axis $\begin{cases} \text{Sum of the moments of } = \text{all simple figures about centroidal axis} \end{cases}$ \int $\begin{cases} \end{cases}$

$$
y = \frac{A}{A} = \frac{1}{22800}
$$

\n= 145.39 mm
\nMoment of inertia of the
\nsection about centroidal axis \int = $\left[\frac{Sum}{all \text{ simple figures about centroidal axis}}\right]$
\n= $\frac{15 \times 600^3}{12} + 600 \times 15(145.39 - 320)^2$
\n+ $\left[\frac{10 \times 140^3}{12} + 1400(145.39 - 100)^2\right] \times 2$
\n+ $\left[\frac{150 \times 10^3}{12} + 1500 \times (145.39 - 10)^2\right] \times 2$
\n+ $\frac{400 \times 20^3}{12} + 400 \times 20 \times (145.39 - 10)^2$
\n= $\frac{I_{xx}}{12} = 7.45156 \times 10^8$ mm⁴.
\nExample 4.21. Compute the moment of the 100 mm × 150 mm rectangle shown is
\nabout *x-x* axis to which it is inclined at an angle
\n $\theta = \sin^{-1} \left(\frac{4}{5}\right)$.
\nSolution: The rectangle is divided into four triangles as
\nshown in the figure. [The lines AE and FC are parallel to
\n*x*-axis].
\nNow $\theta = \sin^{-1} \left(\frac{4}{5}\right) = 53.13^{\circ}$
\nFrom the geometry of the Fig. 4.62,
\n $BK = AB \sin (90^{\circ} - \theta)$
\n= 100 sin (90^{\circ} - 53.13^{\circ})
\n= 60 mm
\n $ND = BK = 60$ mm
\n $FD = \frac{60}{\sin \theta} = \frac{60}{\sin 53.13} = 75$ mm

$$
I_{xx} = 7.45156 \times 10^8
$$
 mm⁴.

Example 4.21. Compute the moment of inertia of the 100 mm × 150 mm rectangle shown in Fig. 4.62 about x-x axis to which it is inclined at an angle

$$
\theta = \sin^{-1}\left(\frac{4}{5}\right).
$$

Solution: The rectangle is divided into four triangles as shown in the figure. [The lines *AE* and *FC* are parallel to *x*-axis].

Now
$$
\theta = \sin^{-1}\left(\frac{4}{5}\right) = 53.13^{\circ}
$$

From the geometry of the Fig. 4.62,

$$
BK = AB \sin (90^\circ - \theta)
$$

= 100 sin (90^\circ - 53.13^\circ)

$$
= 60 \text{ mm}
$$

$$
ND = BK = 60 \text{ mm}
$$

$$
FD = \frac{60}{\sin \theta} = \frac{60}{\sin 53.13} = 75 \text{ mm}
$$

∴ $AF = 150 - FD = 75$ mm Hence $FL = ME = 75 \sin \theta = 60 \text{ mm}$

$$
AE = FC = \frac{AB}{\cos(90^\circ - \theta)} = \frac{100}{0.8} = 125 \text{ mm}
$$

Moment of inertia of the section about *x-x* axis $\begin{cases} \text{Sum of the moments of inertia of indiv}_1\text{dual triangular areas about } x \text{-}x \text{ axis} \end{cases}$ \int $\left\{ \right.$

$$
AE = PC = \frac{1}{\cos(90^\circ - \theta)} = \frac{1}{0.8} = 125 \text{ mm}
$$

\nMoment of inertia of the
\nsection about *x-x* axis
$$
= \left[\frac{\text{Sum of the moments of inertia of in the\nsection about x-x axis
$$
= I_{DFC} + I_{FCZ} + I_{FEA} + I_{AEB}
$$
\n
$$
= \frac{125 \times 60^3}{36} + \frac{1}{2} \times 125 \times 60 \times \left(60 + \frac{1}{3} \times 60 \right)^2 + \frac{125 \times 60^3}{36} + \frac{1}{2} \times 125 \times 60^3 + \frac{1}{2} \times 125 \
$$
$$

Example 4.22. Find moment of inertia of the shaded area shown in the Fig. 4.63 about the axis AB.

Fig. 4.63

Solution: The section is divided into a triangle *PQR,* a semicircle *PSQ* having base on axis *AB* and a circle having its centre on axis *AB.*

Now,

Moment of inertia of the
\nsection about axis
$$
AB
$$

\n
$$
\begin{cases}\n\begin{cases}\n\text{Moment of inertia of triangle } PQR \text{ about } AB + \text{Moment of inertia of semicircle} \\
\hline\nPSQ \text{ about } AB - \text{moment of inertia of } \\
\text{circle about } AB\n\end{cases}
$$
\n
$$
= \frac{80 \times 80^3}{12} + \frac{\pi}{128} \times 80^4 - \frac{\pi}{64} \times 40^4
$$
\n
$$
I_{AB} = 42,92,979 \text{ mm}^4.
$$

Example 4.23. Find the second moment of the shaded portion shown in the Fig. 4.64 about its centroidal axis.

Solution: The section is divided into three simple figures viz., a triangle *ABC*, a rectangle *ACDE* and a semicircle.

section about axis AB =
$$
\frac{1}{2}
$$
 - $\frac{1}{128}$ × 80⁴ - $\frac{\pi}{64}$ × 40⁴
\n
$$
I_{AB} = 42,92,979 \text{ mm}^4.
$$
\nExample 4.23. Find the second moment of the shaded portion shown in the Fig. 4.64 about its
\ncentroidal axis.
\n
$$
\overline{x} = \begin{bmatrix} 10 & 10 & 10 \\ 10 & 10 & 10
$$

$$
= 132203.6
$$

\n
$$
\overline{x} = \frac{A\overline{x}}{A} = \frac{132203.6}{3371.68} = 37.21 \text{ mm}
$$

\nMoment of inertia about
\n
$$
\begin{aligned}\n&\text{Moment of inertia of triangle } ABC \text{ about } \\&\text{entroidal } x \cdot x \text{ axis} + \text{Moment of a rectangle} \\&\text{about } x \cdot x \text{ axis} - \text{moment of semicircle}\n\end{aligned}
$$
\n
$$
\begin{aligned}\n\therefore \qquad I_{xx} = \frac{80 \times 20^3}{36} + \frac{1}{2} \times 80 \times 20 \left(60 - \frac{2}{3} \times 20 - 28.47 \right)^2 \\
&+ \frac{80 \times 40^3}{12} + 80 \times 40 \times (28.47 - 20)^2 \\
&- \left[0.0068598 \times 20^4 + \frac{1}{2} \pi \times 20^2 \left(28.47 - \frac{4 \times 20}{3 \pi} \right)^2 \right] \\
I_{xx} = 6,86,944 \text{ mm}^4.\n\end{aligned}
$$
\nSimilarly,
\n
$$
I_{yy} = \frac{20 \times 30^3}{36} + \frac{1}{2} \times 20 \times 30 \left(39.21 - \frac{2}{3} \times 30 \right)^2 + \frac{20 \times 50^3}{36} + \frac{1}{2} \times 20 \times 50 \times \left[39.21 - \left(30 + \frac{1}{3} \times 50 \right) \right]^2 + \frac{40 \times 80^3}{12} + 40 \times 80(39.21 - 40)^2 - \frac{1}{2} \times \frac{\pi}{64} \times 40^4 - \frac{1}{2} \times \frac{\pi}{4} \times 40^2 (40 - 39.21)^2 + 1868392 \text{ mm}^4.
$$
\n4.8 **THEOREMS OF PAPPUS-GULIMDING**
\nThere are two important theorems, first proposed by Greek scientist (about 340 AD) and then restated by Swiss mathematician Paul Gallotinus (1640) for determining the surface area and volumes generated by rotating a curve and a plane area about a non-intersect

4.8 THEOREMS OF PAPPUS-GULDINUS

There are two important theorems, first proposed by Greek scientist (about 340 AD) and then restated by Swiss mathematician Paul Guldinus (1640) for determining the surface area and volumes generated by rotating a curve and a plane area about a non-intersecting axis, some of which are shown in Fig. 4.65. These theorems are known as Pappus-Guldinus theorems.

Theorem I

The area of surface generated by revolving a plane curve about a non-intersecting axis in the plane of the curve is equal to the length of the generating curve times the distance travelled by the centroid of the curve in the rotation.

Proof: Figure 4.66 shows the isometric view of the plane curve rotated about *x-*axis by angle θ. We are interested in finding the surface area generated by rotating the curve *AB*. Let *dL* be the elemental length on the curve at *D*. Its coordinate be *y*. Then the elemental surface area generated by this element at *D*

Fig. 4.66

Thus we get area of the surface generated as length of the generating curve times the distance travelled by the centroid.

Theorem II

The volume of the solid generated by revolving a plane area about a non-intersecting axis in the plane is equal to the area of the generating plane times the distance travelled by the centroid of the plane area during the rotation.

Proof: Consider the plane area *ABC,* which is rotated through an angle θ about *x-*axis as shown in Fig. 4.67.

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Let *dA* be the elemental area of distance *y* from *x*-axis. Then the volume generated by this area during rotation is given by

 $dV = dA/vθ$ ∴ *V = dA/y*θ $= \theta \int y dA$ = θ *A yc* $= A(y_c \theta)$

Thus the volume of the solid generated is area times the distance travelled by its centroid during the rotation. Using Pappus-Guldinus theorems surface area and volumes of cones and spheres can be calculated as shown below:

(*i*) *Surface area of a cone*: Referring to Fig. 4.68(*a*), Length of the line generating cone $= L$

Distance of centroid of the line from the axis of rotation $= y = 2$ *R*

In one revolution centroid moves by distance = $2\pi y = \pi R$

∴ Surface area = *L* × (π*R*) = π*RL*

(*ii*) *Volume of a cone*: Referring to Fig. 4.68(*b*),

Area generating solid cone = $\frac{1}{2}$ 2 *hR*

Centroid *G* is at a distance $y = \frac{R}{3}$

∴ The distance moved by the centroid in one revolution = $2πy = 2π \frac{1}{3}$ *R*

∴ Volume of solid cone = $\frac{1}{2}$ 2 $hR \times \frac{2}{3}$ 3 πR = 2 3 $\pi R^2 h$

(*iii*) *Surface area of sphere*: Sphere of radius *R* is obtained by rotating a semi circular arc of radius *R* about its diametral axis. Referring to Fig. 4.69(*a*),

Length of the arc = πR

- Centroid of the arc is at $y = \frac{2R}{\pi}$ from the diametral axis (i.e. axis of rotation)
- ∴ Distance travelled by centroid of the arc in one revolution

$$
=2\pi y=2\pi \frac{2R}{\pi}=4R
$$

- ∴ Surface area of sphere = π*R* × 4*R* $= 4\pi R^2$
- (*iv*) *Volume of sphere*: Solid sphere of radius *R* is obtained by rotating a semicircular area about its diametral axis. Referring to Fig. 4.69(*b*).

Area of semicircle =
$$
\frac{\pi R^2}{2}
$$

Distance of centroid of semicircular area from its centroidal axis

$$
= y = \frac{4R}{3\pi}
$$

∴ The distance travelled by the centroid in one revolution

Centroid of the arc is at
$$
y = \frac{2R}{\pi}
$$
 from the diamertal axis (i.e. axis of rotation)

\n7. Distance travelled by centroid of the arc in one revolution

\n
$$
= 2\pi y = 2\pi \frac{2R}{\pi} = 4R
$$
\n8. Surface area of sphere $y = \pi R \times 4R$

\n(iv) Volume of sphere: Solid sphere of radius R is obtained by rotating a semicircular area about its diamertal axis. Referring to Fig. 4.69(b).

\nArea of semicircle $= \frac{\pi R^2}{2}$

\nDistance of centroid of semicircular area from its centroid axis

\n
$$
= y = \frac{4R}{3\pi}
$$
\n9. The distance traveled by the centroid in one revolution

\n
$$
= 2\pi y = 2\pi \frac{4R}{3\pi} = \frac{8R}{3}
$$
\n10. Volume of sphere $= \frac{\pi R^2}{2} \times \frac{8R}{3}$

\n
$$
= \frac{4\pi R^3}{3}
$$
\n11. Volume of sphere $= \frac{\pi R^2 \times 8R}{2 \times 3}$

\n
$$
= \frac{4\pi R^3}{3}
$$
\n12. Figure OF GRAVITY OF SOLIDS

\n13. The distance to gravity of solids may be found using eqn. (4.1) which will be same as those found from eqns. (4.2) and (4.3) if the mass is uniform. Hence center of gravity of solids, centre of gravity of mass or centroid of volumes is the same for all solids with uniform mass. For standard solids, the

4.9 CENTRE OF GRAVITY OF SOLIDS

Centre of gravity of solids may be found using eqn. (4.1) which will be same as those found from eqns. (4.2) and (4.3) if the mass is uniform. Hence centre of gravity of solids, centre of gravity of mass or centroid of volumes is the same for all solids with uniform mass. For standard solids, the centre of gravity may be found from first principle and the results obtained for standard solids may be used to find centre of gravity of composite solids. The procedure is illustrated with examples 4.24 to 4.27.

Example 4.24. Locate the centre of gravity of the right circular cone of base radius r and height h shown in Fig. 4.70.

Solution: Taking origin at the vertex of the cone and selecting the axis as shown in Fig. 4.70, it can be observed that due to symmetry the coordinates of centre of gravity \bar{y} and \bar{z} are equal to zero, *i.e.* the centre of gravity lies on the axis of rotation of the cone. To find its distance \bar{x} from the vertex, consider an elemental plate at a distance *x.* Let the thickness of the elemental plate be *dx*. From the similar triangles *OAB* and *OCD,* the radius of elemental plate *z* is given by

$$
z = \frac{x}{h}r
$$

∴ Volume of the elemental plate *dv*

$$
dv = \pi z^2 dx = \pi x^2 \frac{r^2}{h^2} dx
$$

If γ is the unit weight of the material of the cone, then weight of the elemental plate is given by:

Example 4.24. From the *n* is the *n*th term. The equation is given by the formula for the equation *n* is the *n*th term. The equation is given by the equation
$$
x = \frac{x}{n}
$$
 and *n*th term. The equation is given by the equation $x = \frac{1}{n} \int_{0}^{\pi} \frac{dx}{h^2} dx$ is the *n*th term. The equation is given by the equation $x = \frac{x}{n}$ and *n*th term. The equation is $x = \frac{x}{n}$ is the *n*th term. The equation is $x = \frac{x}{n}$ is the *n*th term. The equation is $x = \frac{x}{n}$ is the *n*th term. The equation is $x = \frac{x}{n}$ is the *n*th term. The equation is $x = \frac{x}{n}$ is the *n*th term. The equation is $x = \frac{x}{n}$ is the *n*th term. The equation is $x = \frac{x}{n}$ is the *n*th term. The equation is $x = \frac{x}{n}$ and *n*th term. The equation is $x = \frac{x}{n}$ is the *n*th term. The equation is $x = \frac{x}{n}$ and *n*th term. The equation is $x = \frac{x}{n}$ and *n*th term. The equation is $x = \frac{x}{n}$ and *n*th term. The equation is $x = \frac{x}{n}$ and *n*th term. The equation is $x = \frac{x}{n}$ and *n*th term. The equation is $x = \frac{x}{n}$ and *n*th term. The equation is $x = \frac{x}{n}$ and *n*th term. The equation is $x = \frac{x}{n}$ and *n*th term. The equation is $x = \frac{x}{n}$ and *n*th term. The equation is $x = \frac{x}{n}$ and *n*th term. The equation is $x = \frac{x}{n}$ and *n*th term. The

Note: $\frac{\pi r^2 h}{3}$ is volume of cone |
| $\overline{}$ $\overline{}$

Now, substituting the value of *dW* in (*i*), above, we get:

$$
\int x \cdot dW = \int_0^h \gamma \frac{\pi r^2}{h^2} x^2 \cdot x \cdot dx
$$

...(*iii*)

$$
= \gamma \frac{\pi r^2}{h^2} \left[\frac{x^4}{4} \right]_0^h
$$

$$
= \gamma \frac{\pi r^2 h^2}{4}
$$

 $2h^2$

 γ πr²h

From eqn. 4.1,

i.e.,

3 $\frac{\pi r^2 h}{2} \bar{x}$ = 4 ∴ $\bar{x} = \frac{3}{4}$ 4 *h*

2

 $W\overline{x} = \int x \ dW$

Thus, in a right circular cone, centre of gravity lies at a distance $\frac{3}{4}$ 4 *h* from vertex along the axis of rotation *i.e.*, at a distance $\frac{h}{4}$ from the base. $= \frac{\pi r^2 h^2}{4}$

From eqn. 4.1,
 $W\overline{x} = \int x dW$
 i.e.
 $\frac{\pi r^2 h^2}{3}x = \frac{7\pi r^2 h^2}{4}$
 $\therefore \qquad \overline{x} = \frac{3}{4}h$

Thus, in a right circular come, centre of gravity lies at a distance $\frac{3}{4}h$ from vertex alo

of rotat

Example 4.25. Determine the centre of gravity of a solid hemisphere of radius r from its diametral axis.

Solution: Due to symmetry, centre of gravity lies on the axis of rotation. To find its distance \bar{x} from the base along the axis of rotation, consider an elemental plate at a distance *x* as shown in Fig. 4.71.

Now,
$$
x^2 + z^2 = r^2
$$

Volume of elemental plate

dv = $πz^2 dx = π(r^2 - x^2)$

∴ Weight of elemental plate

$$
dW = \gamma dv = \gamma \pi (r^2 - x^2) dx
$$

∴ Weight of hemisphere

$$
W = \int dW = \int_{0}^{r} \gamma \pi (r^2 - x^2) dx
$$

$$
= \gamma \pi \left[r^2 x - \frac{x^2}{3} \right]_{0}^{r}
$$

$$
= 2\gamma \pi r^3
$$

 $z^2 = r^2 - x^2$...(*i*) $\dots(ii)$ \dots (*iii*) \circ \overline{x} \overline{x} \overline{x} \overline{x} r z dw z **Fig. 4.71**

...(*iv*)

Moment of weight about *z* axis

$$
= \int_{0}^{r} x dW
$$

3

=

$$
= \int_{0}^{r} x \pi (r^{2} - x^{2}) dx
$$

= $\pi \left[r^{2} \frac{x^{2}}{2} - \frac{x^{4}}{4} \right]_{0}^{r}$
= $\frac{\pi r^{4}}{4}$...(v)

 \overline{x} , the distance of centre of gravity from base is given by:

$$
W\overline{x} = \int\limits_0^r x \, dw
$$

i.e., From (*iv*) and (*v*) above, we get

$$
\frac{2\gamma\pi r^3}{3} \overline{x} = \frac{\gamma\pi r^4}{4} \qquad \overline{x} = \frac{3}{8}r
$$

Thus, the centre of gravity of a solid hemisphere of radius *r* is at a distance $\frac{3}{8}$ 8 *r* from its

diametral axis.

Example 4.26. Determine the maximum height h of the cylindrical portion of the body with hemispherical base shown in Fig. 4.72 so that it is in stable equilibrium on its base.

Solution: The body will be stable on its base as long as its centre of gravity is in hemispherical base. The limiting case is when it is on the plane *x-x* shown in the figure.

Centroid lies on the axis of rotation.

Mass of cylindrical portion

$$
m_1 = \pi r^2 h \rho
$$
, where ρ is unit mass of material.
Its centre of gravity g_1 is at a height

 $z_1 = \frac{h}{2}$ from *x* axis.

Mass of hemispherical portion

$$
m_2 = \rho \frac{2\pi r^3}{3}
$$

and its *CG* is at a distance

$$
z_2 = \frac{3r}{8}
$$
 from x-x plane.

Since centroid is to be on *x*-*x* plane $\overline{z} = 0$

i.e.,
$$
\Sigma m_i z_i = 0
$$

$$
\frac{m_1 h}{2} - m_2 \frac{3}{8} r = 0
$$

$$
\pi r^2 h \rho \frac{h}{2} = \rho \frac{2\pi r^3}{3} \frac{3}{8} r
$$

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 $h^2 = \frac{1}{2}$ 2 *r*2 or $h = \frac{\mathbf{r}}{l}$ **2 = 0.707** *r*

Example 4.27. A concrete block of size 0.60 m × 0.75 m × 0.5 m is cast with a hole of diameter 0.2 m and depth 0.3 m as shown in Fig. 4.73. The hole is completely filled with steel balls weighing 2500 N. Locate the centre of gravity of the body. Take the weight of concrete = 25000 N/m³ .

Solution: Weight of solid concrete block:

 $W_1 = 0.6 \times 0.75 \times 0.5 \times 25000 = 5625$ N Weight of concrete (W_2) removed for making hole:

$$
W_2 = \frac{\pi}{4} \times 0.2^2 \times 0.3 \times 25000 = 235.62 \text{ N}
$$

Taking origin as shown in the figure, the centre of gravity of solid block is (0.375, 0.3, 0.25) and that of hollow portion is (0.5, 0.4, 0.15). The following table may be prepared now:

Table

 $\Sigma W_i = 7889.38$ $\Sigma W_i x_i = 3241.57$ $\Sigma W_i y_i = 2593.25$ $\Sigma W_i z_i = 1745.91$

$$
\overline{x} = \frac{\sum W_i x_i}{W} = \frac{\sum W_i x_i}{\sum W_i} \qquad \overline{x} = \frac{3241.57}{7889.38} = 0.411 \text{ m}
$$

Similarly,

$$
\overline{y} = \frac{2593.25}{7887.38} = 0.329 \text{ m}
$$

$$
\overline{z} = \frac{1745.91}{7889.38} = 0.221 \text{ m}
$$

IMPORTANT FORMULAE

- **1.** Area of sector of a circle = $R^2 \alpha$
- **2.** Area of parabolic spandrel

\n- 2. Area of parabolic spander!
\n- (*i*) if
$$
y = kx^2
$$
, $A = \frac{1}{3}ha = \frac{1}{3} \times$ the area of rectangle of size $a \times h$
\n- (*ii*) if $y^2 = kx$, $A = \frac{2}{3}ha = \frac{2}{3} \times$ the area of rectangle of size $a \times h$
\n- 3. Surface area of the cone $=\pi Rl$
\n- 4. Surface area of the sphere $= 4\pi R^2$
\n- 5. Volume of a cone $= \frac{\pi R^2 h}{3}$
\n- 6. Volume of a sphere $= \frac{4}{3}\pi R^3$
\n- 7. Centroid of a arc of a circle is at $x_c = \frac{R \sin \alpha}{\alpha}$ from the centre of circle on the symmetry equation.
\n- 8. Centroid of a composite figure is given by $x_c = \frac{\sum A_i x_i}{A}$, $y_c = \frac{\sum A_i y_i}{A}$
\n- 9. Centroid of simple figure from the reference axis $\overline{y} = \frac{y \, dA}{A}$
\n- 10. For centroid of standard figures refer Table 4.2.
\n- 11. $I_{yy} = \sum x_i^2 dA_i$ and $I_{xx} = \sum y_i^2 dA_i$, $I_{zz} = \sum r_i^2 dA_i = \int r^2 dA$
\n- 12. Radius of gyration $k = \sqrt{\frac{I}{A}}$ i.e. $I = Ak^2$.
\n- 13. $I_{zz} = I_{xx} + I_{yy}$.
\n- 14. $I_{AB} = I_{GG} + A y_c^2$.
\n- 15. Moment of inertia of standard sections are as shown in Table 4.4.
\n- 16. Pappus-Guldinus Theorems: (i) The area of surface generated by revolving a plane curve about a non-intersect of the plane is equal to the length of the generating curve times the traveling of the plane area during the rotation.
\n- (*ii*) The volume of the solid generated by a plane area about a non-intersecting plane.
\n

 πR^2 3 R^2h

- **3.** Surface area of the cone = πRl
- **4.** Surface area of the sphere = $4\pi R^2$
- **5.** Volume of a cone
- **6.** Volume of a sphere $=$ $\frac{4}{3} \pi R^3$
- **7.** Centroid of a arc of a circle is at $x_c = \frac{R \sin \alpha}{\alpha}$ from the centre of circle on the symmetric axis.
- **8.** Centroid of a composite figure is given by

$$
x_c = \frac{\sum A_i x_i}{A}, \quad y_c = \frac{\sum A_i y_i}{A}.
$$

9. Centroid of simple figure from the reference axis $\overline{y} = \frac{\int y \, dA}{A}$.

$$
\overline{y} = \frac{\int y \, dA}{A}.
$$

10. For centroid of standard figures refer Table 4.2.

\n- 10. For centroid of standard figures refer Table 4.2.
\n- 11.
$$
I_{yy} = \sum x_i^2 dA_i
$$
 and $I_{xx} = \sum y_i^2 dA_i$, $I_{zz} = \sum r_i^2 dA_i = \int r^2 dA$.
\n

- **12.** Radius of gyration $k = \sqrt{\frac{I}{A}}$ *i.e.* $I = Ak^2$.
- **13.** $I_{zz} = I_{xx} + I_{yy}$.
- **14.** $I_{AB} = I_{GG} + Ay_c^2$.
- **15.** Moment of inertia of standard sections are as shown in Table 4.4.
- **16.** Pappus-Guldinus Theorems:
	- (*i*) The area of surface generated by revolving a plane curve about a non-intersecting axis in the plane of the curve is equal to the length of the generating curve times the distance travelled by the centroid of the curve in the rotation.
	- (*ii*) The volume of the solid generated by a plane area about a non-intersecting axis in the plane is equal to the area of the generating plane times the distance travelled by the centroid of the plane area during the rotation.
- **17.** In a right circular cone, the centre of gravity lies at a distance $\frac{3}{4}$ $\frac{1}{4}$ × height from the vertex along the axis of rotation.
- **18.** The centre of gravity of a solid hemisphere of radius *r* is at a distance $\frac{3}{6}$ 8 × *r* from its diametral axis.

THEORY QUESTIONS

- **1.** Determine the centroid of an arc of radius *R* from first principle.
- **2.** Distinguish between centroid and centre of gravity.
- **3.** Determine the centroid of a triangle of base width '*b*' and height '*h*' by the method of integration.
- **4.** Locate the centroid of a semicircle from its diametral axis using the method of integration.
- **5.** Explain the terms moment of inertia and radius of gyration of a plane figure.
- **6.** State and prove
	- (*a*) Perpendicular axis theorem
	- (*b*) Parallel axis theorem of moment of inertia.
- **7.** Determine the moment of inertia of the areas specified below by first principle: (*i*) Triangle of base width *b* and height *h* about its centroid axis parallel to the base. (*ii*) A semicircle about its centroidal axis parallel to the diametral axis.
- **8.** State and explain theorems of Pappus-Guldinus.
- **9.** Locate the centre of gravity of the right circular cone of base radius *R* and height *h.*
- **10.** Determine the centre of gravity of a solid hemisphere of radius *R* from its diametral axis.

PROBLEMS FOR EXERCISE

1. Determine the centroid of the built-up section in Fig. 4.74. Express the coordinates of centroid with respect to *x* and *y* axes shown. [Ans. $\bar{x} = 48.91$ mm; $\bar{y} = 61.30$ mm]

2. Determine the centroid of the reinforced concrete retaining wall section shown in Fig. 4.75. [Ans. $\bar{x} = 1.848 \text{ m}; \ \bar{y} = 1.825 \text{ m}$]

3. Find the coordinates of the centroid of the shaded area with respect to the axes shown in Fig. 4.76. [**Ans.** *x* = 43.98 mm; *y* = 70.15 mm]

4. A circular plate of uniform thickness and of diameter 500 mm as shown in Fig. 4.77 has two circular holes of 40 mm diameter each. Where should a 80 mm diameter hole be drilled so that the centre of gravity of the plate will be at the geometric centre.

[Ans. $x = 50$ mm; $y = 37.5$ mm]

5. With respect to the coordinate axes *x* and *y* locate the centriod of the shaded area shown in Fig. 4.78. **[Ans.** $\bar{x} = 97.47$ mm; $\bar{y} = 70.77$ mm]

Fig. 4.77

6. Locate the centroid of the plane area shown in Fig. 4.79.

[Ans. $\bar{x} = 104.10$ mm; $\bar{y} = 44.30$ mm]

7. Determine the coordinates of the centroid of shaded as shown in Fig. 4.80 with respect to the corner point *O*. Take $x = 40$ mm. [Ans. $\bar{x} = 260.07$ mm; $\bar{y} = 113.95$ mm]

- **8.** *ABCD* is a square section of sides 100 mm. Determine the ratio of moment of inertia of the section about centroidal axis parallel to a side to that about diagonal *AC*. [**Ans.** 1]
- **9.** The cross-section of a rectangular hollow beam is as shown in Fig. 4.81. Determine the polar moment of inertia of the section about centroidal axes.

10. The cross-section of a prestressed concrete beam is shown in Fig. 4.82. Calculate the moment of inertia of this section about the centroidal axes parallel to and perpendicular to top edge. Also determine the radii of gyration. [Ans. $I_{xx} = 1.15668 \times 10^{10} \text{ mm}^4$; $k_{xx} = 231.95 \text{ mm}$; I_{yy} = 8.75729 × 10⁹ mm⁴; k_{yy} = 201.82 mm]

11. The strength of a 400 mm deep and 200 mm wide *I*-beam of uniform thickness 10 mm, is increased by welding a 250 mm wide and 20 mm thick plate to its upper flanges as shown in Fig. 4.83. Determine the moment of inertia and the radii of gyration of the composite section with respect to cetroidal axes parallel to and perpendicular to the bottom edge *AB*.

12. The cross-section of a gantry girder is as shown in Fig. 4.84. It is made up of an *I*-section of depth 450 mm, flange width 200 mm and a channel of size 400 mm \times 150 mm. Thickness of all members is 10 mm. Find the moment of inertia of the section about the horizontal centroid axis. $I_{xx} = 4.2198 \times 10^8 \text{ mm}^4$

13. A plate girder is made up of a web plate of size 400 mm \times 10 mm, four angles of size 100 mm \times 100 mm \times 10 mm and cover plates of size 300 mm \times 10 mm as shown in Fig. 4.85. Determine the moment of inertia about horizontal and vertical centroidal axes.

[Ans. $I_{xx} = 5.35786 \times 10^8 \text{ mm}^4$; $I_{yy} = 6.0850667 \text{ mm}^4$]

14. Determine the moment of inertia and radii of gyration of the area shown in Fig. 4.86 about the base *A-B* and the centroidal axis parallel to *AB*.

15. Determine the moment of inertia of the section shown in Fig. 4.87 about the vertical centroidal axis. **[Ans.** $I_{yy} = 5,03,82,857 \text{ mm}^4$]

16. A semi-circular cut is made in rectangular wooden beam as shown in Fig. 4.88. Determine the polar moment of inertia of the section about the centroidal axis.

17. Determine the moment of inertia of the section shown in the Fig. 4.89 about the horizontal centroidal axis. Also find the moment of inertia of the section about the symmetrical axis. Hence find the polar moment of inertia.

[Ans. $I_{xx} = 54,09,046$ mm⁴; $I_{yy} = 14,55,310$ mm⁴; $I_{zz} = 68,64,356$ mm⁴] **Semicircle PA** $30 - 30$ 100 **Fig. 4.89** 17. Determine the anoment of inertia of the steel and shown in the Fig. 4.89 about the centroidal axis. Also find the moment of inertia of the section about the symmetric Hence find the polar moment of inertia of the sect

18. The cross-section of a machine part is as shown in Fig. 4.90. Determine its moment of inertia and radius of gyration about the horizontal centroidal axis.

[Ans. $I_{xx} = 5249090.85$ mm⁴; $k_{xx} = 27.05$ mm]

19. The cross-section of a plain concrete culvert is as shown in Fig. 4.91. Determine the moment of inertia about the horizontal centroidal axes. [Ans. $I_{xx} = 5.45865 \times 10^{10} \text{ mm}^4$]

20. Determine the centroid of the built-up section shown in Fig. 4.92 and find the moment of inertia and radius of gyration about the horizontal centroidal axis.

[Ans. $I_{xx} = 1267942$ mm⁴; $k_{xx} = 18.55$ mm]

21. Determine the centre of gravity of the pyramid shown in Fig. 4.93.

$$
\left[\mathbf{Ans.}\ x=\frac{3}{4}\ h\right]
$$

- **22.** A steel ball of diameter 150 mm rests centrally over a concrete cube of size 150 mm. Determine the centre of gravity of the system, taking weight of concrete = 25000 N/m³ and that of steel 80000 N/m³. [Ans. 168.94 mm from base]
- **23.** Locate the centre of gravity of the wire shown in Fig. 4.94. Portion *BC* is in *x-y* plane and semicircle *CD* is parallel to *x*-*z* plane.

5

Friction

When a body moves or tends to move over another body, a force opposing the motion develops at the contact surfaces. The force which opposes the movement or the tendency of movement is called the **frictional force** or simply **friction**. So far, in earlier chapters, we had ignored this force and considered contacting surfaces are smooth. Actually in almost all cases the contacting surfaces are not smooth. There are minutely projecting particles which develop frictional force to oppose the tendency to movement of one surface over the other surface. In this chapter, the additional terminology used in connection with frictional forces are explained and laws of dry friction (wet friction excluded) are presented. Applications of these laws to many engineering problems are illustrated.

5.1 COEFFICIENT OF FRICTION

Whenever a resultant force acts in the direction of contacting surfaces frictional force develops to oppose that force. The frictional force, like any other reaction, has a remarkable property of adjusting itself in magnitude to the tangential force. However, there is a limit beyond which the magnitude of the frictional force will not develop. If the applied tangential force is more than this maximum frictional force, there will be movement of one body over the other body with an acceleration as per Newton's second law of mass times acceleration equal to the resultant force. This maximum value of frictional force, which comes into play when the motion is impending is known as **Limiting Friction**. It may be noted that when the applied tangential force is less than the limiting friction, the body remains at rest and such friction is called **Static Friction**, which will have any value between zero and limiting friction. If the value of applied tangential force exceeds the limiting friction, the body starts moving over another body and the frictional resistance experienced while moving is known as **Dynamic Friction**. The magnitude of dynamic friction is found to be less than limiting friction. Dynamic friction may be further classified into two groups: **EPIC**

When a body moves or tends to move over another body, a force opposing the motion devented

contact surfaces. The force which opposes the movement of the funderical

references are simply friction. So far, in earl

(*i*) *Sliding Friction:* It is the friction experienced by a body when it slides over the other body.

(*ii*) *Rolling Friction:* It is the friction experienced by a body when it rolls over another body.

It has been experimentally proved that, between two contacting surfaces, the magnitude of limiting friction bears a constant ratio to the normal reaction between the two and this ratio is called '*Coefficient*

of Friction'. Referring to Fig. 5.1 a body weighing *W* is being pulled by a force *P* and the motion is impending. Let *N* be normal reaction and *F* the limiting frictional force. Then

Coefficient of Friction =
$$
\frac{F}{N}
$$

Coefficient of friction is denoted by µ. Then

$$
\mu = \frac{F}{N}
$$

5.2 LAWS OF FRICTION

The principles discussed in a previous article are mainly due to the experiemental studies by Coulomb (1781) and by Mozin (1831). These principles constitute the laws of dry friction and may be called as *Coulomb's laws of dry friction*. These laws are listed below:

- (*i*) The frictional force always acts in a direction opposite to that in which the body tends to move.
- (*ii*) Till the limiting value is reached, the magnitude of frictional force is exactly equal to the tangential force which tends to move the body.
- (*iii*) The magnitude of the limiting friction bears a constant ratio to the normal reaction between the two contacting surfaces.
- (*iv*) The force of friction depends upon the roughness/smoothness of the surfaces.
- (*v*) The force of friction is independent of the area of contact between the two surfaces.
- (*vi*) After the body starts moving, the dynamic friction comes into play, the magnitude of which is less than that of limiting friction and it bears a constant ratio to the normal force. This ratio is called coefficient of dynamic friction.

5.3 ANGLE OF FRICTION, ANGLE OF REPOSE AND CONE OF FRICTION

Angle of Friction

Consider the block shown in Fig. 5.2 subject to pull *P*. Let *F* be the frictional force developed and *N* the normal reaction. Thus, at contact surface, the reactions are *F* and *N*. They can be combined to get the resultant reaction *R* which acts at angle θ to normal reaction. This angle is given by Coefficient of Friction is denoted by µ. Then
 $\mu = \frac{k}{N}$

Coefficient of friction is denoted by µ. Then
 $\mu = \frac{F}{N}$
 EIG. 5.1
 Co. LAWS OF PRICTON

The principles discussed in a previous article are mainly due to

$$
\tan \theta = \frac{F}{N}
$$

Fig. 5.2

As frictional force increases the angle θ increases and it can reach maximum value α when limiting value of friction is reached. Thus, when motion is impending

FRICTION **163**

$$
\tan \alpha = \frac{F}{N} = \mu \tag{5.2}
$$

and this value of α is called *angle of limiting friction*. Hence, the angle of limiting friction can be defined as the angle between the resultant reaction and the normal to the plane on which the motion of the body is impending.

Angle of Repose

It is very well-known that when grains (food grains, sand, cement, soil etc.) are heaped, there exists a limit for the inclination of the heap. Beyond that the grains start rolling down. The limiting angle up to which the grains repose (sleep) is called *angle of repose*.

Now consider the block of weight *W* shown in Fig. 5.3 which is resting on an inclined plane that makes angle θ with the horizontal. When θ is a small, block rests on the plane. If θ is increased gradually a stage is reached at which the block starts sliding. The angle made by the plane with the horizontal is called angle of friction for the contacting surfaces. Thus, the maximum inclination of the plane on which the body, free from external forces, can repose is called *angle of repose*. definited as the angle between the resultant reaction and the normal to the plane on which the

the body is impending.
 Angle of Repose
 Li is very well-haven that when grains (food grains, sand, cennent, soil (cc) ar

Consider the equilibrium of the block shown in Fig. 5.3. Since the surface of contact is not smooth, not only normal reaction but frictional force also develops. As the body tends to slide down, the frictional resistance will be up the plane.

Σ Forces normal to plane = 0, gives

$$
N = W \cos \theta \qquad ...(i)
$$

Σ Forces parallel to plane = 0, gives

 $F = W \sin \theta$...(*ii*)

Dividing eqn. (*ii*) by eqn. (*i*) we get,

$$
\frac{F}{N} = \tan \theta
$$

Fig. 5.3

If ϕ is the value of θ when motion is impending, frictional force will be limiting friction and hence

$$
\tan \phi = \frac{F}{N}
$$
...(5.3)

$$
\tan \phi = \mu = \tan \alpha
$$

or

$$
\phi = \alpha.
$$

Thus, the value of angle of repose is the same as the value of limiting angle of friction.

Cone of Friction

When a body is having impending motion in the direction of *P*, the frictional force will be the limiting friction and the resultant reaction *R* will make limiting frictional angle α with the normal as shown in

Fig. 5.4. If the body is having impending motion in some other direction, the resultant reaction makes limiting frictional angle α with the normal. Thus, if the direction of force P is gradually changed through 360°, the resultant *R* generates a right circular cone with semicentral angle equal to α .

If the resultant reaction lies on the surface of this inverted right circular cone whose semicentral angle is limiting frictional angle α , the motion of the body is impending. If the resultant is within this cone the body is stationary. This inverted cone with semicentral angle, equal to limiting frictional angle α , is called *cone of friction*.

5.4 PROBLEMS ON BLOCKS RESTING ON HORIZONTAL AND INCLINED PLANES

Analysis of such problems, when motion is impending is illustrated in this article by solving typical problems.

Example 5.1. Block A weighing 1000 N rests over block B which weighs 2000 N as shown in Fig. 5.5(a). Block A is tied to a wall with a horizontal string. If the coefficient of friction between A and B is 1/4 and that between B and the floor is 1/3, what value of force P is required to create impending motion if (a) P is horizontal, (b) P acts 30° upwards to horizontal?

Solution: **(a) When** *P* **is horizontal:** The free body diagrams of the two blocks are shown in Fig. 5.5(*b*). Note the frictional forces are to be marked in the opposite directions of impending relative motion. In this problem, block *B* is having impending motion to the right. Hence on it F_1 and F_2 are towards right. The relative motion of block *A* w.r.t. *B* is to the left. Hence the direction of F_1 in this is towards the right. Another way of thinking for the direction of F_1 in case of block *A* can be 'actions and reactions are equal and opposite'. Hence on block *B* if F_1 is towards left, on *A* it should be towards right.

FRICTION **165**

Now consider the equilibrium of block *A*.

$$
\sum F_V = 0 \rightarrow
$$

$$
N_1 - 1000 = 0
$$
 or $N_1 = 1000$ newton.

Since F_1 is limiting friction,

$$
\frac{F_1}{N_1} = \mu_1 = \frac{1}{4}
$$

$$
\therefore \qquad F_1 = \frac{1}{4} \times 1000
$$

4

$$
\sum F_H = 0 \rightarrow
$$

F₁ - T = 0 or T = F₁, *i.e.* T = 250 newton.

 $= 250$ newton.

Consider the equilibrium of block *B*.

$$
\Sigma F_V = 0 \rightarrow
$$

N₂ - N₁ - 2000 = 0.
N₂ = N₁ + 2000 = 1000 + 2000 = 3000 newton.

Since F_2 is limiting friction,

Since
$$
F_1
$$
 is limiting friction,
\n
$$
\frac{F_1}{N_1} = \mu_1 = \frac{1}{4}
$$
\n
$$
\therefore \qquad F_1 = \frac{1}{4} \times 1000 = 250 \text{ newton.}
$$
\n
$$
\sum F_H = 0 \rightarrow
$$
\n
$$
F_1 - T = 0 \text{ or } T = F_1, \text{ i.e., } T = 250 \text{ newton.}
$$
\nConsider the equilibrium of block *B*.
\n
$$
\sum F_V = 0 \rightarrow
$$
\n
$$
N_2 - N_1 - 2000 = 0.
$$
\n
$$
N_2 = N_1 + 2000 = 1000 + 2000 = 3000 \text{ newton.}
$$
\nSince F_2 is limiting friction,
\n
$$
F_2 = \mu_2 N_2 = \frac{1}{3} \times 3000 = 1000 \text{ newton.}
$$
\n
$$
\sum F_H = 0 \rightarrow
$$
\n
$$
P - F_1 - F_2 = 0
$$
\n
$$
P = F_1 + F_2 = 250 + 1000 = 1250 \text{ newton.}
$$
\n(b) When *P* is inclined: Free body diagrams for this case are shown in Fig. 5.5(c).
\nConsidering equilibrium of block *A*, we get
\n
$$
\sum F_V = 0 \rightarrow N_1 = 1000 \text{ newton.}
$$
\n
$$
\therefore \qquad F_1 = \frac{1}{4} \times 1000 = 250 \text{ newton.}
$$
\nConsider the equilibrium of block *B*.
\n
$$
\sum F_V = 0 \rightarrow T = F_1 = 250 \text{ newton.}
$$

(b) When *P* **is inclined:** Free body diagrams for this case are shown in Fig. 5.5(*c*). Considering equilibrium of block *A*, we get

$$
\sum F_V = 0 \rightarrow N_1 = 1000
$$
 newton.
\n $\therefore F_1 = \frac{1}{4} \times 1000 = 250$ newton.

$$
\sum F_H = 0 \rightarrow T = F_1 = 250
$$
 newton.

Consider the equilibrium of block *B*.

$$
\sum F_V = 0 \quad \rightarrow
$$

Solution: 900 N block is on the verge of moving downward. Hence frictional forces F_1 and F_2 [Ref. Fig. 5.6(b)] act up the plane on 900 N block. Free body diagrams of the blocks are as shown in Fig. 5.6(*b*).

Consider the equilibrium of 300 N block.

 Σ Forces normal to plane = 0 \rightarrow

$$
N_1 - 300 \cos \theta = 0
$$
 or $N_1 = 300 \cos \theta$...(i)

Example 5.3. A block weighing 500 N just starts moving down a rough inclined plane when it is subjected to 200 N force acting up the inclined plane and it is at the point of moving up the plane when pulled up by a force of 300 N parallel to the plane. Find the inclination of the plane and the coefficient of friction between the inclined plane and the block.

Solution: Free body diagram of the block when its motion is impending down the plane is shown in Fig. $5.7(a)$ and that when it is moving up the plane is shown in Fig. $5.7(b)$.

When block starts moving down the plane [Ref. Fig. 5.7(a)]

Example 5.4. Block A weighing 1000 N and block B weighing 500 N are connected by flexible wire. The coefficient of friction between block A and the plane is 0.5 while that for block B and the plane is 0.2. Determine what value of inclination of the plane the system will have impending motion down the plane? [Ref. Fig. 5.8].

 $\mu = \frac{50}{500 \cos 30} = 0.115$

500 μ cos 30 = 300 - 250 = 50

Solution: Let θ be the inclination of the plane for which motion is impending. Free body diagrams of blocks *A* and *B* are as shown in Fig. 5.8(*b*). Considering equilibrium of block *A*,

Σ Forces normal to plane = $0 \rightarrow$

$$
N_1 - 1000 \cos \theta = 0
$$
 or $N_1 = 1000 \cos \theta$...(i)

∴ From law of friction

$$
F_1 = \mu_1 N_1 = 0.5 \times 1000 \cos \theta = 500 \cos \theta \quad ...(ii)
$$

 Σ Forces parallel to plane = 0 \rightarrow

or
\n
$$
F_1 - T - 1000 \sin \theta = 0
$$

\n $T = 500 \cos \theta - 1000 \sin \theta$...(iii)

Consider the equilibrium of block *B*,

 Σ Forces normal to plane = 0 \rightarrow

$$
N_2 - 500 \cos \theta = 0
$$
 or $N_2 = 500 \cos \theta$...(iv)

From law of friction,

$$
F_2 = \mu_2 N_2 = 0.2 \times 500 \cos \theta = 100 \cos \theta \quad ...(v)
$$

Σ Forces parallel to plane = $0 \rightarrow$

$$
F_2 + T - 500 \sin \theta = 0
$$

Using the values of F_2 and *T* from eqn. (*v*) and eqn. (*iii*),

$$
100 \cos \theta + 500 \cos \theta - 1000 \sin \theta - 500 \sin \theta = 0
$$

$$
600 \cos \theta = 1500 \sin \theta
$$

$$
\tan \theta = \frac{600}{1500}
$$

$$
\theta = 21.8^{\circ}
$$

Example 5.5. What is the value of P in the system shown in Fig. 5.9(a) to cause the motion to impend? Assume the pulley is smooth and coefficient of friction between the other contact surfaces is 0.2.

Solution: Free body diagrams of the blocks are as shown in Fig. 5.9(*b*). Consider the equilibrium of 750 N block.

 Σ Forces normal to the plane = 0 \rightarrow

$$
N_1 - 750 \cos 60 = 0 \quad \therefore \quad N_1 = 375 \text{ newton} \tag{1}
$$

Since the motion is impending, from law of friction,

$$
F_1 = \mu N_1 = 0.2 \times 375 = 75
$$
 newton ...(ii)

Σ Forces parallel to the plane = $0 \rightarrow$

 $T - F_1 - 750 \sin 60 = 0$

$$
T = 75 + 750 \sin 60 = 724.5 \text{ newton.} \qquad ...(iii)
$$

Consider the equilibrium of 500 N block.

$$
\sum F_V = 0 \rightarrow
$$

i.e.,
\n
$$
N_2 - 500 + P \sin 30 = 0
$$

\n $N_2 + 0.5P = 500$...(iv)

$$
F_2 = \mu N_2 = 0.2 (500 - 0.5P) = 100 - 0.1P
$$
 ...(v)

From law of friction,

 $\sum F_H = 0 \rightarrow$ *P* cos $30 - T - F_2 = 0$ *i.e.*, $P \cos 30 - 724.5 - 100 + 0.1 P = 0$

∴ $P = 853.5 N$

Example 5.6. Two identical planes AC and BC, inclined at 60° and 30° to the horizontal meet at C as shown in Fig. 5.10. A load of 1000 N rests on the inclined plane BC and is tied by a rope passing over a pulley to a block weighing W newtons and resting on the plane AC. If the coefficient of friction between the load and the plane BC is 0.28 and that between the block and the plane AC is 0.20, find the least and greatest values of W for the equilibrium of the system.

*Solution***:**

(a) Least value of *W***:**

In this case motion of 1000 N block is impending down the plane and block *W* has impending motion up the plane. Hence free body diagrams for the blocks are as shown in Fig. 5.10(*b*). Considering the equilibrium of 1000 N block,

Σ Forces normal to the plane = $0 \rightarrow$

$$
N_1 - 1000 \cos 30 = 0 \quad \therefore N_1 = 866.0 \text{ newton} \tag{i}
$$

From the law of friction

$$
F_1 = \mu_1 N_1 = 0.28 \times 866.0 = 242.5 \text{ newton} \qquad \qquad \dots (ii)
$$

Σ Forces parallel to the plane = $0 \rightarrow$

T – 1000 sin 30 + F_1 = 0

∴ $T = 500 - 242.5 = 257.5$ newton ...(*iii*) Now consider the equilibrium of block weighing *W*. Σ Forces normal to the plane = $0 \rightarrow$ $N_2 - W \cos 60 = 0$ *i.e.*, $N_2 = 0.5$ *W* ...(*iv*) From law of friction $F_2 = \mu_2 N_2 = 0.2 \times 0.5 \text{ W} = 0.1 \text{ W}$...(*v*) Σ Forces parallel to the plane = 0 → $T - F_2 - W \sin 60 = 0$ Substituting the values of *T* and F_2 from eqns. (*iii*) and (*v*), we get $257.5 - 0.1$ *W* – *W* sin 60 = 0 ∴ $W = \frac{257.5}{0.1}$ $0.1 + \sin 60$ **= 266.6 N. (b) For the greatest value of W:** In such case 1000 N block is on the verge of moving up the plane and *W* is on the verge of moving down the plane. For this case free body diagrams of the blocks are as shown in Fig. 5.10(*c*). Considering the block of 1000 N, Σ Forces normal to plane = $0 \rightarrow$ N_1 – 1000 cos 30 = 0 ∴ N_1 = 866.0 newton ...(*vi*) From law of friction, $F_1 = \mu_1 N_1 = 0.28 \times 866.0 = 242.5 \text{ N}$...(*vii*) Σ Forces parallel to the plane = 0 → *T* – 1000 sin 30 – $F_1 = 0$ 2 ενεκε sommation was protoned to $N_2 = W_2E_0$ ($k_1 = 0.5$ W

From law of friction
 $\sum_i = \mu_i N_i = 0.2 \times 0.5$ W = 0.1 W

Σ Forces parallel to the plane = 0 →

Substituting the values of T and F₂ From eqns. (*ii)* and (*v*

∴ $T = 500 + 242.5 = 742.5$ newton ...(*viii*)

Considering the equilibrium of block weighing *W*,

$$
\Sigma
$$
 Forces normal to plane = 0 \rightarrow

$$
N_2 - W \cos 60 = 0 \quad \text{or} \quad N_2 = 0.5 \ W \tag{ix}
$$

$$
F_2 = \mu_2 N_2 = 0.2 \times 0.5 W = 0.1 W \tag{1}
$$

Σ Forces parallel to plane = $0 \rightarrow$

$$
T - W \sin 60 + F_2 = 0 \tag{xi}
$$

Substituting the values of *T* and F_2 from eqns. (*viii*) and (*x*), we get,

$$
742.5 - W \sin 60 + 0.1 W = 0
$$

or
$$
W = \frac{742.5}{\sin 60 - 0.1} = 969.3
$$
 newton

FRICTION **173**

The system of blocks are, in equilibrium for $W = 266.6$ N to 969.3 N.

Example 5.7. Two blocks connected by a horizontal link AB are supported on two rough planes as shown in Fig. 5.11(a). The coefficient of friction on the horizontal plane is 0.4. The limiting angle of friction for block B on the inclined plane is 20°. What is the smallest weight W of the block A for which equilibrium of the system can exist if weight of block B is 5 kN?

 $\sum F_{H} = 0 \rightarrow$

 $F_2 - C = 0$ or $F_2 = C = 4.196$ kN

From law of friction,

i.e., $4.196 = 0.4 \text{ N}_2$

 $F_2 = \mu N_2$

5.5 APPLICATION TO WEDGE PROBLEMS

Wedges are small pieces of hard materials with two of their opposite surfaces not parallel to each other. They are used to slightly lift heavy blocks, machinery, precast beams etc. for making final alignment or to make place for inserting lifting devices. In any problem weight of wedge is very small compared to the weight lifted. Hence in all problems self weight of wedge is neglected. It is found that in the analysis instead of treating normal reaction and frictional force independently, it is advantageous to treat their resultant.

If *F* is limiting friction, then resultant *R* makes limiting angle α with the normal. Its direction should be marked correctly. Note that the tangential component of the resultant reaction *R* is the frictional force and it will always oppose impending motion. Application to wedge problems is illustrated below by solving problems.

Example 5.8. Determine the force P required to start the movement of the wedge as shown in Fig. 5.12(a). The angle of friction for all surfaces of contact is 15°.

Solution: As wedge is driven, it moves towards left and the block moves upwards. When motion is impending limiting friction develops. Hence resultant force makes limiting angle of 15° with normal. The care is taken to mark 15° inclination such that the tangential component of the resultant opposes the impending motion.

The free body diagrams of the block and wedge are shown in Fig. 5.12(*b*). The forces on block and wedge are redrawn in Figs. 5.12(*c*) and (*d*) so that Lami's theorem can be applied conveniently. Applying Lami's theorem to the system of forces on block

Solution: As wedge is driven, it moves towards left and the block moves upwards. When
impending limiting friction develops. Hence resultant force makes limiting angle of 15° w
The care is taken to mark 15° inclination such that the tangential component of the results
the impending motion.
The free body diagrams of the block and wedge are shown in Fig. 5.12(*b*). The forces or
wedge are redrawm in Figs. 5.12(*c*) and (*d*) so that Lami's theorem can be applied co
Applying Lami's theorem to the system of forces on block

$$
\frac{R_1}{\sin(180-15-20)} = \frac{R_2}{\sin(90-15)} = \frac{20}{\sin(15+20+90+15)}
$$

i.e.,

$$
\frac{R_1}{\sin(180-15-20)} = \frac{R_2}{\sin(90-15)} = \frac{20}{\sin(15+20+90+15)}
$$

i.e.,

$$
\frac{R_1}{\sin 145} = \frac{R_2}{\sin 75} = \frac{20}{\sin 140}
$$

 \therefore
 $R_1 = 17.85$ kN
and
 $R_2 = 30.05$ kN
Applying Lami's theorem to system of forces on the wedge, we get

$$
\frac{P}{\sin 130} = \frac{R_2}{\sin 105}
$$

 \therefore
 $P = 23.84$ kN
Example 5.9. A block weighing 160 kN is to be raised by means of the wedges A and B a
Fig. 5.13(a). Find the value of force P for impending motion of block C upwards, if co
friction is 0.25 for all contact surfaces. The self weight of wedges may be neglected.
Solution: Let ϕ be the angle of limiting friction.
 \therefore
 $\phi = \tan^{-1}(0.25) = 14.036^{\circ}$
The free body diagrams of wedges A, B and block are equal. The system of forces on b

Applying Lami's theorem to system of forces on the wedge, we get

$$
\frac{P}{\sin 130} = \frac{R_2}{\sin 105}
$$

∴ *P* **= 23.84 kN**

Example 5.9. A block weighing 160 kN is to be raised by means of the wedges A and B as shown in Fig. 5.13(a). Find the value of force P for impending motion of block C upwards, if coefficient of friction is 0.25 for all contact surfaces. The self weight of wedges may be neglected. Solution: Let φ be the angle of limiting friction.

$$
\phi = \tan^{-1} (0.25) = 14.036^{\circ}
$$

The free body diagrams of wedges *A*, *B* and block *C* are shown in Fig. 5.13(*b*). The problem being symmetric, the reactions R_1 and R_2 on wedges A and B are equal. The system of forces on block C and

on wedge *A* are shown in the form convenient for applying Lami's theorem [Ref. Figs. 5.13(*c*) and (*d*)].

FRICTION **177**

5.6 APPLICATION TO LADDER PROBLEMS

A ladder resting against a wall is a typical case of friction problems in non-concurrent system of forces. Hence we have three equations of equilibrium available. From law of friction we have the

equation $\frac{F}{N} = \mu$. Using equilibrium equations and friction law the problems can be solved. The

procedure is illustrated with the examples below:

Example 5.10. A ladder of length 4 m, weighing 200 N is placed against a vertical wall as shown in Fig. 5.14(a). The coefficient of friction between the wall and the ladder is 0.2 and that between floor and the ladder is 0.3. The ladder, in addition to its own weight, has to support a man weighing 600 N at a distance of 3 m from A. Calculate the minimum horizontal force to be applied at A to prevent slipping.

Solution: The free body diagram of the ladder is as shown in Fig. 5.14(*b*).

$$
\sum M_A = 0 \rightarrow
$$

 $200 \times 2 \cos 60 + 600 \times 3 \cos 60 - F_B \times 4 \cos 60 - N_B \times 4 \sin 60 = 0$

Dividing throughout by 4 and rearranging the terms, we get

$$
0.866 N_B + 0.5 F_B = 275
$$

From law of friction, $F_B = \mu N_B = 0.2 N_B$

$$
\therefore \qquad 0.866 N_B + 0.5 \times 0.2 N_B = 275
$$

or
$$
N_B = 284.7
$$
 newton.

 $F_B = 56.94$ newton.

 $\sum F_V = 0 \rightarrow$ $N_A - 200 - 600 + F_B = 0$ $N_A = 743.06$ newton, since $F_B = 56.94$ ∴ $F_A = \mu_A N_A$ $= 0.3 \times 743.06 = 222.9$ newton $\sum F_H = 0 \rightarrow$ $P + F_A - N_B = 0$ ∴ $P = N_B - F_A = 284.7 - 222.9$ *i.e*., *P* **= 61.8 newton**

Example 5.11. The ladder shown in Fig. 5.15(a) is 6 m long and is supported by a horizontal floor and a vertical wall. The coefficient of friction between the floor and the ladder is 0.25 and between the wall and the ladder is 0.4. The weight of the ladder is 200 N and may be considered as a concentrated load at G. The ladder supports a vertical load of 900 N at C which is at a distance of 1 m from B. Determine the least value of α at which the ladder may be placed without slipping. Determine the *reaction at that stage.*

Solution: Figure 5.15(*b*) shows the free body diagram of the ladder. From law of friction,

coefficient of friction between the ladder and the ground and that between ground and the wall is 0.5 each, what will be the maximum distance on ladder to which a man whose weight is 1.5 times the weight of ladder may ascend before the ladder begins to slip?

Solution: Figure 5.16(*a*) shows the ladder when it is about to slip when the man weighing 1.5 *W* is at a distance, '*a*L' from the end *A*. Its free body diagram is shown in Fig. 5.16(*b*).

Since ladder is on the verge of slipping, from law of friction,

$$
F_A = \mu N_B = 0.5 \ N_B \tag{1}
$$

and
$$
F_B = \mu N_B = 0.5 \ N_B \tag{2}
$$

$$
\Sigma H = 0 \rightarrow
$$

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Thus in this case the man can ascend up to '0.5 *L***' of ladder.**

5.7 BELT FRICTION

The transmission of power by means of belt or rope drives is possible because of friction which exists between the wheels and the belt. Similarly, band brakes stops the rotating discs because of friction between the belt and the disc. All along the contact surface the frictional resistance develops. Hence, the tension in the rope is more on the side it is pulled and is less on the other side. Accordingly, the two sides of the rope may be called as tight side and slack side.

FRICTION **181**

Relationship between Tight Side and Slack Side Forces in a Rope

Figure 5.17 (*a*) shows a load *W* being pulled by a force *P* over a fixed drum. Let the force on slack side be T_1 and on tight side by T_2 [Fig. 5.17(*b*)]. T_2 is more than T_1 because frictional force develops between drum and the rope [Fig. $5.17(c)$]. Let θ be the angle of contact between rope and the drum. Now, consider an elemental length of rope as shown in Fig. 5.17(*d*). Let *T* be the force on slack side and $T + dT$ on tight side. There will be normal reaction *N* on the rope in the radial direction and frictional force $F = \mu N$ in the tangential direction. Then,

exerting a 600 N force at the other end of the rope?

Solution: Angle of contact = $1.25 \times 2\pi = 2.5\pi$ (1) Let the impending motion of the weight be downward. Then,

$$
T_1 = 600 \text{ N}; T_2 = W
$$

$$
\frac{W}{600} = e^{\mu 2.5\pi} = e^{0.3 \times 2.5\pi} = e^{0.75\pi}
$$

$$
W = 6330.43 \text{ N}
$$

(2) Let the impending motion of weight be upwards. Then

$$
T_1 = W;
$$
 $T_2 = 600$ N
\n $T_2 = T_1 e^{\mu\theta}$
\n $600 = W e^{0.75\pi}$
\n $W = 56.87$ N

Thus, a 600 N force can support a range of loads between 56.87 N to 6330.43 N weight on the other side of drum.

Example 5.14. In Fig. 5.19 (a) The coefficient of friction is 0.20 between the rope and the fixed drum, and between other surface of contact µ *= 0.3. Determine the minimum weight W to prevent downward motion of the 1000 N body.*

Fig. 5.18

W

600 N

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Substituting the values from (2), (4) and (5), $0.24 W + 0.24 W + 240 - 1000 \times 0.6 + 2.156 W = 0$

W **= 136.57 N.**

Example 5.15. A torque of 300 N-m acts on the brake drum shown in Fig. 5.20(a). If the brake band is in contact with the brake drum through 250° and the coefficient of friction is 0.3 determine the force P applied at the end of the brake lever for the position shown in the figure.

Solution: Figure 5.20 (*b*) shows free body diagrams of brake drum and the lever arm. Now $T_2 = T_1 e^{\mu \theta}$

^θ ⁼ ²⁵⁰ 180 π radians and µ = 0.3 ∴ µ θ = 1.309 [∴] *T*² = *T*¹ *e*1.309 = 3.7025 *T*¹ Now, (*T*2 – *T*1) *r* = *M* (3.7025 – 1) *T*1 × 250 = 300 × 10³ ∴ *T*¹ = 444.04 N ∴ *T*² = 1644.058 N Civil[datas.blogspo](http://Civildatas.blogspot.in)t.in

Consider the lever arm. Taking moment about the hinge, we get T_2

$$
\times 50 = P \times 300
$$

$$
P = 274.0 \text{ N.}
$$

Example 5.16. Two parallel shafts 3 m apart are to be connected by a belt running over the pulleys of diamter 500 mm and 100 mm respectively. Determine the length of belt required and the angle of contacts between belt and each pulley if it is crossed. What power can be transferred if the larger pulley rotates at 220 revolutions per minute. Given: Maximum permissible tension in the belt = 1 kN and coefficient of friction between the belt and the pulley is 0.25.

Solution: The crossed belt drive system *t* is shown in Fig. 5.21.

From the geometry of the system,

$$
\sin \alpha = \frac{250 + 50}{3000} = 0.1
$$

\n∴ $\alpha = 5.739^{\circ}$
\n∴ $\theta_2 = \theta_1 = 180 + 2\alpha = 191.478^{\circ}$
\n= 3.342 radians.

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of belt and minimum that can be transmitted.

Solution: The parallel drive system is as shown in Fig. 5.22.

$$
\frac{T_2}{T_1} = e^{\mu \theta_2} = e^{0.25 \times 3.008} = 2.1212
$$

$$
T_1 = \frac{T_2}{2.1212} = \frac{1000}{2.1212} = 471.4 \text{ Newton}
$$

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Velocity of belt =
$$
r \times \frac{2\pi \times 220}{60} = 250 \times \frac{2\pi \times 220}{60}
$$

= 5759.6 mm/sec.

∴ Maximum power that can be transferred without slippage

=
$$
(T_2 - T_1)
$$
 × Velocity
= $(1000 - 471.4)$ × 5759.6 = 3044517 N mm/sec

= **3.044517 kW.**

IMPORTANT FORMULAE

1. $\mu = \frac{\text{Limiting } F}{N} = \tan \alpha$

$$
T_2 = T_1 e^{\mu \theta}
$$

THEORY QUESTIONS

- **1.** State the laws of dry friction.
- **2.** Explain the terms: coefficient of friction, angle of friction and cone of friction.
- **3.** Derive the expression for the relationship between tight side and slack side forces in a belt friction problem.

PROBLEMS FOR EXERCISE

- **1.** A pull of 180 N applied upward at 30° to a rough horizontal plane was required to just move a body resting on the plane while a push of 220 N applied along the same line of action was required to just move the same body downwards. Determine the weight of the body and the coefficient of friction. $[Ans. $W = 990 \text{ N}; \mu = 0.1732]$$
- **2.** The block A shown in Fig. 5.23 weighs 2000 N. The cord attached to if passes over a frictionless pulley and supports a weight equal to 800 N. The value of coefficient friction between A and the horizontal plane is 0.35. Determine the horizontal force $P : (i)$ If the motion is impending towards the left. (*ii*) if the motion is impending towards the right.

[**Ans.** (*i*) 1252.82 N (*ii*) 132.82 N]

Frictionless

- **3.** A 3000 N block is placed on an inclined plane as shown in Fig. 5.24. Find the maximum value of *W* for equilibrium if tipping does not occur. Assume coefficient of friction as 0.2. [**Ans.** 2636.15]
- **4.** Find whether block *A* is moving up or down the plane in Fig. 5.25 for the data given below. Weight of block $A = 300$ N. Weight of block $B = 600$ N. Coefficient of limiting friction between plane *AB* and block *A* is 0.2. Coefficient of limiting friction between plane *BC* and block *B* is 0.25. Assume pulley as smooth.

[**Ans.** The block *A* is stationary since F developed $\lt F_{min}$

- **5.** Two identical blocks *A* and *B* are connected by a rod and they rest against vertical and horizontal planes respectively as shown in Fig. 5.26. If sliding impends when $\theta = 45^{\circ}$, determine the coefficient of friction, assuming it to be same for both floor and wall. **[Ans.** 0.414]
- **6.** Determine the force *P* required to start the wedge as shown in Fig. 5.27. The angle of friction for all surfaces of contact is 15°. [**Ans.** 26.6784 kN]

Fig. 5.27

7. Two blocks *A* and *B* weighing 3 kN and 15 kN, respectively, are held in position against an inclined plane by applying a horizontal force *P* as shown in Fig. 5.28. Find the least value of *P* which will induce motion of the block *A* upwards. Angle of friction for all contact surfaces is 12°.

[**Ans.** 14.025 kN]

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FRICTION **189**

- **8.** In Fig. 5.29, *C* is a stone block weighing 6 kN. It is being raised slightly by means of two wooden wedges *A* and *B* with a force *P* on wedge *B*. The angle between the contacting surfaces of the wedge is 5°. If coefficient of friction is 0.3 for all contacting surfaces, compute the value of *P* required to impend upward motion of the block *C*. Neglect weight of the wedges. [**Ans.** 2.344 kN]
- **9.** Find the horizontal force *P* required to push the block *A* of weight 150 N which carries block *B* of weight 1280 N as shown in Fig. 5.30. Take angle of limiting friction between floor and block *A* as 14° and that between vertical wall and block *B* as 13° and coefficient of limiting friction between the blocks as 0.3. **[Ans.** $P = 1294.2 \text{ N}$]
- **10.** The level of precast beam weighing 20,000 N is to be adjusted by driving a wedge as shown in Fig. 5.31. If coefficient of friction between the wedge and pier is 0.35 and that between beam and the wedge is 0.25, determine the minimum force *P* required on

the wedge to make adjustment of the beam. Angle of the wedge is 15°. (**Hint:** Vertical component of reaction on wedge at contact with beam $= 1/2$ vertical load on beam $= 10,000$ kN)

[**Ans.** 9057.4 N]

11. A ladder 5 m long rests on a horizontal ground and leans against a smooth vertical wall at an angle of 70° with the horizontal. The weight of the ladder is 300 N. The ladder is on the verge of sliding when a man weighing 750 N stands on a rung 1.5 m high. Calculate the coefficient of friction between the ladder and the floor.

[Ans. $\mu = 0.1837$]

- **12.** A 4 m ladder weighing 200 N is placed against a vertical wall as shown in Fig. 5.32 as a man weighing 800 N, reaches a point 2.7 m from *A*, the ladder is about to slip. Assuming that the coefficient of friction between the ladder and the wall is 0.2, determine the coefficient of friction between the ladder and the floor. [**Ans.** 0.3548]
- **13.** A uniform ladder of length '*L*' rests against a vertical wall making an angle of 60° with the horizontal. Coefficient of friction between wall and ladder and ground and ladder are 0.3 and 0.25 respectively. A man weighing 650 N ascends the ladder. How
	- high he will be able to go before the ladder slips? Find the magnitude of weight to be put at the bottom of the ladder so as to make it just sufficient to permit the man to go to the top. Take ladder's weight $= 900$ N.

(Hint. Find *P* as found in Example 11. Then $W = \frac{P}{\mu}$) [Ans. (*i*) To a length of 0.4345 *L i.e.*, to a height of 0.435 *L* sin 60° = 0.367 *L* metres. (*ii*) *W* = 777.68 N]

14. Determine the maximum weight that can be lowered by a person who can exert a 300 N pull on rope if the rope is wrapped $2\frac{1}{2}$ turns round a horizontal spur as shown in Fig. 5.33. Coefficient of friction between spur and the rope is 0.3.

15. Determine the minimum value of *W* required to cause motion of blocks *A* and *B* towards right (Ref. Fig. 5.34). Each block weighs 3000 N and coefficient of friction between blocks and the planes is 0.2. Coefficient of friction between the drum and rope is 0.1. Angle of wrap over the drum is 90°. [**Ans.** 3065.18 N]

16. Block *A* shown in Fig. 5.35 weighs 2000 N. The cord attached to *A* passes over a fixed drum and supports a weight equal to 800 N. The value of coefficient of friction between *A* and the horizontal plane is 0.25 and between the rope and the fixed drum is 0.1. Solve for *P*: (1) if motion is impending towards the left, (2) if the motion is impending towards the right.

[**Ans.** (1) 1230.94 N; (2) 143.0 N]

17. The dimension of a brake drum is as shown in Fig. 5.36. Determine the torque *M* exerted on the drum if the load $P = 50$ N. Assume coefficient of kinetic friction between rope and drum to be 0.15. [**Ans.** 747.685 N-m]

Fig. 5.36

- **18.** A belt drive is required to transmit 12 kW power. The velocity of the belt is 9.6 m/sec. If coefficient of friction is 0.25 and the angle of contact is 150º, determine the maximum tension is the belt. **[Ans.** $T_1 = 2.602 \text{ kN}$]
- **19.** Two parallel shafts are having pulleys of diameters 300 mm and 500 mm. The distance between the shafts is 2.5 m. They are connected by crossed belt drive system. The speed of the larger pulley is 500 rpm and coefficient of friction $\mu = 0.3$, determine the maximum power that can be transmitted, if maximum permissible tension in the belt is 2.4 kN. [Ans. $P = 20.3$ kN]

6 Simple Machines

A simple machine is a device with the help of which heavy loads are lifted by applying small effects in a convenient direction. Pulley used to lift water from a well and screw jacks used to lift motor car are some of the common examples of simple machines. In this chapter some of the terms connected with simple machines are explained first followed by the description of the characteristic features of levers, systems of pulleys, wheel and axle, Weston differential pulley block, inclined plane, simple screw jack, differential screw jack and winch crab. **Simple Machinies**
 Simple Machinates
 Simple Machinates is a device with the help of which heavy loads are fifted by applying an a convenient direction. Pulley used to iff where from a well and strew jacks used to co

6.1 DEFINITIONS

The terms commonly used while dealing with simple machines are defined below:

Load: This is the resistance to be overcome by the machine.

Effort: This is the force required to overcome the resistance to get the work done by the machine.

Mechanical Advantage: This is the ratio of load lifted to effort applied. Thus, if *W* is the load and *P* is the corresponding effort, then

Mechanical Advantage =
$$
\frac{W}{P}
$$
 ...(6.1)

Velocity Ratio: This is the ratio of the distance moved by the effort to the distance moved by the load in the same interval of time. Thus,

$$
Velocity Ratio = \frac{D}{d}
$$
...(6.2)

where, $D -$ distance moved by effort

d – distance moved by the load.

Input: The work done by the effort is known as input to the machine. Since work done by a force is defined as the product of the force and the distance moved in the direction of the force, Input = $P \times D$...(6.3)

If force *P* is in newton and distance *D* is in metre, the unit of input will be N-m. One N-m work is also known as one Joule (J).

Output: It is defined as useful work got out of the machine, *i.e.,* the work done by the load. Thus,

Output =
$$
W \times d
$$
 ...(6.4)

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Efficiency: This is defined as the ratio of output to the input. Thus, if we use notation η for efficiency,

$$
\eta = \frac{\text{Output}}{\text{Input}} = \frac{W \times d}{P \times D} = \frac{W}{P} \times \frac{d}{D}
$$
\n= Mechanical Advantage (*MA*) × $\frac{1}{\text{Velocity Ratio (VR)}}$
\n= $\frac{MA}{VR}$...(6.5)
\ni.e.,
\nBfficiency = $\frac{\text{Mechanical Advanced}}{\text{Velocity Ratio}}$...(6.6)
\n $\text{1.2. } \text{1. }$

Ideal Machine: A machine whose efficiency is 1 (*i.e.,* 100%) is called an ideal machine. In other words, in an ideal machine, the output is equal to the input. From eqn. (6.5), in an ideal machine,

Velocity Ratio = Mechanical Advantage

Ideal Effort: Ideal effort is the effort required to lift the given load by the machine assuming the machine to be ideal.

VR = MA

For ideal machine,

If P_i is the ideal effort, then

If
$$
P_i
$$
 is the ideal effort, then
\n
$$
VR = \frac{W}{P_i}
$$
\n
$$
P_i = \frac{W}{VR}
$$
\n...(6.6)

Ideal Load: Ideal load is the load that can be lifted using the given effort by the machine, assuming it to be ideal.

For the ideal machine,

$$
VR = MA
$$

If *Wi* is the ideal load, then

$$
VR = \frac{W_i}{P}
$$

$$
W_i = VR \times P
$$
...(6.7)

6.2 PRACTICAL MACHINES

In practice, it is difficult to get an ideal machine. Friction exists between all surfaces of contacts of movable parts. Some of the work done by the effort is utilised to overcome frictional resistance. Hence, the useful work done in lifting the load is reduced, resulting in reduction of efficiency.

Let $P =$ actual effort required P_i = ideal effort required *W =* actual load to be lifted W_i = ideal load to be lifted

Then,

 $P - P_i$ is called *effort lost in friction* and $W - W_i$ is called *frictional resistance*. Now, $\eta = \frac{MA}{VR}$ *W* $=\frac{W}{P} \times \frac{1}{VR}$ From eqn. (6.6) , $=\frac{W}{VR}$ ∴ $η = \frac{P_i}{P_i}$ *P i* Similarly from eqn. (6.7), $W_i = VR \times P$ ∴ η *⁼^W Wi* Thus, $\eta = \frac{P_i}{P_i}$ *W i* $...(6.8)$

Example 6.1. In a lifting machine, an effort of 500 N is to be moved by a distance of 20 m to raise a load of 10,000 N by a distance of 0.8 m. Determine the velocity ratio, mechanical advantage and efficiency of the machine. Determine also ideal effort, effort lost in friction, ideal load and frictional resistance.

W

i

P

```
Solution: Load, W = 10,000 N
                     Effort P = 500 N
                 Distance moved by the effort D = 20 m
                 Distance moved by the load d = 0.8 m
                 Mechanical advantage, MA = \frac{W}{P} = \frac{10,000}{500},
                                                             = 20
                 Velocity ratio,\frac{D}{d} = \frac{20}{0.8}= 25
                 Efficiency, \eta = \frac{MA}{VR} = \frac{20}{25}= 0.8 = 80 per cent
                 Ideal effort, P_i = \frac{W}{VR} = \frac{10,000}{25} = 400 \text{ N}Effort lost in friction = P - P_i= 500 - 400= 100 N
                 Ideal load,
                                                        W_i = P \times VR = 500 \times 25= 12,500 N
From eqn. (6.6),<br>
P_1 = \frac{W}{VR} = \frac{W}{VR}<br>
From eqn. (6.6),<br>
P_1 = \frac{W}{VR}<br>
Similarly from eqn. (6.7), W_1 = VR \times P<br>
\therefore \Pi = \frac{W}{W_1}<br>
Thus,<br>
\Pi = \frac{W}{W_1}<br>
Thus,<br>
\Pi = \frac{W}{W_1}<br>
\therefore \Pi = \frac{W}{W_1}<br>
\therefore \Pi = \frac{W}{
```
Frictional resistance

 $= W_i - W$ $= 12,500 - 10,000$ **= 2500 N**

6.3 LAW OF MACHINE

The relationship between the load lifted and the effort required in a machine is called the **law of machine.** This is found by conducting experiments in which efforts required for lifting different loads are determined and then load *versus* effort graph as shown in Fig. 6.1 is plotted. This is generally a straight line which does not pass through the origin.

The law of machine can be expressed mathematically in the form:

$$
P = mW + C \tag{6.9}
$$

where, *C* is the intercept *OA* and $m = \tan \theta$, the slope of *AB*. For the ideal machine

$$
MA = VR
$$

$$
\frac{W}{P} = VR
$$

This is a straight line relationship passing through the origin and is shown by line *OC* in Fig. 6.1.

After plotting the law for the actual machine (AB) and the law for ideal machine (OC) , it is easy to determine efficiency at any given load. The vertical line *DEF* corresponding to given load *OD* is drawn. Then,

$$
DE = P_i
$$
 effort required in ideal machine
\n
$$
DF = P
$$
, effort required in actual machine
\nFriction loss = $P - P_i$
\n
$$
= DF - DE
$$

\n
$$
= EF
$$

\nEfficiency,
\n
$$
\eta = \frac{P_i}{P}
$$

\n
$$
= \frac{DE}{DF}
$$

6.4 VARIATION OF MECHANICAL ADVANTAGE

Mechanical Advantage (*MA*) is given by:

$$
MA = \frac{W}{P}
$$

From the law of machine,

$$
P = mW + C
$$

\n
$$
MA = \frac{W}{mW + C}
$$

\n
$$
= \frac{1}{m + \frac{C}{W}}
$$
...(6.10)

As the load increases, *C W* which is in denominator, decreases and hence mechanical advantage increases. In limiting case when *W* tends to infinity, $\frac{C}{T}$ *W* = 0 and hence *maximum mechanical advantage equals* ¹ *m* . The variation of mechanical advantage with respect to load is as shown in Fig. 6.2. From the law of machine.
 $P = mW + C$
 \therefore $MA = \frac{W}{mW + C}$
 $= \frac{1}{m + \frac{C}{W}}$

As the load increases, $\frac{C}{W}$ which is in denominator, decreases and hence mechanical

increases. In limiting case when W tends to infinity,

6.5 VARIATION OF EFFICIENCY

From eqn. (6.5), the efficiency of the machine is given by $\eta = \frac{MA}{VR}$. Using the eqn. (6.10),

$$
\eta = \frac{1}{VR} \times \frac{1}{m + \frac{C}{W}}
$$
...(6.11)

Since the velocity ratio (*VR*) is constant for a machine, variation of efficiency with load is similar to the variation of mechanical advantage with the load.

The *maximum efficiency* is approached as the load approaches infinity ($W \rightarrow \infty$) and its value is equal to $\frac{1}{VR} \times \frac{1}{m}$. The variation of the efficiency with load is shown in Fig. 6.3.

Example 6.2. In a simple machine, whose velocity ratio is 30, a load of 2400 N is lifted by an effort of 150 N and a load of 3000 N is lifted by an effort of 180 N. Find the law of machine and calculate the load that could be lifted by a force of 200 N. Calculate also:

- *(1) The amount of effort wasted in overcoming the friction,*
- *(2) Mechanical advantage, and*

(3) The efficiency.

Solution: Let the law of machine be

In the first case, $P = 150$ N, and $W = 2400$ N In the second case, $P = 180$ N and $W = 3000$ N $150 = 2400 \, m + C$...(1) $180 = 3000 \text{ m} + C$...(2) Subtracting eqn. (1) from eqn. (2), we get 30 = 600 *m* $m = 0.05$ Substituting this value in eqn. (1), we get $150 = 120 + C$ ∴ $C = 30$ Hence, **the law of machine is** $P = 0.05 \text{ W} + 30$...(3) When a force of 200 N is applied: From the law of machine (3), $200 = 0.05 W + 30$ $W = 3400 N$ **Example 6.2.** In a simple machine, whose velocity ratio. See the **Example 6.2.** In a simple machine, whose velocity ratio is 90. a load of 2400 N is lifted by the light by a force of 200 N. Catalate about the lind condit

 $P = mW + C$

Ideal effort is given by:

$$
P_i = \frac{W}{VR} = \frac{3400}{30} = 113.33 \text{ N}
$$

$$
\therefore \text{ Effect wasted in overcoming the friction}
$$
\n
$$
= P - P_i = 200 - 113.33
$$
\n
$$
= 86.67 \text{ N}
$$
\nMechanical advantage = $\frac{W}{P} = \frac{3400}{200}$

\n
$$
= 17
$$
\n
$$
= \frac{MA}{VR} = \frac{17}{30} = .5667
$$
\n
$$
= 56.67\%
$$

Example 6.3. In a lifting machine an effort of 150 N raised a load of 7700 N. What is the mechanical advantage? Find the velocity ratio if the efficiency at this load is 60%. If by the same machine, a load of 13,200 N is raised by an effort of 250 N, what is the efficiency? Calculate the maximum mechanical advantage and the maximum efficiency.

Solution: Effort, *P* = 150 N

Load, $W = 7700$ N

∴ **Mechanical advantage:**

$$
MA = \frac{W}{P} = \frac{7700}{150} = 51.33
$$

If the efficiency is 60%,

$$
\eta = 0.6
$$

\n
$$
\eta = \frac{MA}{VR}
$$

\n
$$
0.6 = \frac{51.33}{VR}
$$

\nor
\n
$$
VR = \frac{51.33}{0.6} \quad i.e., \quad VR = 85.55
$$

When an effort of 250 N raised a load of 13,200 N,
 $W = 13,200$

$$
MA = \frac{W}{P} = \frac{13,200}{250} = 52.8
$$

 $\therefore \quad \eta = \frac{MA}{VR} = \frac{52.8}{85.55} = 0.6172$
i.e., $\eta = 61.72\%$
I et the law of machine be

Let the law

\n**Mechanical advantage** =
$$
\frac{W}{P} = \frac{3400}{200}
$$

\n $= 17$ \n

\n\n**Efficiency** = $\frac{MA}{VR} = \frac{17}{30} = .5667$
\n $= 56.67\%$ \n

\n\n**Example 6.3.** In a lifting machine and effort of 150 N raised a load of 7700 N. What is the mechanical advantage? Find the velocity ratio if the efficiency at this load is 60%. If by the same machine, a mechanical advantage and the maximum efficiency.\n

\n\n**Solution:** Effort, $P = 150 \text{ N}$
\n $= 150 \text{ N}$
\n

∴ *m =* 0.01818 ∴ **Maximum mechanical advantage**

Example 6.4. The efforts required for lifting various loads by a lifting machine are tabulated below: Load lifted in N 100 200 300 400 500 600 Effort required in N 16.0 22.5 28.0 34.0 40.5 46.5 Determine the law of machine. If the velocity ratio is 25, calculate efficiency at each load and plot efficiency versus load curve. From this curve, determine the maximum efficiency.

Solution: Figure 6.4 shows the graph of effort *versus* load. From this figure, *C =* 10 N and slope 30 L

$$
m = \frac{30}{500} = 0.06
$$

∴ **The law of machine is**

 $\eta = \frac{MA}{VR}$ *W P VR W* $=\frac{W}{P} \times \frac{1}{VR} = \frac{W}{P} \times \frac{1}{25} = \frac{W}{25P}$

From the graph (Fig. 6.5) maximum efficiency is seen as 57%. Actually if it is plotted for infinitely large load, **maximum efficiency** will be equal to

$$
\frac{1}{m} \times \frac{1}{VR} = \frac{1}{0.06} \times \frac{1}{25} = 0.6667 = 66.67\%
$$

6.6 REVERSIBILITY OF A MACHINE

If the removal of effort while lifting results in lowering of the load, the machine is said to be **reversible.** The machine is said to be **self-locking** if the load is not lowered on removal of the effort.

For example, while lifting water from the well, the pot falls back if the effort to pull it up is removed whereas the screw jack used to lift the motor car will hold the car at the same position even if the application of the effort is stopped. Hence, the former is a reversible and later is a selflocking type simple lifting machine.

A simple lifting machine will be reversible or self-locking solely based on its efficiency. It can be shown that a lifting machine is reversible if its efficiency is greater than 50 per cent and selflocking if its efficiency is less than 50 per cent.

Let *W* – load being lifted

P – effort required

VR – Velocity ratio

D – distance moved by the effort

Then,

d – distance moved by the load

Input = $P \times D$

Output = $W \times d$

∴ Work lost in friction = *PD – Wd*

When effort is removed, the load can start moving down if it can overcome the frictional resistance $= PD - Wd$. Hence the condition for the reversibility is:

Wd > (*PD – Wd*) ∴ 2*Wd* > *PD W P d* $\left(\frac{W}{P}\right)\left(\frac{d}{D}\right) > \frac{1}{2}$ $\overline{1}$ $\Big($ $\overline{}$ $\overline{1}$ $\bigg($ 2 *MA ×* 1 1 $\frac{1}{VR}$ > $\frac{1}{2}$ *i.e.*, $\eta > \frac{1}{2}$ $\frac{1}{2}$ or 50%.

Hence, a machine is reversible if its efficiency is greater than 50%.

Example 6.5. In a lifting machine in which velocity ratio is 30, a load of 5000 N is lifted with an effort of 360 N. Determine whether it is self-locking or reversible machine. How much is the frictional resistance?

∴ Work lost in friction =
$$
PD - Wd
$$

\nWhen effort is removed, the load can start moving down if it can overcome the resistance = $PD - Wd$. Hence the condition for the reversibility is:

\n $Wd > (PD - Wd)$

\n $\left(\frac{W}{P}\right)\left(\frac{d}{D}\right) > \frac{1}{2}$

\n $\frac{MA}{VR} > \frac{1}{2}$

\ni.e., $\eta > \frac{1}{2}$ or 50%.

\nHence, a machine is reversible if its efficiency is greater than 50%.

\nExample 6.5. In a lifting machine in which velocity ratio is 30, a load of 5000 N is lift of 500 N. Determine whether it is self-locking or reversible machine. How m frictional resistance?

\nSolution: VR = 30

\n $W = 5000 \text{ N}$

\n $P = 360 \text{ N}$

\n $MA = \frac{W}{P} = \frac{5000}{360} = 13.889$

\nefficiency, $\eta = \frac{M}{VR} = \frac{13.889}{30} = 0.4630$

\n $= 46.30\%$

\nSince the efficiency is less than 50%, it is self-locking machine.

\nJdeal load, $W_i = P \times V$ is 10,800 - 5000 = 5,800 N

\n∴ Frictional resistance = $W_i - W = 10,800 - 5000$ = 5,800 N

\n6.7 **LEVER ARM**

\nProbability is less than 50% is inserted by man. It is a strong rod used to slightly if the for pushing A body is inserted by man. It is a strong rod used to slightly if the for pushing A body is inserted very close to the body (say at the top of the body) is 50% to be slightly a heavy body. Figure 6.6 shows a lever arm *AB* used for sliding a block. Let *W* be to be slightly if the for pushing A body as a three about which the lever is to be

Ideal load, $W_i = P \times VR$ $= 360 \times 30$ $= 10,800$ N

∴ **Frictional resistance** = $W_i - W = 10,800 - 5000$

$$
= 5,800
$$
 N

6.7 LEVER ARM

Probably lever arm is the first machine to be invented by man. It is a strong rod used to push or lift slightly a heavy body. Figure 6.6 shows a lever arm *AB* used for sliding a block. Let *W* be the weight to be slightly lifted for pushing. A rigid body is inserted very close to the body (say at *C*) between the rod (lever) and floor. The rigid body acts as a hinge about which the lever is to be rotated. If '*a*' is the effort arm and '*b*' is the load arm, obviously.

6.8 PULLEYS

A systematic arrangement of one or more pulleys may provide a simple and convenient lifting machine. In its simplest form, it consists only one pulley over which a rope or chain passes as shown in the Fig. 6.7. In this case, velocity ratio is equal to one since distance moved by effort is equal to the distance moved by the load. It just performs changing the direction of the applied force.

Depending on the arrangement, pulleys are classified as:

- (1) First order pulley system
- (2) Second order pulley system
- (3) Third order pulley system.

At times, it may be difficult or may be detour to find velocity ratio directly. In such cases ideal conditions may be assumed (neglecting friction) and mechanical advantage may be found first. Then applying $VR = MA$ for ideal machine, the velocity ratio is found. This method of determining velocity ratio is used for various pulley systems considered here.

First Order Pulley System

A first order pulley system is shown in the Fig. 6.8. Pulley No. 1 is fixed in position to a support at top. A rope passes over this pulley and one end of this rope is tied to the support at the top, making a loop, in which pulley No. 2 is suspended and effort is applied at the other end. One end of another rope is tied to pulley No. 2 and the other end to the top support which makes a loop in which pulley No. 3 is suspended. Similarly, a number of pulleys can be arranged as shown, when an effort is applied to lift the load except first pulley all other pulleys move vertically. Therefore, first pulley is termed as fixed pulley and the others as movable pulleys. a loop, in which pulley No. 2 is supposeded and effort is applied at the other end. One ead,
one is tied to pulley No. 2 and the other end to the top support which makes a loop in w/
No. 3 is supposed. Similarly, a namber

Let an effort P be applied to lift a load W . In an ideal pulley system (friction = 0), the rope which passes over pulley No. 4 is subjected to a tension *^W* 2 . Then tension in rope which passes over

pulley No. 3 is $\frac{W}{A}$ 4 and tension in the rope which passes over pulley No. 2 is $\frac{W}{I}$ 8 . Hence, an effort

equal to $\frac{W}{8}$ is required to lift a load *W*.

$$
\therefore \qquad P = \frac{W}{8} \qquad \text{and hence } MA = \frac{W}{P} = \frac{W}{W} = 8 = 2^3
$$

But in an ideal machine, $VR = MA$ ∴ $VR = 2³$

8

It is to be noted that in the system considered, there are three movable pulleys and the velocity ratio is $2³$. If there are only two movable pulleys, then velocity ratio would be 4 (*i.e.*, $2²$). In general, in the first order pulley system, velocity ratio (VR) is given by 2^n , where, *n* is the number of movable pulleys present in the system. Thus, in first order pulley system

$$
VR = 2^n \tag{6.13}
$$

Second Order Pulley System

Figure 6.9 shows a second order pulley system. This system consists of a top pulley block and a bottom pulley block. In a pulley block pulleys may be arranged side by side or may be one below the other as shown in Fig. 6.9. The top pulley block is fixed in position to the top support whereas bottom pulley block can move vertically along with the load which is attached to it. One end of the rope is attached to the hook provided at the bottom of the top pulley block and the effort is applied at the other end. A single rope goes round all the pulleys. Let an effort *P* be applied to lift a load *W.*

Neglecting frictional losses, the tension in the rope all along the length is *P.* Take the section along (1)–(1) and consider the equilibrium of the bottom pulley block.

The load *W* is lifted using six ropes having equal tension *P*

$$
W = 6P, \quad \therefore \quad P = \frac{W}{6}, \quad \therefore \quad MA = \frac{W}{P} = \frac{W}{\frac{W}{6}} = 6.
$$

But *VR* = *MA* in ideal condition

$$
VR = 6.
$$

In general, in the second order pulley system velocity ratio is equal to twice the number of movable pulleys in the system.

That is, $VR = 2n$ (6.14)

where, *n* is total number of movable pulley in the system.

Third Order Pulley System

The arrangement of the pulleys in the third order system is shown in the Fig. 6.10. In this system a pulley (No. 1) is fixed to the top support, over which a rope passes. One end of the rope is attached to a rigid base at the bottom. The other end is attached to a second pulley. Over this pulley another rope passes, whose one end is attached to the same rigid base and the other end to a third pulley as shown. Likewise a series of pulleys can be arranged. The load to be lifted will be attached to the rigid base.

Referring to the Fig. 6.10, let the effort required be *P* to lift a load *W*. Then neglecting friction, Tension in the rope which passes over pulley No. $3 = P$ Tension in the rope which passes over pulley No. 2 = 2*P* Tension in the rope which passes over pulley No. 1 = 4*P*

∴ A total force of 7*P* is acting on the base.

∴ Lifting force produced on the base = 7*P*

Considering the equilibrium of rigid base,

$$
7P = W \qquad \therefore \qquad MA = \frac{W}{P} = 7
$$

But in an ideal machine, $VR = MA$, and hence $VR = 7$.

It can be easily seen that, if there are only two pulleys, $VR = 3$ and if there is only one pulley, $VR = 1$. Therefore, in general, for the third order pulley system:

$$
VR = 2^n - 1 \tag{6.15}
$$

where, $n =$ number of pulleys.

6.9 WHEEL AND AXLE

This machine consists of an axle *A* having diameter *d* and a wheel *B* having diameter *D* ($D > d$) co-axially fitted as shown in Fig. 6.11. The whole assembly is mounted on ball bearing so that wheel and axle can be rotated.

One end of a rope is tied to the pin provided on the wheel and the rope is wound around the wheel. The other end of the rope provides the means for the application of the effort. One end of another rope is tied to the pin provided on the axle and wound around the axle in the opposite direction to that of rope wound to the wheel. To the other end of this rope the load is attached. If the whole assembly is rotated, one rope gets wound up and the other gets unwound.

Suppose the assembly is moved by one complete revolution, then the distance moved by the effort = π *D* and distance moved by the load = π *d*

6.10 WHEEL AND DIFFERENTIAL AXLE

An improvement over wheel and axle machine is made by using one more wheel of bigger diameter and it is called wheel and differential axle.

This consists of a differential axle of diameter d_1 and d_2 ($d_1 < d_2$) and a wheel of diameter *D*, fixed uniaxially as shown in Fig. 6.12. One end of the rope is tied to the pin provided on the axle portion having diameter d_1 and a part of the rope is wound around it. The other end of the rope is wound around axle on the portion having diameter d_2 in the opposite direction. This pattern of winding forms a loop and a simple pulley is installed in this loop as shown in the figure. The load is attached to this pulley. The second rope is wound to the wheel in such a direction that if it is unwound, the rope around the bigger diameter axle gets wound up and the rope around smaller diameter axle gets unwound. One end of this second rope provides means for application of the effort. winding from a loop and a simple pulley's installed in his loop as Sown in the figure
is stated to this pulley. The second rope is wound to the wheel in such a direction
unwaund, the rupe around the higger diameter sade g

Suppose the whole system makes one complete revolution due to the applied effort, then Total distance moved by the effort at the differential axle = πD

Length of winding of the rope = πd_2 Length of unwinding of rope = πd_1 \therefore Net wound length $= \pi d_2 - \pi d_1$ $= \pi(d_2 - d_1)$

But, the rope is continuous and the load is to be lifted by the pulley block in the loop.

∴ Total height over which pulley is lifted = $\frac{\pi(d_2 - d_1)}{2}$ 2 −

and hence the distance moved by the load = $\frac{\pi(d_2 - d_1)}{2}$ −

$$
VR = \frac{\pi D}{\frac{\pi (d_2 - d_1)}{2}}
$$

$$
VR = \frac{2D}{d_2 - d_1}
$$

Hence, velocity ratio in wheel and differential axle is given by:

$$
VR = \frac{2D}{d_2 - d_1} \tag{6.17}
$$

6.11 WESTON DIFFERENTIAL PULLEY BLOCK

This is a special type of simple pulley system. It is shown in the Fig. 6.13.

This system consists of two pulley blocks, one at the top attached to the support and the other at the bottom hanging in the chain loop. The top block consists of two wheels of different diameters, but fixed co-axially. The bottom block is a simple pulley to which the load *W* is attached. An endless chain is wound around the pulley system as shown in the figure. All the wheels are made with teeth so as to accommodate the links of the chain. The chain is essentially used to avoid slipping.

To determine velocity ratio of the system, let us consider pulley block as an ideal machine and determine its mechanical advantage first. In ideal machine $VR = MA$. Let the diameter of the larger wheel of the top block be *D* and the diameter of the smaller wheel of the top block be *d* [Fig. 6.14(*a*)]. Let the effort required to lift the load *W* be *P.*

Then the tension in the chain loop in which pulley is hanging is *W* $\frac{1}{2}$. Now, taking moment about the axis of top block [Fig. 6.14(*b*)],

$$
\frac{W}{2} \times \frac{D}{2} = \frac{W}{2} \times \frac{d}{2} + P \times \frac{D}{2}
$$

$$
\frac{W}{4}(D-d) = \frac{PD}{2}
$$

$$
\frac{W}{P} = \frac{2D}{(D-d)}
$$

i.e.,

$$
MA = \frac{2D}{(D-d)}
$$

In an ideal machine,

∴

$$
VR = MA = \frac{2D}{(D - d)}
$$
...(6.18)

6.12 INCLINED PLANE

Inclined plane is a very simple lifting device. The lift is essentially accomplished with the horizontal displacement.

An inclined plane consists of a plane surface at a definite angle over which the load is to be lifted.

Figure 6.15 shows a typical inclined plane. Here the load is a roller which is to be lifted to a higher elevation. One end of a rope is tied to the roller and the rope is passed over a pulley attached at the top of the inclined plane. At the other end of the rope the effort is applied.

Let the angle of inclination of the plane be θ and the length of the inclined plane be *L*. Then, if the roller is made to roll from bottom to top, applying an effort *P,* the load is lifted through a height of *L* sin θ. In this process the effort *P* moves through a distance *L* vertically downwards.

[∴] *VR* ⁼ distance moved by the effort distance through which load is lifted ⁼*^L ^L* sin sin θ θ ⁼ ¹ [∴] *VR* ⁼ ¹ sin ^θ ...(6.19) Civil[datas.blogspo](http://Civildatas.blogspot.in)t.in

Example 6.6. In a first order system of pulleys there are three movable pulleys. What is the effort required to raise a load of 6000 N ? Assume efficiency of the system to be 80%.

If the same load is to be raised using 520 N, find the number of movable pulleys that are necessary.

Assume a reduction of efficiency of 5% for each additional pulley used in the system.

Solution: $VR = 2^n$, where *n* is the number of movable pulleys. $VR = 2^3 = 8$ Now, $MA = \eta \times VR$ $= 0.8 \times 8$ $= 6.4$ *i.e.,* $\frac{W}{P} = 6.4$ ∴ $P = \frac{W}{6.4} = \frac{6000}{6.4}$

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i.e., **P = 937.5 N** In the second case, Effort = 520 N Efficiency $\eta = 0.80 - n_1 \times 0.05$ where n_1 = number of additional pulleys required and equal to $(n - 3)$. $MA = \eta \times VR$ *i.e., W ^P* = η *× VR* ∴ $W = P \times \eta \times 2^n$ $= P(0.8 - n_1 \times 0.05) \times 2^n$ $= P[0.8 - (n-3) \times 0.05] 2^n$ Efficiency

where $n_1 = \text{number of additional pulleys required and equal to } (n - 3)$.
 $M = 1 \times VR$
 $\frac{W}{P} = 1 \times VR$
 \therefore $W = P \times 1 \times 2^n$
 $= P(0.8 - n_1 \times 0.05) \times 2^n$

By going for a trial and error solution, starting with one additional pulley *i.e.*, totall;

pull

By going for a trial and error solution, starting with one additional pulley *i.e.,* totally with four pulleys,

$$
W = 520 [0.8 - (4 - 3) \times 0.05] 2^4 = 6240 N
$$

i.e., if four pulleys are used, a load of 6240 N can be raised with the help of 520 N effort.

∴ **Number of movable pulleys required = 4**

Example 6.7. What force is required to raise the load W shown in Fig. 6.16 ? Assume efficiency of the system to be 85%.

Solution: The pulley system shown in the Fig. 6.16 is a variation of the second order pulley system.

Fig. 6.16

VR = 2 × number of movable pulleys = 6
\n*MA* =
$$
\eta \times VR
$$

\n= 0.85 × 6 = 5.1
\ni.e., $\frac{W}{P} = 5.1$
\n $P = \frac{W}{5.1} = \frac{12,000}{5.1} = 2352.94 N$

Example 6.8. Find the pull required to lift the load W shown in Fig. 6.17(a) assuming the efficiency of the system to be 78%.

Solution: The pulley system shown in Fig. 6.17(*a*) is a combination of a first order system and a second order system as shown in Figs. 6.17(*b*) and (*c*).

Let load *W* be lifted by a distance *x*. Consider the first order system portion [Fig. 6.17(*b*)]. Here there are two movable pulleys. Hence

 $VR = 2^2 = 4$

In this portion *P* moves by 4*x*.

Now, consider the second order pulley system portion [Fig. 6.17(*c*)]. Here there are two movable pulleys. Hence $VR = 2 \times 2 = 4$.

∴ Distance moved by the effort in this system = 4*x*.

Hence, the total distance moved by the effort in the given system $= 4x + 4x = 8x$

$$
VR = \frac{8x}{x} = 8
$$

Now,

$$
MA = \eta \times VR = 0.78 \times 8
$$

$$
= 6.24
$$

i.e.,

$$
\frac{W}{P} = 6.24
$$

$$
P = \frac{12,000}{6.24} = 1923.08 \text{ N}
$$

Example 6.9. A lifting machine consists of pulleys arranged in the third order system. There are three pulleys in the system. A load of 1000 N is lifted by an effort of 180 N. Find the efficiency of the machine and the effort lost in friction.

Solution: For the third order system of pulleys,

 $VR = 2^n - 1$ where, n is the number of pulleys in the system. $VR = 2^3 - 1 = 7$

Now, *MA =* η × *VR i.e.,* $\frac{W}{P} = \eta \times VR$ ∴ $η = \frac{W}{P} \times \frac{1}{VR}$ $=\frac{1000}{180}$ 1 $\times \frac{1}{7} = 0.7937$ *i.e.*, $η = 79.37%$ Now, Ideal effort, $=\frac{W}{VR}$ $=\frac{1000}{7}$ $= 142.86$ N ∴ **Effort lost in friction** $= P - P_i$ $= 180 - 142.86$ **= 37.14 N**

Example 6.10. What force P is required to raise a load of 2500 N in the system of pulleys shown in Fig. 6.18(a). Assume efficiency of the system to be equal to 70%.

Solution: Figure 6.18(*a*) can be split into two simple systems as shown in Figs. 6.18(*b*) and 6.18(*c*). What is shown in Fig. 6.18(*b*) is a third order pulley system having two pulleys.

Figure 6.18(*c*) is also a third order system, having two pulleys.

Example 6.11. In a wheel and axle, diameter of the wheel is 500 mm and that of the axle is 200 mm. The thickness of the cord on the wheel is 6 mm and that of the axle is 20 mm. Find the velocity ratio of the machine. If the efficiency when lifting a load of 1200 N with a velocity of 10 metres per minute is 70%, find the effort necessary.

Solution: Effective wheel diameter =
$$
\frac{6}{2} + 500 + \frac{6}{2}
$$

= 506 mm
Effective axle diameter = $\frac{20}{2} + 200 + \frac{20}{2}$
= 220 mm.

For a wheel and axle, the velocity ratio is given by *^D d*

$$
\therefore \qquad \qquad VR = \frac{506}{220} = 2.30
$$

Mechanical advantage = Efficiency \times velocity ratio

$$
= 0.7 \times 2.30 = 1.61
$$

$$
MA = \frac{W}{P}
$$

$$
P = \frac{1200}{1.61}
$$

i.e.,

$$
P = 745.34 N
$$

Example 6.12. A load of 20 kN is to be lifted by a differential wheel and axle. It consists of differential axle of 250 mm and 300 mm diameter and the wheel diameter is 800 mm. Find the effort required if the efficiency of the machine is 55%.

Solution: Differential axle diameters,

Load, $W = 5000 \text{ N}$ Efficiency, $\eta = 60\%$ Velocity ratio = 2*D D* − *d* $=\frac{2 \times 500}{500 - 200}$ $\frac{1}{0}$ \times 500
 $\frac{1}{0}$ = 3.33

Mechanical advantage = Efficiency \times Velocity ratio

$$
= 0.6 \times 3.33 = 2
$$

Effort required, $P = \frac{W}{MA} = \frac{5000}{2}$
i.e., $\mathbf{P} = 2500 \text{ N}$

6.13 SCREW JACK

This is a device commonly used to lift heavy loads. Screw jack works on the principle same as that of inclined plane. A typical section of the screw jack is shown in the Fig. 6.19.

The device consists of a nut and a screw. Monolithically cast nut and stand form the body of the jack. The load is carried by the screw head fitted onto the screw as shown in the figure. The body (consisting of nut) is fixed and the screw is rotated by means of a lever.

The axial distance moved by the nut (or by the screw, relative to each other) when it makes one complete revolution is known as **lead of the screw head.** The distance between consecutive threads is called *pitch* (of a screw thread). If the screw is single threaded, then lead of the screw is equal to the pitch. If the screw is double threaded then lead of the screw is twice the pitch.

Let R be the length of the lever and d be the mean diameter of the screw.

Let a load *W* be lifted using an effort *P.*

If an effort *P* is applied at the lever end, it is equivalent to an effort P_1 at the screw [Fig. 6.20(*a*)] and P_1 is given by the condition:

$$
P \times R = P_1 \times \frac{d}{2}
$$

$$
P_1 = \frac{2PR}{d}
$$

Now, consider one complete revolution of the lever. The load *W* is lifted up by a distance *p* equal to the lead of the screw.

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 $\ldots(6.20(b))$

where, tan $\theta = \frac{p}{\pi d}$

If the load is descending, then the friction will be acting in the reverse direction so that the resultant reaction *R* shifts as shown in Fig. 6.20(*d*).

Then eqn. 6.20 changes to

$$
P = \frac{d}{2R}W \tan (\theta - \phi)
$$

Torque required, $= T = PR$

$$
= \frac{d}{2}W \tan{(\theta + \phi)}
$$

Hence torque required while ascending

$$
T = \frac{d}{2} W \tan (\theta + \phi) \tag{6.21}
$$

and torque required while descending

$$
T = \frac{d}{2} W \tan (\theta - \phi) \tag{6.21(a)}
$$

Now,

$$
VR = \frac{\text{Distance moved by the effort}}{\text{Distance moved by the load}}
$$

= $\frac{2\pi R}{p}$...(6.22)

Example 6.14. A screw jack raises a load of 40 kN. The screw is square threaded having three threads per 20 mm length and 40 mm in diameter. Calculate the force required at the end of a lever 400 mm long measured from the axis of the screw, if the coefficient of friction between screw and nut is 0.12. resultant reaction R shifts as shown in Fig. 6.20(d)

Then eqn. 6.20 changes to
 $P = \frac{d}{2R}W \tan (0 - \theta)$

Torque required. = $T = PR$
 $= \frac{d}{2}W \tan (0 + \phi)$

Hence torque required. = $T = \frac{d}{2} W \tan (0 + \phi)$

Hence torque required w

Solution: Screw diameter, $d = 40$ mm

Lead of the screw,
$$
p = \frac{20}{3} = 6.667 \, \text{mm}
$$

\nLoad, $W = 40 \, \text{kN}$

\nLower length, $R = 400 \, \text{mm}$

\n $\mu = 0.12$

We have

$$
P = \frac{d}{2R} W \frac{\mu + \tan \theta}{1 - \mu \tan \theta}
$$

and $\tan \theta = \frac{p}{\pi D}$

$$
= \frac{6.667}{\pi \times 40} = 0.05305
$$

 \therefore $P = \frac{40}{2 \times 400} \times 40,000 \left[\frac{0.12 + 0.05305}{1 - (0.12 \times 0.05305)} \right]$
i.e., $\mathbf{P} = 348.32 \text{ kN}$

Example 6.15. A screw jack has square threads 50 mm mean diameter and 10 mm pitch. The load on the jack revolves with the screw. The coefficient of friction at the screw thread is 0.05.

 $\overline{}$ $\overline{}$

(*1*) *Find the tangential force required at the end of 300 mm lever to lift a load of 6000 N.*

(*2*) *State whether the jack is self-locking. If not, find the torque which must be applied to keep the load from descending.*

and tan θ =
$$
\frac{p}{\pi D}
$$

\n= $\frac{6.667}{\pi \times 40} = 0.05305$
\n \therefore $P = \frac{40}{2 \times 400} \times 40,000 \left[\frac{0.12 + 0.05305}{1 - (0.12 \times 0.05305)} \right]$
\ni.e., **P** = 348.32 kN
\nExample 6.15. A screw jack has square threads 50 mm mean diameter and 10 mm pitch. The load
\non the jack revolves with the care. The coefficient of friction at the score where is 0.05.
\n(1) Find the tangential force required at the end of 300 mm level to load of 6000 N.
\n(2) State whether the jack is self-locking. If not, find the torque which must be applied to keep the load from descending.
\nSolution:
\n $\tan \theta = \frac{p}{\pi d} = \frac{10}{\pi \times 50} = 0.0637$
\n $\therefore \qquad \theta = 3.6426^{\circ}$
\n $\therefore \qquad \theta = 2.8624^{\circ}$
\n(1) $P = \frac{d}{2R} \times W \tan (\theta + \phi)$
\n $= \frac{50}{2 \times 300} \times 6000 \tan (3.6426^{\circ} + 2.8624^{\circ})$
\n $P = 57.01 \text{ N}$
\n(2) We have
\n $WR = \frac{2\pi R}{p} = \frac{2\pi \times 300}{10}$
\n= 188.496
\n $MA = \frac{W}{p} = \frac{6000}{57.01}$
\n= 105.245
\nEfficiency = $\frac{MA}{VR} = \frac{105.245}{188.496}$
\n= 0.5583

i.e., Efficiency = 55.83\% > 50.

Hence the screw jack is not self-locking.

∴ The torque required to keep the load from descending

$$
= \frac{d}{2} W \tan (\theta - \phi)
$$

= $\frac{50}{2} \times 600 \times \tan (3.6426^\circ - 2.8624^\circ)$
T = 204.3 N-mm

6.14 DIFFERENTIAL SCREW JACK

Differential screw jack is an improvement over simple screw jack. A typical differential screw jack is shown in Fig. 6.21. It consists of two threaded elements *A* and *B.* Both *A* and *B* have threads in the same direction (right-handed). The element *A* is a cylinder which has threads on both its outer and inner surfaces. The threads on the outer surface of the element *A* fits into the nut *C* which also functions as the base of the whole mechanism. The threads on the element *B* fit into the threads cut on the inner surface of *A.* Thus, the element *A* acts as a screw for the nut *C* and also as a nut for the element *B.* With the help of a lever inserted in the holes made on the top of the block *D,* which is attached to the element *B,* block *D* can be rotated. When *D* is rotated, *A* rotates with it. Rotation of *B* is prevented by suitable arrangement.

Let *D* and *d* be the mean diameters of the screws *A* and *B,* respectively.

Let p_A and p_B be the pitch of the screws *A* and *B*, respectively and p_A be greater than p_B .

If the lever is rotated through one complete revolution, the height through which the element *A* moves $up = p_A$

In the mean time, the element *B* moves down with respect to *C.*

The distance through which *B* comes down = p_B

∴ Net height through which load is lifted = $(p_A - p_B)$

Let *R* be the radial distance (from the centre line of *A* and *B*) at which an effort *P* is applied

Now,
$$
VR = \frac{\text{Distance moved by the effort}}{\text{Distance moved by the load}} = \frac{2\pi R}{p_A - p_B}
$$
...(6.23)

It can be seen from eqn. (6.23) that the velocity ratio in the differential screw jack is increased as compared to that of simple screw jack (eqn. 6.22).

Example 6.16. The following are the specifications for a differential screw jack:

(*1*) *Pitch of smaller screw, 5.0 mm*

(*2*) *Pitch of larger screw, 10.0 mm*

(*3*) *Lever arm length from centre of screw = 500 mm.*

The screw jack raises a load of 15 kN with an effort of 185 N. Determine the efficiency of the differential screw jack at this load.

If the above jack can raise a load of 40 kN with an effort of 585 N, determine the law of machine.

Solution: Now, $p_A = 10.0$ mm

$$
p_B = 5.0 \text{ mm}
$$

Lever arm length, $R = 500$ mm

It can be seen from eqn. (6.23) that the velocity ratio in the differential screw jack is
\nas compared to that of simple screw jack (eqn. 6.22).
\nExample 6.16. The following are the specification of a differential screw jack:
\n(1) Pitch of smaller screw, 5.0 mm
\n(2) Pitch of larger screw, 10.0 mm
\n(3) Lower arm length from centre of screw = 500 mm.
\nThe screw jack raises a load of 15 kN with an effort of 185 N. Determine the ej
\nthe differential screw index at this load.
\n
$$
If the above jack can raise a load of 40 kN with an effort of 585 N, determine\nmachine.\nSolution: Now, $p_A = 10.0$ mm
\n $p_B = 5.0$ mm
\nLever arm length, $R = 500$ mm
\n
$$
VR = \frac{2\pi R}{p_A - p_B} = \frac{2\pi \times 500}{10 - 5.0} = 628.32
$$
\n
$$
= \frac{M}{10 - 5.0} = 628.32
$$
\n
$$
= 0.129 = 12.9
$$
 per cent.
\nTo find law of machine:
\nLet law of machine:
\nLet law of machine:
\nLet law of machine is $P = mW + C$
\nFrom second case: 185 = m × 50,000 + C
\n(2) minus (1) gives, 400 = 35,000 m.
\nor $m = \frac{4}{350}$
\nSubstituting this value in eqn. (1), we get
\n $185 = 171.43 + C$
\n $C = 13.57$ N
\n \therefore Law of machine is $P = \frac{4}{350}$ W + 13.57
\n6.15 WINCH CRABS
\nWinch reads are lifting machines in which velocity ratio is increased by a gear system.
\nNote of years is used the which crash is called a sine perrebase wind, graph and if two se
$$

To find law of machine:

Let law of machine be $P = mW + C$ From first case: $185 = m \times 15,000 + C$...(1) From second case: $585 = m \times 50,000 + C$...(2) (2) minus (1) gives, 400 = 35,000 *m.*

or
$$
m = \frac{4}{350}
$$

Substituting this value in eqn. (1), we get

$$
185 = 171.43 + C
$$

$$
C = 13.57 \text{ N}
$$

$$
\therefore
$$
 Law of machine is P = $\frac{1}{350}$ W + 13.57

6.15 WINCH CRABS

Winch crabs are lifting machines in which velocity ratio is increased by a gear system. If only one set of gears is used, the winch crab is called a **single purchase winch crab** and if two sets are used it is called **double purchase winch crab.**

Single Purchase Winch Crab

Line diagram of a single purchase winch crab is shown in Fig. 6.22. It consists of a load drum of radius *r* connected to an axle by gears. The toothed wheel on load drum is called **spur wheel** and the toothed wheel on axle is called **pinion.** Pinion is always smaller in size and it contains less number of teeth as compared to that on the spur wheel. The axle is provided with a handle of arm length *R*. Let the number of teeth on pinion and spur wheel be T_1 and T_2 , respectively. Let the effort be applied at the end of the handle. When one revolution is made, the distance moved by the effort is given by:

$D = 2πR$

When axle makes one revolution, due to gear arrangement load drum moves by T_1 number

∴ The distance over which the load moves:

$$
d = 2\pi r \times \frac{T_1}{T_2}
$$

∴ Velocity ratio,

Double Purchase Winch Crab

Velocity ratio of a winch crab can be increased by providing another axle with a pair of pinion and gear as shown in Fig. 6.23. Since two pairs of pinion and gear are used it is called a double purchase winch crab. This is used for lifting heavier loads.

Let the number of teeth on various wheels be T_1 , T_2 , T_3 and T_4 as shown in Fig. 6.23. Let the handle makes one revolution.

Distance moved by effort *P* is given by:

 $D = 2\pi R$...(6.25)

When axle *A* makes one revolution, axle *B* is moved by T_1 teeth, *i.e.*, it makes $\frac{T_1}{T_2}$ 1 2 revolutions.

The number of teeth by which spur wheel is moved is $\frac{7}{5}$ *T* 1 2 $\times T_3$ and hence load drum makes

T T 1 $\left(\frac{T_1}{T_2}\right)$ \overline{a} $\overline{\mathcal{C}}$ $\bigg\} \times \bigg(\frac{T}{T} \bigg)$ $\overline{1}$ *T* 3 $\left(\frac{T_3}{T_4}\right)$ \overline{a} $\overline{\mathcal{L}}$  revolutions. $\overline{}$

∴ The distance moved by the load

The number of teeth is moved by the load
\n
$$
\frac{1}{\sqrt{1-\frac{1}{2}}}
$$
\nWhen axle *A* makes one revolution, axle *B* is moved by *T*₁ teeth, i.e., it makes $\frac{T_1}{T_2}$ revolutions.
\nThe number of teeth by which spur wheel is moved is $\frac{T_1}{T_2} \times T_3$ and hence load drum makes $\left(\frac{T_1}{T_2}\right) \times \left(\frac{T_3}{T_4}\right)$ revolutions.
\nThe distance moved by the load
\n
$$
d = 2\pi r \times \left(\frac{T_1}{T_2}\right) \times \left(\frac{T_3}{T_4}\right)
$$
\n
$$
V = \frac{P}{d} = \frac{2\pi R}{2\pi r \times \left(\frac{T_1}{T_2}\right) \times \left(\frac{T_3}{T_4}\right)}
$$
\ni.e.,
$$
V = \frac{R}{r} \times \left(\frac{T_2}{T_1}\right) \times \left(\frac{T_3}{T_4}\right)
$$
\ni.e.,
$$
V = \frac{R}{r} \times \left(\frac{T_2}{T_1}\right) \times \left(\frac{T_3}{T_4}\right)
$$
\n
$$
L = \frac{R}{r}
$$
\n
$$
L = \frac{R}{
$$

Example 6.17. Following are the specifications of a single purchase crab:

Diameter of the load drum = 200 mm *Length of lever arm R* $= 1.2 m$

Number of teeth on pinion, T_1 = 10 *Number of teeth on spur wheel,* $T_2 = 100$. *Find the velocity ratio of this machine. On this machine efforts of 100 N and 160 N are required to lift the load of 3 kN and 9 kN, respectively. Find the law of the machine and the efficiencies at the above loads.* **Solution:** Radius of the load drum, $r = \frac{200}{2}$ 2 $= 100$ mm Length of lever arm, $R = 1.2$ m = 1200 mm Velocity ratio of the single purchase crab is given by: $VR = \frac{R}{A}$ *r T* $\times \frac{I_2}{T_1}$ 1 $=\frac{1200}{100}$ 100 $\times \frac{100}{10}$ **VR = 120** Let the law of machine be $P = mW + C$ In first case: $P = 100 \text{ N}; W = 3 \text{ kN} = 3000 \text{ N}$ ∴ $100 = m \times 3000 + C$...(1) In the second case: $P = 160$ N; and $W = 9$ kN = 9000 N ∴ $160 = m \times 9000 + C$...(2) Subtracting eqn. (1) from eqn. (2), we get 60 = 6000 *m* ∴ $m = \frac{1}{100} = 0.01$ Substituting this value of m in eqn. (1), we get to the above loads.

The above loads.

Solution: Radius of the load drum, $r = \frac{200}{2} = 100$ mm

Length of lever arm, $R = 1.2$ m = 1200 mm

Length of lever arm, $R = 1.2$ m = 1200 mm

Velocity ratio of the single purchase

$$
100 = \frac{1}{100} \times 3000 + C
$$

$$
C = 70
$$

Hence, the law of machine is

$$
P = 0.01 W + 70
$$

Efficiencies:

In the first case,

$$
MA = \frac{W}{P} = \frac{3000}{100} = 30
$$

$$
\eta = \frac{MA}{VR} = \frac{30}{120} = 0.25 = 25\%
$$

In the second case,

$$
MA = \frac{W}{P} = \frac{9000}{160} = 56.25
$$

Now,
\n
$$
\eta = \frac{MA}{VR}
$$
\n
$$
0.40 = \frac{MA}{32}
$$
\nor
\n
$$
MA = 12.8
$$
\n
$$
\frac{W}{P} = 12.8, \text{ but } P = 250 \text{ N}
$$
\n
$$
\therefore \qquad W = 12.8 \times 250
$$
\ni.e.,
\n
$$
W = 3200 \text{ N}
$$

 $= 32$

Applied effort lifts a load of 3200 N

IMPORTANT FORMULAE

1. Mechanical advantage = $\frac{\text{Load}}{\text{Edd}}$

Effort 2. Velocity ratio = Distance moved by effort

2. Velocity ratio =
$$
\frac{\text{Distance moved by elior}}{\text{Distance moved by load}}
$$

- **3.** Input = Effort \times Distance moved by effort.
- **4.** Output = Load \times Distance moved by the load.

5. Efficiency =
$$
\frac{\text{Output}}{\text{1}} = \frac{\text{Mechanical advantage}}{\text{1}}.
$$

Efficiency =
$$
\frac{S_{\text{target}}}{\text{Input}} = \frac{M_{\text{e}}}{\text{Velocity ratio}}
$$
.

6. Law of machine:

$$
P = mW + C.
$$

7. A machine is reversible, if $\eta > 50\%$.

- **8.** In case of first order pulleys, $VR = 2^n$, where *n* is the number of movable pulleys.
- **9.** In case of second order pulleys, $VR = 2n$, is the where *n* is number of pulleys.
- **10.** In case of third order pulleys $VR = 2^n - 1$, where *n* is number of pulleys.
- **11.** In wheel and axle,

$VR =$ $\frac{\text{Diameter of wheel}}{\text{Diameter of axle}}$.

12. In case of wheel and differential axle,

$$
VR = \frac{2D}{d_2 - d_1}.
$$

13. In Weston differential pulley block,

$$
VR = \frac{2D}{D - d}.
$$

14. In inclined plane

 $VR = \frac{1}{1}$ $sin θ$

15. In screw jack

 $VR = \frac{2\pi R}{p}$.

16. In double purchase winch crab

$$
VR = \frac{R}{r} \times \frac{T_2}{T_1} \times \frac{T_4}{T_3}.
$$

17. In single purchase winch crab

$$
VR = \frac{R}{r} \times \frac{T_2}{T_1}.
$$

THEORY QUESTIONS

- **1.** Define the following terms:
	- (*i*) Mechanical advantage

(*ii*) Velocity ratio

- (*iii*) Efficiency of a lifting machine.
- **2.** Explain the law of machine. How it is determined ?
- **3.** What is meant by reversibility of machine ? Prove that a machine is reversible if its efficiency is more than 50%. *VR* = 2*n*, the where *n* is analised to the whole control in the case of third order pulleys.

10. In case of wheel and either a jump here of pulleys.

11. In wheel and either of subset of axis

12. In case of wheel and

- **4.** With neat sketches explain the different types of pulley systems and derive the expressions for their efficiency.
- **5.** With neat sketches explain the following simple machines:
	- (*i*) Wheel and axle
	- (*ii*) Wheel and differential axle
	- (*iii*) Weston differential pulley block.
- **6.** Show that in a screw jack effort *P* required to lift a load *W* is given by *P R W* $2R-1$ μ + tan θ μ tan $θ$ + − tan tan

where ϕ is angle of friction and tan $\theta = \frac{p}{q}$ π*d* , *p* being the pitch of screw.

7. Give a neat sketch of

(*i*) Single purchase winch crab

(*ii*) Double purchase winch crab

PROBLEMS FOR EXERCISE

1. In a lifting machine an effort of 1.5 kN is to be moved through a distance of 7.2 m to move a load of 24 kN through a distance of 300 mm. Determine: (1) mechanical advantage, (2) velocity ratio, (3) efficiency, (4) ideal effort, (5) effort lost in friction, (6) ideal load and (7) frictional resistance.

[**Ans.** (1) 16; (2) 24; (3) 66.67%; (4) 1.0 kN; (5) 0.5 kN; (6) 36 kN and (7) 12 kN]

- **2.** In a lifting machine an effort of 400 N is required to raise a load of 3000 N and an effort of 640 N, to raise a load of 5000 N. How much load can be lifted with an effort of 760 N? If the velocity ratio is 16, determine the efficiency of the machine when an effort of 760 N is applied. Is it a reversible machine? [Ans. $W = 6000 \text{ N}; \eta = 49.34$; It is not reversible] (*ii*) Where and after

(*iii*) Where and differential axie

(*iii*) Weston differential axie

(*iii*) Weston differential pulley block.

6. Show that in a serew jack effort *P* required to lift a load *W* is given by $P =$
	- **3.** The following observations were made in an experiment on a lifting machine:

Load in N 500 1000 1500 2000 2500 3000

Draw the load *versus* effort graph and determine the law of machine. If the velocity ratio is 30 determine the efficiency while lifting a load of 1800 N.

What is the maximum efficiency of the machine ?

[**Ans.** *P =* 0.04*W +* 6; η = 51.28%; maximum η = 83.33%]

4. A lifting machine having velocity ratio 28 starts raising a load of 6420 N on applying an effort of 450 N to it. If suddenly the effort is removed find whether the load comes down or not?

[Ans. $\eta = 50.95\%$, hence the load comes down]

- **5.** In the first order pulley system having three movable pulleys, how much effort is required to raise a load of 5780 N if the same system raises a load of 1200 N with an effort of 200 N? Assume the efficiency to be constant for the pulley system. [Ans. $P = 963.33 \text{ N}$]
- **6.** For the arrangement of pulleys shown in Fig. 6.24, find the effort required to raise the given load 7280 N. Assume efficiency of the system as 75%. [**Ans.** *P =* 2436.7 N]

Fig. 6.24

7. For the combination of first order and second order pulley system shown in Fig. 6.25 what will be velocity ratio?

Assuming efficiency to be 80%, calculate what effort is required to raise the load of 8200 N.

[Ans. $VR = 8$; $P = 1281.25$ N]

- **8.** For a third order pulley system having six movable pulleys, an effort of 720 N is required to raise a load of 30 kN. Calculate the efficiency of the system. [**Ans.** 66.14%]
- **9.** For a wheel and axle, the following details are available: Diameter of the wheel = 540 mm Diameter of the $axle = 270$ mm Thickness of the cord on the wheel $= 6$ mm Thickness of the cord on the $axle = 20$ mm

Calculate the efficiency of the device if an effort of 725 N is required to lift a load of 1000 N.

[Ans.
$$
\eta = 73.26\%]
$$

- **10.** A differential axle and wheel consists of a differential axle of 240 mm and 320 mm diameter; and a wheel of diameter 750 mm. Assuming efficiency to be 57%, find the effort required to raise a load of 24 kN. [**Ans.** *P =* 2.2456 kN]
- **11.** A Weston differential pulley block of diameters 400 mm and 800 mm is used to lift a load of 40,000 N. Find the effort required if efficiency of the system is 60%. [**Ans.** 26,667 N]
- **12.** The following particulars refer to a screw jack:
	- (1) Diameter of the screw rod = 62.5 mm
	- (2) Length of the handle $= 250$ mm
	- (3) Pitch of the square threads $= 12.5$ mm
	- (4) Coefficient of friction = 0.05

Find the effort required to lift a load of 5000 N. Find also the efficiency of the jack.

[Ans.
$$
P = 71.27
$$
 N; $\eta = 55.83\%$]

- **13.** A screw jack carries a load of 10 kN. It has a square threaded screw of pitch 25 mm and mean diameter 60 mm. The coefficient of friction between screw and nut is 0.20. Calculate the torque required to raise the load and the efficiency of the screw. Find also the force required at the end of the handle 500 mm long to lower the load. [**Ans.** *T =* 102.5802 kN-mm; *P =* 0.2052 kN] (3) Pitch of the square linearises 12.5 mm

(4) Coefficient of friction = 0.05

Find the efficiency of the signs.

[Ans. $P = 71.27$ N; n)

13. A serew jack extracts a load of 10 kN. It has a square threaded serew of pitch
	- **14.** The following are the specifications of a single purchase crab:

Diameter of the load drum = 220 mm

Length of lever arm, $R = 1.0$ m

Number of teeth on pinion, $T_1 = 10$

Number of teeth on spur wheel, $T_2 = 120$

Determine the velocity ratio of the machine. On this machine, it is found in one trial that an effort of 120 N is required to lift a load 4.2 kN and in another trial an effort of 150 N is required to lift a load of 11.4 kN. Establish the law of machine. Also calculate efficiencies corresponding to 120 N and 150 N efforts.

$$
\left(\text{Ans. } VR = 109.091; \, P = \frac{1}{240}W + 102.5; \, \eta_1 = 32.08\%; \, \eta_2 = 69.67\%
$$

15. In a double purchase crab, the pinions have 12 and 18 teeth while the spur wheels have 24 and 36 teeth. The effective diameter of the drum is 142 mm. The effort handle is 450 mm. What load can be lifted by an effort of 220 N applied at the end of handle if efficiency of winch is 42%.

 $[Ans. $P = 2342.54 \text{ N}].$$

7

Physical and Mechanical Properties of Structural Materials

The principle materials used in building structures and machines include metals and their alloys, concrete, wood, stones, bricks and other clay products, bituminous materials and plastics. The products finally produced with these materials should be strong enough to resist the forces coming on it without undergoing deformation so as to put it out of use. The product should be durable. The required properties for different usage is different. Hence an engineer must know the required properties for his products. Various properties of materials may be grouped into: **Physical and Mechani Material Constrainer (Separation 1991) The principle materials used in building structurates and matching include translate matching products. The principle matching products with the materials show**

- (*a*) Physical Properties
- (*b*) Mechanical Properties.

7.1 PHYSICAL PROPERTIES

The various physical properties of the materials are briefly explained below:

- **1. Density:** It is defined as mass of a material per unit volume. Kilogram is the unit of mass in SI system. It is found by weighing the material without voids.
- **2. Bulk density or unit weight:** It is defined as weight of material per unit volume. Its unit in SI is kN/m3 . Volume of the material is found from the natural state of material *i.e.* from volume of voids and solids together.
- **3. Specific gravity:** It is the ratio of density of materials to the unit weight of water. It may be noted that for water density and unit weight are one and the same. Specific gravity is a dimensionless quantity.
- **4. Porosity:** It is the ratio of the volume of voids to the volume of granular materials.
- **5. Hygroscopy:** It is the property of the material to absorb water from the atmosphere. Timber possesses this property.
- **6. Water absorption:** It is the porperty of the material to absorb water and retain it. It is expressed as percentage in weight of absorbed water to weight of dry material. Concrete and bricks possess this property to significant extent.
- **7. Permeability:** It is the property of material by virtue of which it allow water to flow through it. Soil possesses this property and materials like bitumen and metals are impervious.
- **8. Weather resistance:** It is the ability to sustain alternating changes in wet and dry conditions of weather. We look for materials with good weather resistance without undergoing considerable changes in dimensions and loosing mechanical strength.
- **9. Frost resistance:** It is the property to resist repeated freezing and throwing of absorbed water without losing mechanical strength.

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- **10. Thermal conductivity:** It is the ability of the material to transfer heat. Metals have higher thermal conductivity. Moist materials have higher heat conductivity compared to dry porous materials. Rubber is having very low thermal conductivity.
- **11. Thermal resistivity:** It is the reciprocal of thermal conductivity and is defined as time taken for the flow of unit heat.
- **12. Fire resistance:** It is the property by virtue of which a material resists the action of high temperature without undergoing substantial changes in shape and loss of strength. Steel has poor fire resistance. Concrete is better than steel in resisting fire. Bricks are having very good fire resistance.
- **13. Sound absorption:** It is the property of reducing the reflection of sound waves. Porous materials have better sound absorption property. In auditoriums and cinema halls we have to look for materials with good sound absorption properties.
- **14. Chemical resistance:** It is the ability of the material to withstand the action of acids, alkalies, gases etc. Materials are subjected to the action of these chemicals, which are present in air or water or land. Stone and wood have poor resistance to chemical actions while bricks have good chemical resistance.
- **15. Corrosion:** It is the destruction of the material due to slow oxidation. Steel is susceptible to corrosion. The corrosion rate is high in marine environment.
- **16. Soundness:** A material is said to be sound, if it has good resistance to heat, alternate freezing and thawing and the other destructive actions of the atmosphere.
- **17. Durability:** It is the ability to resist the combined effects of atmosphere, rain and other effects and maintain the original strength characters for a long period.

7.2 MECHANICAL PROPERTIES

Resistance to various applied forces is referred as mechanical properties. Some of these properties are briefly presented below:

1. Properties related to axial loading: For designing structural elements to resist the applied loads, the following strength properties should be known:

(*a*) **Tensile strength:** If an element is subjected to pulling forces, the resistance developed per unit area is termed as tensile stress and the maximum trensile stress the material can resist is termed as tensile strength. To find the resisting forces developed a section plane may be taken through the element and the equilibrium of any one part may be considered. Each part is in equilibrium under the action of applied forces and internal resisting forces. Fig. 7.1 From the two sections of the interpretation of the matrix and the section of the flow of unit heat.

2. First estation the property by interest of which a matterial desist the action

for the flow of unit metallity of the

shows a bar under tension and stresses at a cross-section and also an element of bar under tensile stress. In this *P* is the tensile force applied, *p* is the tensile stress (resistance per unit area) and *A* is the cross-sectional area.

∴ $P = pA$...(7.1)

(*b*) **Compressive stress:** Instead of pull, if push '*P*' acts on the element, it tries to shorten the bar and the internal resistance developed per unit area is called compressive stress

(Ref. Fig. 7.2). Here also

2. Properties related to shear loading: If the applied force is trying to shear off a particular section of the element, the resistance developed for unit area in such case along that section is called shearing stress. The maximum stress that material can resist is called shearing strength of the material. Figure 7.3 shows a section of bolt under shearing stress.

If *Q* is the shear force, *q* is the shear stress, then the equilibrium condition of any one part considered shows,

 $Q = qA$, where *A* is the cross-sectional area of the bolt.

$$
q = \frac{Q}{A} \tag{7.2}
$$

3. Properties related to torsional moment: A member is said to be in torsion when it is subjected to a moment about its axis, Fig. 7.5 shows a shaft in torsion. The effect of a torsional member is to twist it and hence a torsional moment is also called as a twisting moment. In engineering problems many members are subjected to torsion. Shafts transmitting power from engine to rear axle of automobile, from a motor to machine tool and from a turbine to electric motor are the common examples of members in torsion. Ring beams of circular water tanks and beams of grid floors are also the examples of members in torsion. The applied torsion is resisted by the material of the member by developing shear stresses. This shear stress varies linearly from zero at centre of the shaft to the maximum value at the extreme fibre.

4. Properties related to bending: When a member is supported at two or more points and subjected to transverse load it bends and develops resistance to the load. The cross-sections of the members are subjected to bending moment and shear force. Finally the load gets transferred to the support by end shear. The shear force introduces shear stresses in the material while bending moment introduces tension in some parts and compression in other parts as shown in Fig. 7.5.

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5. Fatigue: Many structural memberes are subjected to fluctuation and reversal of stresses as shown in Fig. 7.6 (*a*) and (*b*). A bridge deck is a common example of fluctuation of stresses and the shaft is the mechanical component subjected to reversal of stresses. This type of stress when applied repeatedly reduces the strength of material and this effect is called fatigue. This property is shown in Fig. 7.6 (*c*). The maximum stress at which even a billion reversal of stresses cannot cause failure of the material is called endurance limit.

6. Abrasive resistance: This is the property by virtue of which a material resists the forces acting at contact surfaces when one material rubs/moves over the other. This is an important property to be studied for the materials used for road surfaces and flooring.

7. Impact strength: This property refers to the ability of the material to resist shock loads due to heavy loads falling on the surface.

8. Hardness: It is the ability of the material to resist penetration from another material. A number called hardness number is used to measure hardness of various materials. It is based on hardnesses of ten minerals arranged in the increasing hardness. Diamond has hardness number 10. Sometimes hardness is measured by size of indentation of steel balls under standard pressure using hardness testing machine.

9. Elasticity: It is the property of the material by virtue of which it regains its original shape and size after the removal of external load. The maximum stress level before which if the load is removed the material regains its shape and size fully is called its elastic limit.

10. Plasticity: It is the property of the material to retain its changed shape and size after the loads are removed. It is a required property when a material is to be moulded into different shape.

11. Creep: It is the property of the material by virtue of which it undergoes changes in size with time under the action of constant load. Concrete possesses this property.

12. Toughness: It is the property of a material whereby it absorbs energy due to straining actions by undergoing plastic deformation.

The other terms proof stress, factor of safety, working stress and load factors are explained in the next chapter after explaining the stress strain curves. There are standard tests specified by codes to quantity various properties of the materials. Knowing the requirements of materials properties for different purposes, the engineer has to select or reject the materials to be used in construction. is a manuary a to see property one manuar is been also be mounded into different the temporal of the material is to be mounded into different intercept to the property of the material is to be mounded into different interc

THEORY QUESTIONS

1. Explain the terms:

Weather resistance, frost resistance, thermal resistance and fire resistance.

- **2.** What do you understand by the terms: sound absorption, soundness, corrosion resistance and durability ?
- **3.** Differentiate between:
	- (*a*) Impact strength and toughness (*b*) Elasticity and plasticity

(*c*) Tensile stress and shear stress.

8

Simple Stresses and Strains

In this chapter general meaning of stress is explained. Expressions for stresses and strains is derived with the following assumptions:

- 1. For the range of forces applied the material is elastic *i.e.* it can regain its original shape and size, if the applied force is removed.
- 2. Material is homogeneous *i.e.* every particle of the material possesses identical mechanical properties.
- 3. Material is isotropic *i.e.* the material possesses identical mechanical property at any point in any direction.

Presenting the typical stress-strain curve for a typical steel, the commonly referred terms like limits of elasticity and proportionality, yield points, ultimate strength and strain hardening are explained.

Linear elastic theory is developed to analyse different types of members subject to axial, shear, thermal and hoop stresses.

8.1 MEANING OF STRESS

When a member is subjected to loads it develops resisting forces. To find the resisting forces developed a section plane may be passed through the member and equilibrium of any one part may be considered. Each part is in equilibrium under the action of applied forces and internal resisting forces. The resisting forces may be conveniently split into normal and parallel to the section plane. The resisting force parallel to the plane is called *shearing resistance*. The intensity of resisting force normal to the sectional plane is called *intensity of Normal Stress* (Ref. Fig. 8.1).

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In practice, intensity of stress is called as ''stress'' only. Mathematically

Normal Stress =
$$
p = \lim_{\Delta A \to 0} \frac{\Delta R}{\Delta A}
$$

= $\frac{dR}{dA}$...(8.1)

where R is normal resisting force.

The intensity of resisting force parallel to the sectional plane is called *Shearing Stress* (*q*).

Shearing Stress =
$$
q = \lim_{\Delta A \to 0} \frac{\Delta Q}{\Delta A} = \frac{dQ}{dA}
$$
 ...(8.2)

where *Q* is Shearing Resistance.

Thus, *stress at any point may be defined as resistance developed per unit area.* From equations (8.1) and (8.2) , it follows that

$$
dR = pdA
$$
\n
$$
R = \int pdA
$$
\n...(8.3*a*)\n
$$
Q = \int qdA
$$
\n...(8.3*b*)

At any cross-section, stress developed may or may not be uniform. In a bar of uniform crosssection subject to axial concentrated loads as shown in Fig. 8.2*a*, the stress is uniform at a section away from the applied loads (Fig. 8.2*b*); but there is variation of stress at the section near the applied loads (Fig. 8.2*c*).

Fig. 8.2

Similarly stress near the hole or at fillets will not be uniform as shown in Figs. 8.3 and 8.4. It is very common that at some points in such regions maximum stress will be as high as 2 to 4 times the average stresses.

Fig. 8.3. Stresses in a Plate with a Hole **Fig. 8.4**

8.2 UNIT OF STRESS

When Newton is taken as unit of force and millimetre as unit of area, unit of stress will be N/mm². The other derived units used in practice are kN/mm², N/m², kN/m² or MN/m². A stress of one $N/m²$ is known as Pascal and is represented by Pa.

Hence, 1 MPa = 1 MN/m² = 1 × 10⁶ N/(1000 mm)² = 1 N/mm².

Thus one Mega Pascal is equal to 1 N/mm². In most of the standard codes published unit of stress has been used as Mega Pascal (MPa or N/mm^2).

8.3 AXIAL STRESS

Consider a bar subjected to force *P* as shown in Fig. 8.5. To maintain the equilibrium the end forces applied must be the same, say *P*.

(b) Resisting Force Developed

Fig. 8.5. Tensile Stresses

The resisting forces acting on a section are shown in Fig. 8.5*b*. Now since the stresses are uniform

$$
R = \int pdA = p \int dA = pA \qquad ...(8.4)
$$

where *A* is the cross-sectional area.

Considering the equilibrium of a cut piece of the bar, we get $P = R$ (8.5) From equations (8.4) and (8.5), we get $P = pA$ or $p = \frac{P}{A}$ $\frac{1}{A}$...(8.6)

Thus, in case of axial load 'P' the stress developed is equal to the load per unit area. Under this type of normal stresses the bar is being extended. Such stress which is causing extension of the bar is called tensile stress.

A bar subjected to two equal forces pushing the bar is shown in Fig. 8.6. It causes shortening of the bar. *Such forces which are causing shortening, are known as compressive forces and corresponding stresses as compressive stresses.*

Now $R = \int pdA = p \int dA$ (as stress is assumed uniform) For equilibrium of the piece of the bar

$$
P = R = pA
$$

or

$$
p = \frac{P}{A}
$$
 as in equation 8.6

Thus, whether it is tensile or compressive, the stress developed in a bar subjected to axial forces, is equal to load per unit area.

8.4 STRAIN

No material is perfectly rigid. Under the action of forces a rubber undergoes changes in shape and size. This phenomenon is very well known to all since in case of rubber, even for small forces deformations are quite large. Actually all materials including steel, cast iron, brass, concrete, etc. undergo similar deformation when loaded. But the deformations are very small and hence we cannot see them with naked eye. There are instruments like extensometer, electric strain gauges which can measure extension of magnitude 1/100th, 1/1000th of a millimetre. There are machines like universal testing machines in which bars of different materials can be subjected to accurately known forces of magnitude as high as 1000 kN. The studies have shown that the bars extend under tensile force and shorten under compressive forces as shown in Fig. 8.7. *The change in length per unit length is known as linear strain.* Thus, **E** = pd

Thus, in case of actal loved $\frac{P}{P}$

Thus in case of actal loved $\frac{P}{P}$ the stress developed is equal to the load give unit of

the pre of normal arreases the art is being extended. Such stress which is ca

Linear Strain $=$ $\frac{\text{Change in Length}}{\text{Original Length}}$

When changes in longitudinal direction is taking place changes in lateral direction also take place. The nature of these changes in lateral direction are exactly opposite to that of changes in longitudinal direction *i.e.*, if extension is taking place in longitudinal direction, the shortening of lateral dimension takes place and if shortening is taking place in longitudinal direction extension takes place in lateral directions (See Fig. 8.7). *The lateral strain may be defined as changes in the lateral dimension per unit lateral dimension.* Thus,

Lateral Strain =
$$
\frac{\text{Change in Lateral Dimension}}{\text{Original Lateral Dimension}}
$$

$$
= \frac{b' - b}{b} = \frac{\delta b}{b}
$$
...(8.8)

8.5 STRESS-STRAIN RELATION

The stress-strain relation of any material is obtained by conducting tension test in the laboratories on standard specimen. Different materials behave differently and their behaviour in tension and in compression differ slightly.

8.5.1 Behaviour in Tension

Mild steel. Figure 8.8 shows a typical tensile test specimen of mild steel. Its ends are gripped into universal testing machine. Extensometer is fitted to test specimen which measures extension over the length *L*1, shown in Fig. 8.8. The length over which extension is mesured is called *gauge length.* The load is applied gradually and at regular interval of loads extension is measured. After certain load, extension increases at faster rate and the capacity of extensometer to measure extension comes to an end and, hence, it is removed before this stage is reached and extension is measured from scale on the universal testing machine. Load is increased gradually till the specimen breaks.

Load divided by original cross-sectional area is called as nominal stress or simply as stress. Strain is obtained by dividing extensometer readings by gauge length of extensometer (L_1) and by dividing scale readings by grip to grip length of the specimen (L_2) . Figure 8.10 shows stress *vs* strain diagram for the typical mild steel specimen. The following salient points are observed on stress-strain curve:

- (*a*) *Limit of Proportionality (A): It is the limiting value of the stress up to which stress is proportional to strain.*
- (*b*) *Elastic Limit: This is the limiting value of stress up to which if the material is*

stressed and then released (unloaded) strain disappears completely and the original length is regained. This point is slightly beyond the limit of proportionality.

- (*c*) *Upper Yield Point (B): This is the stress at which, the load starts reducing and the extension increases.* This phenomenon is called yielding of material. At this stage strain is about 0.125 per cent and stress is about 250 N/mm².
- (*d*) *Lower Yield Point (C): At this stage the stress remains same but strain increases for some time.*
- (*e*) *Ultimate Stress (D):* This is the maximum stress the material can resist. This stress is about 370–400 N/mm². At this stage cross-sectional area at a particular section starts reducing very fast (Fig. 8.9). This is called neck formation. After this stage load resisted and hence the stress developed starts reducing.
- (*f*) *Breaking Point (E): The stress at which finally the specimen fails is called breaking point.* At this strain is 20 to 25 per cent.

If unloading is made within elastic limit the original length is regained *i.e.*, the stress-strain curve

Fig. 8.8. Tension Test Specimen **Fig. 8.9.** Tension Test Specimen after Breaking

follows down the loading curve shown in Fig. 8.6. If unloading is made after loading the specimen beyond elastic limit, it follows a straight line parallel to the original straight portion as shown by line *FF*′ in Fig. 8.10. Thus if it is loaded beyond elastic limit and then unloaded a permanent strain (*OF*) is left in the specimen. This is called *permanent* set.

Stress-strain relation in aluminium and high strength steel. In these elastic materials there is no clear cut yield point. The necking takes place at ultimate stress and eventually the breaking point is lower than the ultimate point. The typical stress-strain diagram is shown in Fig. 8.11. The stress *p* at which if unloading is made there will be 0.2 per cent permanent set is known as 0.2 per cent proof stress and this point is treated as yield point for all practical purposes.

Aluminium and High Strength Steel for Brittle Material

Stress-strain relation in brittle material. The typical stress-strain relation in a brittle material like cast iron, is shown in Fig. 8.12.

In these material, there is no appreciable change in rate of strain. There is no yield point and no necking takes place. Ultimate point and breaking point are one and the same. The strain at failure is very small.

*Percentage elongation and percentage reduction in area***.** Percentage elongation and percentage reduction in area are the two terms used to measure the ductility of material.

(*a*) *Percentage Elongation: It is defined as the ratio of the final extension at rupture to original length expressed, as percentage*. Thus,

Percentage Elongation =
$$
\frac{L'-L}{L} \times 100
$$
 ...(8.9)

where L – original length, L' – length at rupture.

The code specify that original length is to be five times the diameter and the portion considered must include neck (whenever it occurs). Usually marking are made on tension rod at every '2.5 *d*' distance and after failure the portion in which necking takes place is considered. In case of ductile material percentage elongation is 20 to 25.

(*b*) *Percentage Reduction in Area: It is defined as the ratio of maximum changes in the crosssectional area to original cross-sectional area, expressed as percentage.* Thus,

Percentage Reduction in Area =
$$
\frac{A - A'}{A} \times 100
$$
 ...(8.10)

where *A*–original cross-sectional area, *A*′–minimum cross-sectional area. In case of ductile material, *A*′ is calculated after measuring the diameter at the neck. For this, the two broken pieces of the specimen are to be kept joining each other properly. For steel, the percentage reduction in area is 60 to 70.

8.5.2 Behaviour of Materials under Compression

As there is chance to bucking (laterally bending) of long specimen, for compression tests short specimens are used. Hence, this test involves measurement of smaller changes in length. It results into lesser accuracy. However precise measurements have shown the following results:

- (*a*) In case of ductile materials stress-strain curve follows exactly same path as in tensile test up to and even slightly beyond yield point. For larger values the curves diverge. There will not be necking in case of compression tests.
- (*b*) For most brittle materials ultimate compresive stress in compression is much larger than in tension. It is because of flows and cracks present in brittle materials which weaken the material in tension but will not affect the strength in compression.

8.6 NOMINAL STRESS AND TRUE STRESS

So far our discussion on direct stress is based on the value obtained by dividing the load by original cross-sectional area. That is the reason why the value of stress started dropping after neck is formed in mild steel (or any ductile material) as seen in Fig. 8.10. But actually as material is stressed its cross-sectional area changes. We should divide load by the actual cross-sectional area to get true stress in the material. To distinguish between the two values we introduce the terms nominal stress and true stress and define them as given below:

Fig. 8.13. Nominal Stress-Strain Curve and True Stress-Strain Curve for Mild Steel.
8.7 FACTOR OF SAFETY

In practice it is not possible to design a mechanical component or structural component permitting stressing up to ultimate stress for the following reasons:

- 1. Reliability of material may not be 100 per cent. There may be small spots of flaws.
- 2. The resulting deformation may obstruct the functional performance of the component.
- 3. The loads taken by designer are only estimated loads. Occasionally there can be overloading. Unexpected impact and temperature loadings may act in the lifetime of the member.
- 4. There are certain ideal conditions assumed in the analysis (like boundary conditions). Actually ideal conditions will not be available and, therefore, the calculated stresses will not be 100 per cent real stresses.

Hence, *the maximum stress to which any member is designed is much less than the ultimate stress, and this stress is called Working Stress. The ratio of ultimate stress to working stress is called factor of safety*. Thus

Factor of Safety =
$$
\frac{\text{Ultimate Stress}}{\text{Working Stress}}
$$
...(8.12)

In case of elastic materials, since excessive deformation create problems in the performance of the member, working stress is taken as a factor of yield stress or that of a 0.2 proof stress (if yield point do not exist).

Factor of safety for various materials depends up on their reliability. The following values are commonly taken in practice:

- 1. For steel 1.85
- 2. For concrete 3
- 3. For timber -4 to 6

8.8 HOOKE'S LAW

Robert Hooke, an English mathematician conducted several experiments and concluded that *stress is proportional to strain up to elastic limit.* This is called Hooke's law. Thus Hooke's law is, up to elastic limit

$$
\propto e
$$

p ∝ *e* ...(8.13*a*)

where *p* is stress and *e* is strain

Hence, *p* = *Ee* ...(8.13*b*)

where E is the constant of proportionality of the material, known as modulus of elasticity or Young's modulus, named after the English scientist Thomas Young (1773–1829).

However, present day sophisticated experiments have shown that for mild steel the Hooke's law holds good up to the proportionality limit which is very close to the elastic limit. For other materials, as seen in art. 1.5, Hooke's law does not hold good. However, in the range of working stresses, assuming Hooke's law to hold good, the relationship does not deviate considerably from actual behaviour. Accepting Hooke's law to hold good, simplifies the analysis and design procedure considerably. Hence Hooke's law is widely accepted. The analysis procedure accepting Hooke's law is known as Linear Analysis and the design procedure is known as the working stress method. 1. Reliability of material may note by 100 per cent. These may be small spots of the complex and spots the complex and super-
and spectra control and spectra of the complex of the back taken by designer are only estimated

8.9 EXTENSION/SHORTENING OF A BAR

Consider the bars shown in Fig. 8.14

From equation (8.6), Stress $p = \frac{P}{P}$ *A*

From equation (8.7), Strain, $e = \frac{\Delta}{L}$

From Hooke's Law we have,

$$
E = \frac{\text{Stress}}{\text{Strain}} = \frac{p}{e} = \frac{P/A}{\Delta/L} = \frac{PL}{A\Delta}
$$

or

$$
\Delta = \frac{PL}{AE}
$$
...(8.14)

 $Solution:$

*Example 8.1***.** *A circular rod of diameter 16 mm and 500 mm long is subjected to a tensile force 40 kN. The modulus of elasticity for steel may be taken as 200 kN/mm² . Find stress, strain and elongation of the bar due to applied load.*

Load
$$
P = 40 \text{ kN} = 40 \times 1000 \text{ N}
$$

\n $E = 200 \text{ kN/mm}^2 = 200 \times 10^3 \text{ N/mm}^2$
\n $L = 500 \text{ mm}$

Diameter of the rod $d = 16$ mm

 $Therefore, sectional area$ $A = \frac{\pi d^2}{4} = \frac{\pi}{4} \times 16^2$ $= 201.06$ mm² **Stress** $p = \frac{P}{A} = \frac{40 \times 1000}{201.06} = 198.94 \text{ N/mm}^2$ **Strain** $e = \frac{p}{E} = \frac{198.94}{200 \times 10^{10}}$ 200×10^{3} $\frac{.94}{10^{3}}$ = 0.0009947 Elongation $\Delta = \frac{PL}{AE} = \frac{4.0 \times 1000 \times 5}{201.06 \times 200 \times 1000}$ $\frac{4.0 \times 1000 \times 500}{201.06 \times 200 \times 10^{3}}$ = **0.497** mm P

Fig. 8.14

From equation (8.6), Stress $p = \frac{p}{A}$

Fig. 8.14

From equation (8.5), Stress, $p = \frac{p}{A}$

From Hocke's Law we have,
 $E = \frac{8 \text{ terms}}{8 \text{ min}} = \frac{p}{e} = \frac{P/A}{\Delta L}$

or

Example 8.1. A circulate red of diameer 1 *Example 8.2***.** *A Surveyor's steel tape 30 m long has a cross-section of 15 mm × 0.75 mm. With this, line AB is measure as 150 m. If the force applied during measurement is 120 N more than the force applied at the time of calibration, what is the actual length of the line?*

Take modulus of elasticity for steel as 200 kN/mm² .

Solution: $A = 15 \times 0.75 = 11.25 \text{ mm}^2$

 $P = 120$ N, $L = 30$ m = 30×1000 mm $E = 200 \text{ kN/mm}^2 = 200 \times 10^3 \text{ N/mm}^2$

Elongation
$$
\Delta = \frac{PL}{AE} = \frac{120 \times 30 \times 1000}{11.25 \times 200 \times 10^3} = 1.600
$$
 mm

Hence, if measured length is 30 m.

Actual length is 30 m + 1.600 mm = 30.001600 m

:. Actual length of line
$$
AB = \frac{150}{30} \times 30.001600 = 150.008
$$
 m

*Example 8.3***.** *A hollow steel tube is to be used to carry an axial compressive load of 160 kN. The yield stress for steel is 250 N/mm² . A factor of safety of 1.75 is to be used in the design. The following three class of tubes of external diameter 101.6 mm are available.*

$$
A = \frac{160 \times 10^3}{142.857} = 1120
$$
 mm²

120

For hollow section of outer diameter '*D*' and inner diameter '*d*'

$$
A = \frac{\pi}{4}(D^2 - d^2) = 1
$$

$$
\frac{\pi}{4}(101.6^2 - d^2) = 1120
$$

π

$$
d^{2} = 8896.53 \quad \therefore \quad d = 94.32 \text{ mm}
$$
\n
$$
t = \frac{D - d}{2} = \frac{101.6 - 94.32}{2} = 3.63 \text{ mm}
$$

Hence, use of light section is recommended.

Example 8.4. A specimen of steel 20 mm diameter with a gauge length of 200 mm is tested to destruction. It has an extension of 0.25 mm under a load of 80 kN and the load at elastic limit is 102 kN. The maximum load is 130 kN. Hence, use of light section is recommended.

Hence, the A. A specimen of steel 20 non diameter with a gauge length of 200 non is

determined by A. The maximum tood is 130 kN mulder a load of 80 kN and the load at elast

1

The total extension at fracture is 56 mm and diameter at neck is 15 mm. Find

- (*i*) *The stress at elastic limit.*
- (*ii*) *Young's modulus.*

(*iii*) *Percentage elongation.*

(*iv*) *Percentage reduction in area.*

(*v*) *Ultimate tensile stress.*

Solution: Diameter
$$
d = 20
$$
 mm

Area
$$
A = \frac{\pi d^2}{4} = 314.16 \text{ mm}^2
$$

(*i*) **Stress at elastic limit** $=$ Load at elastic limit

Area

 $=\frac{102 \times 10^3}{314.16} = 324.675$ N/mm² (*ii*) **Young's modulus** $E = \frac{\text{Stress}}{\text{Stress}}$ within elastic limit

Strain

$$
= \frac{P/A}{\Delta/L} = \frac{80 \times 10^3 / 314.16}{0.25 / 200}
$$

$$
= 203718 \text{ N/mm}^2
$$

 (iii) **Percentage elongation** = $\frac{\text{Final extension}}{\text{Final extension}}$ Original length

$$
=\frac{56}{200} \times 100 = 28
$$

(*iv*) **Percentage reduction in area**

 $=\frac{\text{Initial area} - \text{Final area}}{2} \times 100$ Initial area

$$
= \frac{\frac{\pi}{4} \times 20^2 - \frac{\pi}{4} \times 15^2}{\frac{\pi}{4} \times 20^2} \times 100 = 43.75
$$

(*v*) **Ultimate Tensile Stress** = Ultimate Load Area

$$
= \frac{130 \times 10^3}{314.16} = 413.80 \text{ N/mm}^2.
$$

8.10 BARS WITH CROSS-SECTIONS VARYING IN STEPS

A typical bar with cross-sections varying in steps and subjected to axial load is as shown in Fig. 8.15(*a*). Let the length of three portions be L_1 , L_2 and L_3 and the respective cross-sectional areas of the portion be A_1 , A_2 , A_3 and E be the Young's modulus of the material and P be the applied axial load.

Figure 8.15(*b*) shows the forces acting on the cross-sections of the three portions. It is obvious that to maintain equilibrium the load acting on each portion is *P* only. Hence stress, strain and extension of each of these portions are as listed below:

Fig. 8.15. Typical Bar with Cross-section Varying in Step

Hence total change in length of the bar

$$
\Delta = \Delta_1 + \Delta_2 + \Delta_3 = \frac{PL_1}{A_1E} + \frac{PL_2}{A_2E} + \frac{PL_3}{A_3E} \tag{8.15}
$$

Example 8.5. The bar shown in Fig. 8.16 is tested in universal testing machine. It is observed that at a load of 40 kN the total extension of the bar is 0.280 mm. Determine the Young's modulus of the material.

Example 8.6. The stepped bar shown in Fig. 8.17 is made up of two different materials. The material 1 has Young's modulus = 2×10^5 *N/mm, while that of material* 2 is 1 \times 10⁵ *N/mm*². Find the *extension of the bar under a pull of 30 kN if both the portions are 20 mm in thickness.*

8.11 BARS WITH CONTINUOUSLY VARYING CROSS-SECTIONS

When the cross-section varies continuously, an elemental length of the bar should be considered and general expression for elongation of the elemental length derived. Then the general expression should be integrated over entire length to get total extension. The procedure is illustrated with Examples 8.8 and 8.9.

Example 8.8. A bar of uniform thickness 't' tapers uniformly from a width of $b₁$ at one end to $b₂$ *at other end in a length 'L' as shown in Fig. 8.18. Find the expression for the change in length of the bar when subjected to an axial force P.*

Fig. 8.19

*Solution***:** Consider an elemental length *dx* at a distance *x* from larger end. Rate of change of breadth is $\frac{b_1 - b}{a_1}$ $a_1 - b_2$.

$$
L
$$

Hence, width at section *x* is $b = b_1 - \frac{b_1 - b_2}{L} x = b_1 - kx$

where
$$
k = \frac{b_1 - b_2}{L}
$$

 \therefore Cross-section area of the element = $A = t(b_1 - kx)$ Since force acting at all sections is *P* only,

Extension of element $= \frac{Pdx}{AE}$ [where length = dx] $=\frac{Pdx}{\sqrt{2\pi}}$ $(b_1 - kx)tE$ **Total extension of the bar** $= \int_0^L \frac{P dx}{(b_1 - kx)tE} = \frac{P}{tE} \int_0^L \frac{1}{(b_1 - kx)^2} dt$ $b_1 - kx$)tE *P tE dx* $\int_0^L \frac{P dx}{(b_1 - kx)tE} = \frac{P}{tE} \int_0^L \frac{dx}{(b_1 - kx)}$ $=\frac{P}{tE}\left(\frac{1}{-k}\right)$ $b_1 - kx$ $\left[\begin{matrix}1\end{matrix}\right]_{1\leq i\leq k}$ $\left(\frac{1}{-k}\right)\left[\log\left(b_{1}-kx\right)\right]_{0}^{k}$ \overline{a} $\overline{\mathcal{C}}$ $\overline{}$ $\left[\log \left(b_1 - kx \right) \right]$ $\overline{\mathsf{L}}$ $\overline{}$ $=\frac{P}{P}$ *tEk* $b_1 - \frac{b_1 - b_2}{a_1 - b_2}$ *L x* $-\log\left(b_1 - \frac{b_1 - b_2}{L}x\right)\Big]_0^L$ $\frac{1}{2}$ $\overline{\mathbb{I}}$ $\overline{\mathsf{L}}$ $\left[-\log\left(b_1-\frac{b_1-b_2}{L}x\right)\right]$ $\overline{}$ 0 $=\frac{P}{tEk}[-\log b_2 + \log b_1] = \frac{P}{tEk} \log \frac{b}{b}$ 1 2 $=\frac{PL}{tE(b_1-b_2)}\log\frac{b}{b}$ **1 .** ...(8.16) Solution: Consider an elemental length dx at a distance x from larger end. Rate of change

Solution: Consider an elemental length dx at a distance x from larger end. Rate of change

is $\frac{b_1-b_2}{L}$.

Hence, width at sec

 $-{\bf b}_2$ ^{\cdots} ${\bf b}_2$

Example 8.9. A tapering rod has diameter d_1 at one end and it tapers uniformly to a diameter d_2 *at the other end in a length L as shown in Fig. 8.20. If modulus of elasticity of the material is E, find its change in length when subjected to an axial force P.*

Solution: Change in diameter in length *L* is $d_1 - d_2$

∴ Rate of change of diameter, $k = \frac{d_1 - d_2}{d_1}$ *L* $a_1 - d_2$

Consider an elemental length of bar *dx* at a distance *x* from larger end. The diameter of the bar at this section is

$$
d = d_1 - kx.
$$

Cross-sectional area

$$
A = \frac{\pi d^2}{4} = \frac{\pi}{4} (d_1 - kx)^2
$$

∴ Extension of the element

$$
= \frac{P dx}{\frac{\pi}{4} (d_1 - kx)^2 E}
$$

Extension of the entire bar

Solution: Change in diameter in length L is
$$
d_1 - d_2
$$

\n
$$
\therefore
$$
 Rate of change of diameter, $k = \frac{d_1 - d_2}{L}$
\nConsider an elemental length of bar dx at a distance x from larger end. The diameter of t at this section is
\n
$$
d = d_1 - kx.
$$

\nCross-sectional area
\n
$$
\therefore
$$
 Extension of the element
\n
$$
= \frac{\pi d^2}{4} = \frac{\pi}{4}(d_1 - kx)^2
$$

\nExtension of the entire bar
\n
$$
\Delta = \int_0^L \frac{P dx}{\frac{\pi}{4}(d_1 - kx)^2 E}
$$

\n
$$
= \frac{4P}{\pi E} \int_0^L \frac{dx}{(d_1 - kx)^2}
$$

\n
$$
= \frac{4P}{\pi E} \left(\frac{1}{d_1 - kx}\right)_0^L
$$

\n
$$
= \frac{4P}{\pi E(d_1 - d_2)} \left(\frac{1}{d_2} - \frac{1}{d_1}\right), \text{ since } d_1 - kL = d_2
$$

\n
$$
\Delta = \frac{4PL}{\pi E(d_1 - d_2)} \times \frac{(d_1 - d_2)}{d_1 d_2} = \frac{4PL}{\pi E(d_1 d_2)}.
$$

 $_{1} - u_{2}$ $1^{\mathcal{U}}2$

 $1 - u_2$

 $\pi E (d_1 - d_2)$ *d*₁*d*

 $E(d_1 - d$

$$
\mathcal{L}^{\mathcal{L}}_{\mathcal{L}}
$$

. ...(8.17)

Note: For bar of uniform diameter extension is
$$
\frac{PL}{\pi d^2 E}
$$
 and for the tangenting rod it is $\frac{PL}{\pi d_1 d_2 E}$.

Solution: Now, $t = 10 \text{ mm}$ $b_1 = 60 \text{ mm}$ $b_2 = 40 \text{ mm}$

Thus the change is, instead of d^2 for uniform bar, $d_1 d_2$ term appears if the bar tapers uniformly. *Example 8.10***.** *A steel flat of thickness 10 mm tapers uniformly from 60 mm at one end to 40 mm at other end in a length of 600 mm. If the bar is subjected to a load of 80 kN, find its extension.* Take $E = 2 \times 10^5$ MPa. What is the percentage error if average area is used for calculating *extension?*

 $L = 600$ mm $P = 80$ kN = 80000 N Now, $1 \text{ MPa} = 1 \text{ N/mm}^2$ Hence $E = 2 \times 10^5$ N/mm² Extension of the tapering bar of rectangular section $\Delta = \frac{PL}{tE(b_1 - b_2)} \log \frac{b}{b}$ 1 2 $=\frac{80000\times600}{10000\times600}$ $10 \times 2 \times 10^5$ (60 – 40 $\frac{80000\times600}{\times 2\times10^5(60-40)}$ log $\frac{60}{40}$ = **0.4865 mm** If averages cross-section is considered instead of tapering cross-section, extension is given by $\Delta = \frac{PL}{I}$ $A_{av}E$ Now $A_{av} = \frac{60 \times 10 + 40 \times 10}{2}$ $\frac{\times 10 + 40 \times 10}{0}$ = 500 mm² $\Delta = \frac{80000 \times 600}{25}$ $500 \times 2 \times 10^5$ × \times 2 \times $= 0.480$ mm \therefore **Percentage error** = $\frac{0.4865 - 0}{0.4865}$ $.4865 - 0.48$ $\frac{65-0.48}{0.4865} \times 100$ Thus the change is, instead of d^2 for uniform bar, d_1 d₂ term appears if the bartagers
 Circle Batter and the maximal fieldness. *10 mm appears infforms by tom 60 mm at anne end*
 Circle E = 2 × 10 MPu. What is

= **1.348**

*Example 8.11***.** *A 2.0 m long steel bar is having uniform diameter of 40 mm for a length of 1 m and in the next 0.5 m its diameter gradually reduces from 40 mm to 20 mm as shown in Fig. 8.21. Determine the elongation of this bar when subjected to an axial tensile load of 200 kN. Given* $E = 200 \text{ GN/m}^2$.

Solution: Now, P

$$
P = 200 \times 10^3 \text{ N}
$$

\n
$$
E = 200 \text{ GN/m}^2 = \frac{200 \times 10^9}{(1000)^2} \text{ N/mm}^2
$$

\n
$$
= 200 \times 1000 \text{ N/mm}^2
$$

\n
$$
= 2 \times 10^5 \text{ N/mm}^2
$$

Extensions of uniform portion and tapering portion are worked out separately and then added to get extension of the given bar.

Extension of uniform portion

$$
\Delta_1 = \frac{PL}{AE} = \frac{200 \times 10^3 \times 1500}{\frac{\pi}{4} \times 40^2 \times 2 \times 10^5} = 1.194 \text{ mm}.
$$

Extension of tapering portion

$$
\Delta_2 = \frac{4 \, PL}{E \pi d_1 d_2} = \frac{4 \times 200 \times 10^3 \times 500}{2 \times 10^5 \times \pi \times 60 \times 40}
$$

= 0.265 mm
= $\Delta_1 + \Delta_2 = 1.194 \pm 0.265 = 1.459$ m

 $Total extension$

$$
= \Delta_1 + \Delta_2 = 1.194 + 0.265 = 1.459 \text{ mm}
$$

Example 8.12. The extension of a bar uniformly tapering from a diameter of d + a to d – a in a length L is calculated by treating it as a bar of uniform cross-section of average diameter d. What is the percentage error? = 200 × 1000 N/mm²

Extensions of uniform pertion and taperals portion are worked out separately and then

get extension of the given bar.

Extension of uniform pertion
 $\Delta_1 = \frac{PL}{Ak} = \frac{200 \times 10^3 \times 1500}{\pi \times 40^2 \times 2 \$

*Solution***:** Actual extension under load

$$
P = \frac{4PL}{\pi d_1 d_2 E} = \frac{4PL}{\pi (d+a)(d-a)E}
$$

$$
= \frac{4PL}{\pi E (d^2 - a^2)}
$$

If it is treated as a bar of uniform diameter '*d*', erroneous extension calculated

$$
= \frac{PL}{\frac{\pi}{4} \times d^2 E} = \frac{4 PL}{\pi d^2 E}
$$

∴ **Percentage error** =

$$
= \frac{\frac{4\,PL}{\pi E(d^2 - a^2)} - \frac{4\,PL}{\pi Ed^2}}{\frac{4\,PL}{\pi E(d^2 - a^2)}} \times 100
$$

$$
= \frac{\frac{1}{(d^2 - a^2)} - \frac{1}{d^2}}{\frac{1}{(d^2 - a^2)}} \times 100
$$

=

Thus, the direct stress is equal to shearing force per unit area.

8.13 SIMPLE SHEAR

A material is said to be in a state of simple shear if it is subjected to only shearing stress. A shearing force tries to shear off the cross-section of the body. Consider a bolt subjected to pure shear as shown in Fig. 8.24*a*. A rectangular element at this section is shown in Fig. 8.24*b*.

Let the intensity of shear stress be q_{ab} and thickness of element be '*t*'. Consider the equilibrium of the element.

Vertical force on $AB = q_{ab} \times AB \times t$

This can be balanced by vertical downward force on *CD*.

(a) Bolt Subjected to Shearing

qab

A

q

 q_{bc} B

C

D q_{cc}

 α

 \overline{q}_{π}

(b) Element under Shear

(c) State of Simple Shear **Fig. 8.24**

$q_{ab} \times AB \times t = q_{cd} \times CD \times t$

 $q_{ab} = q_{cd} = q$ (say)

Since q_{ab} and q_{cd} are giving rise to equal and opposite forces $q \times AB \times t$, with moment arm = *AD*, they form a couple of magnitude equal to

$$
q \times AB \times t \times AD \qquad \qquad ...(1)
$$

This can be balanced by another couple only. *i.e.* q_{bc} and q_{cd} should form an equal and opposite couple. Hence their direction will be as shown in Fig. 8.24*b*.

Let $q_{bc} = q_{ad} = q'$

The couple formed by these forces

 $= q' \times AD \times t \times AB$...(2)

Equating (1) and (2) we get, $q = q'$.

Thus if a section is subjected to pure shear, the state of stress in any element at that section is as shown in Fig. 8.24*c*.

State of stress at 45° to shearing direction is of engineering interest. To study it, for simplicity, consider a square element of side '*a*' and thickness '*t*', under simple shear as shown in Fig. 8.25*a*.

Since it is square element $AC = a\sqrt{2}$. Consider section along AC and let 'p' be the stress on this section. From equilibrium condition of system of forces in the direction normal to *AC*, we get

$$
p \times AC \times t = q CD t \cos 45^\circ + q AD t \sin 45^\circ
$$

$$
pa\sqrt{2}t = qat \frac{1}{\sqrt{2}} + qat \frac{1}{\sqrt{2}}
$$

$$
= qat\sqrt{2}
$$

i.e., p = *q*

Thus in case of simple shear tensile stress of the same magnitude as shearing stress develops at 45° to shearing plane. By taking stresses on diagonal *BD* it can be shown that compressive stress of magnitude *q* acts on this plane. *Thus simple shear gives rise to tensile and compressive stresses across planes inclined at 45° to the shearing planes, the intensity of direct stresses being of same magnitude as shearing stresss*.

Shearing strain. Shearing stress has a tendency to distort the element to position *AB*′*C*′*D* from the original position *ABCD* as shown in Fig. 8.26.

This deformation is expressed in terms of angular displacement and is called shear strain. Thus

Shear strain =
$$
\frac{BB'}{AB}
$$
 = tan ϕ

 $= \phi$, since angle ϕ is small

8.14 POISSON'S RATIO

When a material undergoes changes in length, it undergoes changes of opposite nature in lateral directions. For example, if a bar is subjected to direct tension in its axial direction it elongates and at the same time its sides contract (Fig. 8.27).

Fig. 8.27. Changes in Axial and Lateral Directions

If we define the ratio of change in axial direction to original length as linear strain and change in lateral direction to the original lateral dimension as lateral strain, it is found that *within elastic limit there is a constant ratio between lateral strain and linear strain. This constant ratio is called Poisson's ratio.* Thus,

$$
Poisson's ratio = \frac{\text{Lateral strain}}{\text{Linear strain}} \qquad \qquad \dots (8.19)
$$

^V ...(8.20)

It is denoted by $\frac{1}{1}$ *m* , or µ. For most of metals its value is between 0.25 to 0.33. Its value for steel is 0.3 and for concrete 0.15.

8.15 VOLUMETRIC STRAIN

When a member is subjected to stresses, it undergoes deformation in all directions. Hence, there will be change in volume. The *ratio of the change in volume to original volume is called volumetric strain.*

 $e_v = \frac{\delta V}{V}$

where e_V = Volumetric strain

 δ_V = Change in volume

 $V =$ Original volume

It can be shown that volumetric strain is sum of strains in three mutually perpendicular directions.

i.e., $e_y = e_x + e_y + e_z$

For example consider a bar of length *L*, breadth *b* and depth *d* as shown in Fig. 8.28.

In general for any shape *volumetric strain may be taken as sum of strains in three mutually perpendicular directions*.

8.16 ELASTIC CONSTANTS

Modulus of elasticity, modulus of rigidity and bulk modulus are the three elastic constants. Modulus of elasticity (Young's Modulus) '*E*' has been already defined as the ratio of linear stress to linear strain within elastic limit. Rigidity modulus and Bulk modulus are defined in this article.

Modulus of Rigidity: It is defined as the *ratio of shearing stress to shearing strain within elastic limit and is usually denoted by letter G or N*. Thus

$$
G = \frac{q}{\phi} \tag{8.21}
$$

where $G =$ Modulus of rigidity

 $q =$ Shearing stress

and ϕ = Shearing strain

Bulk Modulus: When a body is subjected to identical stresses p in three mutually perpendicular directions, (Fig. 8.30), the body undergoes uniform changes in three directions without undergoing distortion of shape. The ratio of change in volume to original volume has been defined as volumetric strain (e_v) . Then the bulk modulus, *K* is defined as

$$
K = \frac{p}{e_v}
$$

where $p =$ identical pressure in three mutually perpendicular directions

 $e_v = \frac{\Delta_v}{v}$, Volumetric strain

 Δ ^{*v*} = Change in volume

 $v =$ Original volume

Thus bulk modulus may be defined as the ratio of identical pressure 'p' acting in three mutually perpendicular directions to corresponding volumetric strain.

Fig. 8.30

Figure 8.30 shows a body subjected to identical compressive pressure '*p*' in three mutually perpendicular directions. Since hydrostatic pressure, the pressure exerted by a liquid on a body within it, has this nature of stress, such a pressure '*p*' is called as hydrostatic pressure.

8.17 RELATIONSHIP BETWEEN MODULUS OF ELASTICITY AND MODULUS OF RIGIDITY

Consider a square element *ABCD* of sides '*a*' subjected to pure shear '*q*' as shown in Fig. 8.31. *AEC*′*D* shown is the deformed shape due to shear *q*. Drop perpendicular *BF* to diagonal *DE*. Let φ be the shear strain and *G* modulus of rigidity.

Now, strain in diagonal
$$
BD = \frac{DE - DF}{DF}
$$

= $\frac{EF}{DB}$
= $\frac{EF}{AB\sqrt{2}}$

Since angle of deformation is very small we can assume ∠*BEF* = 45°, hence *EF* = *BE* cos 45°

$$
\therefore \quad \text{Strain in diagonal} \qquad BD = \frac{EF}{BD} = \frac{BE \cos 45^{\circ}}{AB\sqrt{2}}
$$
\n
$$
= \frac{a \tan \phi \cos 45^{\circ}}{a\sqrt{2}}
$$
\n
$$
= \frac{1}{2} \tan \phi = \frac{1}{2} \phi \qquad \text{(Since } \phi \text{ is very small)}
$$
\n
$$
= \frac{1}{2} \times \frac{q}{G}, \text{ since } \phi = \frac{q}{G} \qquad \text{(1)}
$$

Now, we know that the above pure shear gives rise to axial tensile stress *q* in the diagonal direction of *DB* and axial compression q at right angles to it. These two stresses cause tensile strain along the diagonal *DB*.

Tensile strain along the diagonal
$$
DB = \frac{q}{E} + \mu \frac{q}{E} = \frac{q}{E}(1 + \mu)
$$
 ...(2)

From equations (1) and (2), we get

1 $\overline{2}$

$$
\times \frac{q}{G} = \frac{q}{E} (1 + \mu)
$$

$$
E = 2G(1 + \mu)
$$
...(8.22)

8.18 RELATIONSHIP BETWEEN MODULUS OF ELASTICITY AND BULK MODULUS

Consider a cubic element subjected to stresses *p* in the three mutually perpendicular direction *x*, *y*, *z* as shown in Fig. 8.32. Now the stress p in x direction causes tensile $\frac{z}{\sqrt{2}}$ strain $\frac{p}{E}$ in *x* direction while the stress *p* in *y* and *z* direction cause compressive strains $\mu \frac{p}{q}$ *E* in *x* direction. Hence, $e_x = \frac{p}{E}$ *p E* $-\mu \frac{p}{E}$ – $\mu \frac{p}{E}$ $=\frac{p}{E}(1 - 2\mu)$ Similarly $e_y = \frac{p}{E} (1 - 2\mu)$ **Fig. 8.32** y x p p p p p p Since angle of deformation is very small we can assume $\angle DEF = 45^\circ$, hence $EF = B$
 \therefore Strain in diagonal $BD = \frac{EF}{BD} = \frac{BE \cos 45^\circ}{AB \sin 25^\circ}$
 $= \frac{a \tan \phi \cos 45^\circ}{\pi \sqrt{2}}$
 $= \frac{a \tan \phi \cos 45^\circ}{\pi \sqrt{2}}$
 $= \frac{1}{2} \tan \phi = \frac{1}{2} \phi$

$$
e_z = \frac{p}{E}(1 - 2\mu)
$$
...(1)

$$
\therefore \quad \text{Volumetric strain} \qquad e_v = e_x + e_y + e_z = \frac{3p}{E}(1 - 2\mu)
$$

From definition, bulk modulus *K* is given by

$$
K = \frac{p}{e_v} = \frac{p}{\frac{3p(1-2\mu)}{E}}
$$

or

$$
E = 3K(1 - \mu)
$$
...(2)

Relationship between EGK:

and $E = 3K(1 - 2\mu)$...(*b*) By eliminating µ between the above two equations we can get the relationship between *E*, *G*,

We know $E = 2G(1 + \mu)$...(*a*)

 K , free from the term μ .

From equation (a)
$$
\mu = \frac{E}{2G} - 1
$$

Substituting it in equation (*b*), we get

From definition, bulk modulus K is given by
\n
$$
K = \frac{p}{e_v} = \frac{p}{3p(1-2\mu)}
$$
\nor
\n
$$
E = 3K(1 - \mu)
$$
\n
$$
E = 2G(1 + \mu)
$$
\nand
\nBy eliminating μ between the above two equations we can get the relationship between E, G,
\nK, free from the term μ .
\nFrom equation (a) $\mu = \frac{E}{2G} - 1$
\nSubstituting it in equation (b), we get
\n
$$
E = 3K[1 - 2(\frac{E}{2G} - 1)]
$$
\n
$$
= 3K(1 - \frac{E}{G} + 2) = 3K(3 - \frac{E}{G})
$$
\n
$$
= 9K - \frac{3KE}{G}
$$
\n
$$
\therefore
$$
\n
$$
E\left(1 + \frac{3K}{G}\right) = 9K
$$
\nor
\n
$$
E\left(\frac{G + 3K}{G}\right) = 9K
$$
\nor
\n
$$
E = \frac{9KG}{G + 3K}
$$
\n
$$
\therefore
$$
\n
$$
E = \frac{9KG}{2G}
$$
\n
$$
\therefore
$$
\n
$$
E = \frac{9KG}{2G}
$$
\n
$$
\therefore
$$
\n
$$
\frac{p}{E} = \frac{3}{K} + \frac{1}{K}
$$
\n
$$
\therefore
$$
\n
$$
\therefore
$$
\n
$$
\frac{p}{E} = \frac{3}{K} + \frac{1}{K}
$$
\n
$$
\therefore
$$
\n
$$
E = \frac{9KG}{2G}
$$
\n
$$
\therefore
$$
\n
$$
E = \frac{9KG}{2G}
$$
\n
$$
\therefore
$$
\n
$$
E = \frac{9KG}{2G} \quad ...(8.23a)
$$
\nEquation (c) may be expressed as
\n
$$
\frac{9}{E} = \frac{3 + 1}{6}
$$
\n
$$
\therefore
$$
\n
$$
P = 60 \text{ N} = 60000 \text{ N}
$$
\nArea $A = \frac{\pi}{4} \times 25^2 = 156.25\pi \text{ mm}^2$

Equation (*c*) may be expressed as

i.e.,
$$
\frac{9}{E} = \frac{G + 3K}{KG}
$$

$$
\frac{9}{E} = \frac{3}{G} + \frac{1}{K}
$$
...(8.23b)

*Example 8.13***.** *A bar of 25 mm diameter is tested in tension. It is observed that when a load of 60 kN is applied, the extension measured over a guage length of 200 mm is 0.12 mm and contraction in diameter is 0.0045 mm. Find Poisson's ratio and elastic constants E, G, K.* **Solution:** Now, $P = 60 \text{ kN} = 60000 \text{ N}$

Area
$$
A = \frac{\pi}{4} \times 25^2 = 156.25\pi
$$
 mm²

Guage length $L = 200$ mm

Solution: From the relationship

$$
E = 2G(1 + \mu) = 3k(1 - 2\mu)
$$

We get,

$$
G = \frac{E}{2(1 + \mu)} = \frac{2 \times 10^5}{2(1 + 0.3)} = 0.7692 \times 10^5 \text{ N/mm}^2
$$

and

$$
K = \frac{E}{3(1 + 2\mu)} = \frac{2 \times 10^5}{3(1 - 2 \times 0.3)} = 1.667 \times 10^5 \text{ N/mm}^2
$$

Longitudinal stress $= \frac{P}{A} = \frac{60 \times 10^3}{\frac{\pi}{4} \times 25^2} = 122.23 \text{ N/mm}^2$
Linear strain $= \frac{\text{Stress}}{E} = \frac{122.23}{2 \times 10^5} = 61.115 \times 10^{-5}$
Lateral strain $= e_y = -\mu e_x$ and $e_z = -\mu e_x$
Volumeetric strain $e_y = e_x + e_y + e_z$
 $= e_x(1 - 2\mu)$
 $= 61.115 \times 10^{-5} (1 - 2 \times 0.3)$
 $= 24.446 \times 10^{-5}$
but

$$
\frac{\text{Change in volume}}{v} = e_y
$$

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Now,

$$
\therefore
$$
 Change in volume = $e_v \times v$

= 24.446 × 10⁻⁵ ×
$$
\frac{\pi}{4}
$$
 × (25²) × 500
= 60 mm³

Example 8.15. A 400 mm long bar has rectangular cross-section 10 mm × 30 mm. This bar is subjected to

- *(i) 15 kN tensile force on 10 mm × 30 mm faces,*
- *(ii) 80 kN compressive force on 10 mm × 400 mm faces, and*
- *(iii) 180 kN tensile force on 30 mm × 400 mm faces.*

Find the change in volume if $E = 2 \times 10^5$ *N/mm² and* $\mu = 0.3$.

*Solution***:** The forces on the bar are as shown in Fig. 8.33. Let the *x*, *y*, *z* be the mutually perpendicular directions as shown in the figure.

$$
p_x = \frac{15 \times 10^3}{10 \times 30} = 50 \text{ N/mm}^2
$$
 (Tensile)

$$
p_y = \frac{80 \times 10^3}{10 \times 400} = 20 \text{ N/mm}^2
$$
 (Comp.)

$$
p_z = \frac{180 \times 10^3}{30 \times 400} = 15 \text{ N/mm}^2
$$
 (Tensile)

Noting that a stress produces a strain of *p*/*E* in its own direction, the nature being same as that of stress and $\mu \frac{p}{q}$ *E* in lateral direction of opposite nature, and taking tensile stress as +ve, we can write expression for strains e_x , e_y , e_z as

$$
e_x = \frac{50}{E} + \mu \frac{20}{E} - \mu \frac{15}{E}
$$

= $\frac{1}{E} (50 + 0.3 \times 20 - 0.3 \times 15) = \frac{1}{E} (51.5)$
 $e_y = -\mu \frac{50}{E} - \frac{20}{E} - \mu \frac{15}{E}$
= $\frac{1}{E} (-0.3 \times 50 - 20 - 0.3 \times 15) = -\frac{39.5}{E}$

$$
e_z = -\mu \frac{50}{E} + \mu \frac{20}{E} + \frac{15}{E}
$$

\n
$$
= \frac{1}{E}(-0.3 \times 50 + 20 \times 0.3 + 15) = \frac{6}{E}
$$

\nVolume
\n
$$
= \frac{1}{E}(51.5 - 39.5 + 6) = \frac{18}{E}
$$

\nBut
\n
$$
\frac{\text{Change in volume}}{\text{Volume}} = e_v
$$

\n
$$
\therefore \text{Change in volume} = e_v \times \text{Volume}
$$

\n
$$
= \frac{18}{2 \times 10^5} \times 10 \times 30 \times 400
$$

\n
$$
= 10.8 \text{ mm}^3
$$

*Example 8.16***.** *A bar of rectangular section shown in Fig. 8.34 is subjected to stresses px*, *py and pz in x, y and z directions respectively. Show that if sum of these stresses is zero, there is no change in volume of the bar.*

Solution: A stress *p* produces strain of magnitude p/E in its direction and a strain $\mu \frac{p}{p}$ *E* of opposite nature at right angles to its direction.

Hence $=\frac{p_x}{E}$ *p E p E* $\frac{x}{z}$ – $\mu \frac{P_y}{Z}$ – $\mu \frac{P_z}{Z}$ $e_y = -\mu \frac{p_x}{E} + \frac{p_y}{E} - \mu$ *p E p E* $\frac{x}{z} + \frac{p_y}{z} - \mu \frac{p_z}{z}$ $e_z = -\mu \frac{p_x}{E} - \mu$ *p E p E* $\frac{x}{z}$ – $\mu \frac{p_y}{z}$ + $\frac{p_z}{z}$

Now,
\n
$$
e_y = e_x + e_y + e_z
$$
\n
$$
= \frac{p_x}{E} (1 - 2\mu) + \frac{p_y}{E} (1 - 2\mu) + \frac{p_z}{E} (1 - 2\mu)
$$
\ni.e.,
\n
$$
\frac{dV}{V} = \frac{(p_x + p_y + p_z)}{E} (1 - 2\mu)
$$
\nHence
\n
$$
p_x + p_y + p_z = 0
$$
\n
$$
e_y = \frac{dV}{V} = 0
$$
\n
$$
\therefore \quad dV = 0
$$
\ni.e., there is no volumetric change, if $p_x + p_y + p_z = 0$.

*Example 8.17***.** *In a laboratory, tensile test is conducted and Young's modulus of the material is* found to be 2.1 \times 10⁵ N/mm². On the same material torsion test is conducted and modulus of rigidity is found to be 0.78 \times 10⁵ N/mm². Determine Poisson's Ratio and bulk modulus of the material.

 ϵ 0

[**Note:** This is usual way of finding material properties in the laboratory].

to 0.8 × 10⁵ N/mm² . Find its Young's Modulus and Poisson's Ratio.

Solution:
\n
$$
G = 0.4 \times 10^5 \text{ N/mm}^2
$$
\n
$$
K = 0.8 \times 10^5 \text{ N/mm}^2
$$
\nUsing the relation
$$
E = \frac{9 \times 0.4 \times 10^5 \times 0.8 \times 10^5}{3 \times 0.8 \times 10^5 + 0.4 \times 10^5}
$$
\n
$$
E = \frac{9 \times 0.4 \times 10^5 \times 0.8 \times 10^5}{3 \times 0.8 \times 10^5 + 0.4 \times 10^5}
$$
\nFrom the relation
$$
E = 2G(1 + \mu)
$$
\nwe get
$$
1.0286 \times 10^5 = 2 \times 0.4 \times 10^5 (1 + \mu)
$$
\nor
$$
1.2857 = 1 + \mu
$$
\n
$$
\mu = 0.2857
$$

8.19 COMPOSITE/COMPOUND BARS

Bars made up of two or more materials are called composite/compound bars. They may have same length or different lengths as shown in Fig. 8.35. The ends of different materials of the bar are held together under loaded conditions.

Consider a member with two materials. Let the load shared by material 1 be P_1 and that by material 2 be P_2 . Then

(*i*) From equation of equilibrium of the forces, we get

$$
P = P_1 + P_2 \tag{8.24a}
$$

(*ii*) Since the ends are held securely, we get

$$
\Delta l_1 = \Delta l_2
$$

where Δl_1 and Δl_2 are the extension of the bars of material 1 and 2 respectively

$$
\frac{P_1 L_1}{A_1 E_1} = \frac{P_2 L_2}{A_2 E_2} \tag{8.24b}
$$

Using equations 8.24(*a*) and (*b*), P_1 and P_2 can be found uniquely. Then extension of the system

can be found using the relation
$$
\Delta l = \frac{P_1 L_1}{A_1 E_1}
$$
 or $\Delta l = \frac{P_2 L_2}{A_2 E_2}$ since $\Delta l = \Delta l_1 = \Delta l_2$.

The procedure of the analysis of compound bars is illustrated with the examples below:

Example 8.19. A compound bar of length 600 mm consists of a strip of aluminium 40 mm wide and 20 mm thick and a strip of steel 60 mm wide × 15 mm thick rigidly joined at the ends. If elastic modulus of aluminium and steel are 1×10^5 *N/mm² and* 2×10^5 *N/mm², determine the stresses developed in each material and the extension of the compound bar when axial tensile force of 60 kN acts.*

i.e.

Extension of the compound bar =
$$
\frac{P_a L}{A_a E_a} = \frac{18462 \times 600}{800 \times 1 \times 10^5}
$$

∆*l* **= 0.138 mm.**

Example 8.20. A compound bar consists of a circular rod of steel of 25 mm diameter rigidly fixed into a copper tube of internal diameter 25 mm and external diameter 40 mm as shown in Fig. 8.37. If the compound bar is subjected to a load of 120 kN, find the stresses developed in the two materials.

> *Take* $E_s = 2 \times 10^5 \text{ N/mm}^2$ *and* $E_c = 1.2 \times 10^5 \text{ N/mm}^2$.

Solution: Area of steel rod
$$
A_x = \frac{\pi}{4} \times 25^2 = 490.87 \text{ mm}^2
$$

\nArea of copper tube $A_c = \frac{\pi}{4} (40^2 - 25^2) = 765.76 \text{ mm}^2$
\nFrom equation of equilibrium,
\nwhere P_s is the load shared by steel rod and P_c is the load shared by the copper tube.
\nFrom compatibility condition, we have
\n
$$
\Delta_s = \Delta_c
$$
\n
$$
\frac{P_s I}{A_f E_c} = \frac{P_c}{A_c E_c}
$$
\n
$$
\frac{P_s}{490.87 \times 2 \times 10^5} = \frac{P_c}{765.76 \times 1.2 \times 10^5}
$$
\n
$$
\therefore P_s = 1.068 \quad P_c
$$
\nFrom eqns. (1) and (2), we get
\n
$$
1.068 \quad P_c + P_c = 120 \times 1000
$$
\n
$$
\therefore P_c = \frac{120 \times 1000}{20.68} = 58027 \text{ N}
$$
\n
$$
\therefore P_c = \frac{120 \times 1000}{20.68} = 58027 \text{ N}
$$
\n
$$
\therefore P_c = \frac{120 \times 1000}{20.68} = 58027 \text{ N}
$$
\n
$$
\therefore P_c = \frac{120 \times 1000}{20.68} = 58027 \text{ N}
$$
\n
$$
\therefore \text{Stress in copper} = \frac{58027}{9765.76} = 75.78 \text{ N/mm}^2
$$
\n
$$
\therefore \text{Stress in steel } \frac{10373}{400.87} = 126.25 \text{ N/mm}^2
$$
\n
$$
\therefore \text{Stress in steel } \frac{10373}{400.87} = 126.25 \text{ N/mm}^2
$$
\n
$$
\therefore \text{Stress in steel } \frac{10373}{400.67} = 126.25 \text{ N/mm}^2
$$
\n
$$
\therefore \text{Stress in steel } \frac{10373}{400.67} = 126.2
$$

Example 8.21. A reinforced concrete column of size 300 mm × 500 mm has 8 steel bars of 16 mm diameter as shown in Fig. 8.38. If the column is subjected to an axial compressive force of 800 kN, find the

stresses developed in steel and concrete. Take E E s c = 18.

Solution: Area of steel = $8 \times \frac{\pi}{4} \times 16^2 = 1608.5$ mm²

∴ Area of concrete = $300 \times 500 - 1608.5 = 148391.5$ mm² From equilibrium condition,

$$
P_c + P_s = 800 \times 1000 \tag{1}
$$

From compatibility condition,

$$
\Delta_s = \Delta_c
$$

$$
\frac{P_s L}{A_s E_s} = \frac{P_c L}{A_c E_c}
$$

$$
\frac{P_s}{1608.5 E_s} = \frac{P_c}{148391.5 E_c}
$$

Fig. 8.38

300

600

ö

∙

 \bullet

$$
P_c = \frac{148391.5}{1608.5} \times \frac{E_c}{E_s} P_s
$$

= $\frac{148391.5}{1608.5} \times \frac{1}{18} P_s$
= 5.125 P_s ...(2)

From eqns. (1) and (2), we get

 P_s + 5.125 P_s = 800 × 1000 ∴ $P_s = 130612$ N. Hence $P_c = 5.125 P_s = 669388 \text{ N}$

Hence **stress in concrete** = $\frac{P_a}{A_b}$ *c c* $=\frac{669388}{148391.5}$ = **4.51** N/mm²

and **stress in steel** = $\frac{P_s}{A}$ *s s* $=\frac{130612}{1608.5}$ = **81.2** N/mm²

Example 8.22. Three pillars, two of aluminium and one of steel support a rigid platform of 250 kN as shown in Fig. 8.39. If area of each aluminium pillar is 1200 mm2 and that of steel pillar is 1000 mm² , find the stresses developed in each pillar. From eqns. (1) and (2), we get
 $P_2 + 5.125$ $P_3 = 130612$ N.

Hence $P_4 = 5.125$ $P_5 = 609388$ N.

Hence arress in concrete $= \frac{R}{A_0} = \frac{130612}{148391.5} = 4.51$ Ntmm²

and stress in sected $= \frac{R}{A_0} = \frac{130612}{1608.$

Take
$$
E_s = 2 \times 10^5
$$
 N/mm² and $E_a = 1 \times 10^6$ N/mm².
Solution: Let force shared by each aluminum pillar be P_a and that

shared by steel pillar be P_s .

∴ The forces in vertical direction = $0 \rightarrow$

$$
P_a + P_s + P_a = 250
$$

2P_a + P_s = 250 ...(1)

From compatibility condition, we get

$$
\Delta_s = \Delta_a
$$

$$
\frac{P_s L_s}{A_s E_s} = \frac{P_a L_a}{A_a E_a}
$$

$$
\frac{P_s \times 240}{1000 \times 2 \times 10^5} = \frac{P_a \times 160}{1200 \times 1 \times 10^5}
$$

∴ $P_s = 1.111 P_a$...(2)

From eqns. (1) and (2), we get

$$
P_a (2 + 1.111) = 250
$$

$$
P_a = 80.36 \text{ kN}
$$

Hence from eqn. (1),

$$
P_s = 250 - 2 \times 80.36 = 89.28
$$
 kN

∴ Stresses developed are

$$
\sigma_s = \frac{P_s}{A_s} = \frac{89.28 \times 1000}{1000} = 89.28 \text{ N/mm}^2
$$

$$
\sigma_a = \frac{80.36 \times 1000}{1200} = 66.97 \text{ N/mm}^2
$$

Example 8.23. A steel bolt of 20 mm diameter passes centrally through a copper tube of internal diameter 28 mm and external diameter 40 mm. The length of whole assembly is 600 mm. After tight fitting of the assembly, the nut is over tightened by quarter of a turn. What are the stresses introduced in the bolt and tube, if pitch of nut is 2 mm? Take $E_s = 2 \times 10^5$ *N/mm² and* $E_c = 1.2 \times 10^5$ *N/mm².*

Solution: Figure 8.40 shows the assembly. Let the force shared by bolt be P_s and that by tube be P_c . Since there is no external force, static equilibrium condition gives

$$
P_s + P_c = 0 \quad \text{or} \quad P_s = -P_c
$$

i.e., the two forces are equal in magnitude but opposite in nature. Obviously bolt is in tension and tube is in compression.

Let the magnitude of force be *P*. Due to quarter turn of the nut, the nut advances by $\frac{1}{4}$ $\frac{1}{4}$ × pitch

 $=\frac{1}{4}$ 4 $\times 2 = 0.5$ mm.

[**Note.** Pitch means advancement of nut in one full turn]

 P_sL $A_s E$

 sL_s

During this process bolt is extended and copper tube is shortened due to force *P* developed. Let ∆*s* be extension of bolt and ∆*c* shortening of copper tube. Final position of assembly be ∆, then

$$
\Delta_s + \Delta_c = \Delta
$$

$$
P_c L_c = 0
$$

 $+\frac{4cE_c}{4E}$ = 0.5

 $P_c L$ *A E*

i.e.

$$
A_s L_s
$$

\n
$$
P \times 600
$$

\n
$$
\frac{P \times 600}{(\pi/4) \times 20^2 \times 2 \times 10^5} + \frac{P \times 600}{(\pi/4)(40^2 - 28^2) \times 1.2 \times 10^5} = 0.5
$$

\n
$$
\frac{P \times 600}{(\pi/4) \times 10^5} \left[\frac{1}{20^2 \times 2} + \frac{1}{(40^2 - 28^2) \times 1.2} \right] = 0.5
$$

\n
$$
P = 28816.8 \text{ N}
$$

$$
p_s = \frac{P_s}{A_s} = \frac{28816.8}{(\pi/4) \times 20^2} = 91.72 \text{ N/mm}^2
$$

$$
p_c = \frac{P_c}{A_c} = \frac{28816.8}{(\pi/4)(40^2 - 28^2)} = 44.96 \text{ N/mm}^2
$$

8.20 THERMAL STRESSES

Every material expands when temperature rises and contracts when temperature falls. It is established experimentally that the change in length ∆ is directly proportional to the length of the member *L* and change in temperature *t.* Thus

$$
\Delta \propto tL
$$

= α tL ...(8.25)

The constant of proportionality α is called coefficient of thermal expansion and is defined as change in unit length of material due to unit change in temperature. Table 8.1 shows coefficient of thermal expansion for some of the commonly used engineering materials:

If the expansion of the member is freely permitted, as shown in Fig. 8.41, no temperature stresses are induced in the material.

Fig. 8.41 Free Expansion Permitted

If the free expansion is prevented fully or partially the stresses are induced in the bar, by the support forces. Referring to Fig. 8.42,

The expansion is prevented by developing compressive force *P* at supports

$$
\therefore \quad \frac{PL}{AE} = \Delta = \alpha \, tL - \delta. \tag{8.27}
$$

Example 8.24. A steel rail is 12 m long and is laid at a temperature of 18°C. The maximum temperature expected is 40°C.

- (*i*) *Estimate the minimum gap between two rails to be left so that the temperature stresses do not develop.*
- (*ii*) *Calculate the temperature stresses developed in the rails, if:*
	- (*a*) *No expansion joint is provided.*
	- (*b*) *If a 1.5 mm gap is provided for expansion.*
- (*iii*) *If the stress developed is 20 N/mm2 , what is the gap provided between the rails? Take E* = 2 × 10^5 *N*/mm² and α = 12 × 10^{-6} /°C.

Solution:

(*i*) The free expansion of the rails

$$
= \alpha tL = 12 \times 10^{-6} \times (40 - 18) \times 12.0 \times 1000
$$

= 3.168 mm

∴ **Provide a minimum gap of 3.168 mm between the rails, so that temperature stresses do not develop.**

(*ii*) (*a*) If no expansion joint is provided, free expansion prevented is equal to 3.168 mm.

∴

i.e.
\n
$$
\frac{PL}{AE} = 3.168 \text{ mm}
$$
\n
$$
\frac{PL}{AE} = 3.168
$$
\n
$$
\therefore \quad p = \frac{P}{A} = \frac{3.168 \times 10^{-14} \text{ m}}{100}
$$

$$
= \frac{P}{A} = \frac{3.168 \times 2 \times 10^5}{12 \times 1000} = 52.8 \text{ N/mm}^2
$$

(*b*) If a gap of 1.5 mm is provided, free expansion prevented $\Delta = \alpha tL - \delta = 3.168 - 1.5 =$ 1.668 mm.

∴ The compressive force developed is given by *PL* $\frac{12}{AE}$ = 1.668

or
$$
p = \frac{P}{A} = \frac{1.668 \times 2 \times 10^5}{12 \times 1000} = 27.8 \text{ N/mm}^2
$$

(*iii*) If the stress developed is 20 N/mm², then $p = \frac{P}{P}$ *A* $= 20$

If δ is the gap, ∆ = α *tL* – δ ∴ *PL AE* = 3.168 – δ *i.e.* 20 × 12 1000 2 10⁵ × × = 3.168 – δ ∴ δ = 3.168 – 1.20 = **1.968 mm** Civil[datas.blogspo](http://Civildatas.blogspot.in)t.in

Example 8.25. The temperature of a steel ring is raised through 150°C in order to fit it on a wooden wheel of 1.2 m diameter. Find the original diameter of the steel ring and also the stresses developed in the ring, when it cools back to normal temperature. Assume $E_s = 2 \times 10^5$ *N/mm² and* $\alpha_s = 12$ $\times 10^{-6}$ /°C.

Solution: Let *D* be the diameter of ring after heating and '*d*' be its diameter before heating

∴ $D = 1.2$ m = 1200 mm. Circumference of ring after heating $= \pi D$ Circumference of ring before heating $(L) = \pi d$ $πD = πd + α_s t (πd)$ $D = d + \alpha_s t \, d = (1 + \alpha_s t) d$ *i.e.*, $1200 = [1 + 12 \times 10^{-6} \times 150] d$ ∴ $d = 1197.84$ mm. when it cools expansion prevented $Δ = π(D - d) = α_s$ tπ*d* = 12 × 10⁻⁶ × 150 π*d* $= 1800 \times 10^{-6}$ πd ∴ *PL* $\frac{I}{AE}$ = 1800 × 10⁻⁶ πd *P A d* $\left(\frac{P}{A}\right) \times \frac{\pi \times d}{E} = 1800 \times 10^{-6}$ π*d* $\overline{1}$ $\overline{}$ **Stress** *p* = $\frac{P}{A}$ = 1800 × 10⁻⁶ × 2 × 10⁵ = 360 N/mm² *Example 8.26. The composite bar shown in Fig. 8.44 is rigidly fixed at the ends A and B. Determine the reaction developed at ends when the temperature is raised by 18°C. Given Ea = 70 kN/mm²* $E_s = 200 \text{ kN/mm}^2$ $\alpha_a = 11 \times 10^{-6} / \text{°C}$ $\alpha_s = 12 \times 10^{-6} / \text{°C}$ **Solution:** Let D be the diameter of ring after heating and 'd' be its diameter before heat
 \therefore

Circumference of ring forte heating = πD

Circumference of ring forte heating (1) = πd
 \therefore $\pi D = \pi d + \alpha$, t ($\$

Fig. 8.44

Free expansion =
$$
\alpha_a tL_a + \alpha_s tL_s
$$

= 11 × 10⁻⁶ × 18 × 1500 + 12 × 10⁻⁶ × 18 × 3000
= 0.945 mm

Since this is prevented

$$
\Delta = 0.945 \text{ mm.}
$$

E_a = 70 kN/mm² = 70000 N/mm² ;

$$
E_s = 200 \text{ kN/mm}^2 = 200 \times 1000 \text{ N/mm}^2
$$

If *P* is the support reaction,

$$
\Delta = \frac{PL_a}{A_a E_a} + \frac{PL_s}{A_s E_s}
$$

i.e.
$$
0.945 = P \left[\frac{1500}{600 \times 70000} + \frac{3000}{400 \times 200 \times 1000} \right]
$$

$$
0.945 = 73.214 \times 10^{-6} P
$$
or
$$
P = 12907.3 \text{ N}
$$

Example 8.27. The steel bar AB shown in Fig. 8.45 varies linearly in diameter from 25 mm to 50 mm in a length 500 mm. It is held between two unyielding supports at room temperature. What is the stress induced in the bar, if temperature rises by $25^{\circ}C$ *? Take* $E_s = 2 \times 10^5$ *N/mm² and* $\alpha_s = 12 \times 10^{-6}$ /°C.

Solution: Free expansion of bar = α *tL*

$$
= 12 \times 10^{-6} \times 25 \times 500
$$

$$
= 0.15 \text{ mm}.
$$

This expansion is prevented ∴ $\Delta = 0.15$ mm.

If *P* is the force developed by supports, then it can cause a contraction of

$$
\frac{4PL}{\pi d_1 d_2 E} = \frac{4 \times P \times 500}{\pi \times 25 \times 50 \times 2 \times 10^5}
$$

Equating it to the contraction prevented, for satisfying the compatibility, we get

$$
\frac{4P \times 500}{\pi \times 25 \times 50 \times 2 \times 10^5} = 0.15
$$

$$
\sim
$$

$$
P = 58905 \text{ N}
$$

$$
\therefore \text{ Corresponding maximum stress} = \frac{P}{A_{\min}} = \frac{58905}{\frac{\pi}{4} \times 25^2}
$$

$$
= 120 \text{ N/mm}^2.
$$

4

8.21 THERMAL STRESSES IN COMPOUND BARS

When temperature rises the two materials of the compound bar experience different free expansion. Since they are prevented from seperating, the two bars will have common position. This is possible only by extension of the bar which has less free expansion and contraction of the bar which has more free expansion. Thus one bar develops tensile force and another develops the compressive force. In this article we are interested to find such stresses.

Consider the compound bar shown in Fig. 8.46(*a*). Let α_1 , α_2 be coefficient of thermal expansion and E_1 , E_2 be moduli of elasticity of the two materials respectively. If rise in temperature is '*t*',

Free expansion of bar $1 = \alpha_1$ *tL*

Free expansion of bar $2 = \alpha_2 tL$

Let $\alpha_1 > \alpha_2$. Hence the position of the two bars, if the free expansions are permitted are at *AA* and *BB* as shown in Fig. 8.46(*b*)

Since the two bars are rigidly connected at the ends, the final position of the end will be somewhere between *AA* and *BB*, say at *CC*. It means Bar–1 will experience compressive force P_1 which contracts it by Δ_1 and Bar–2 experience tensile force P_2 which will expand it by Δ_2 .

The equilibrium of horizontal forces gives,

$$
P_1 = P_c
$$
, say P

From the Fig. 8.46 (*b*), it is clear,

$$
\alpha_1 tL - \Delta_1 = \alpha_2 tL + \Delta_2
$$

$$
\Delta_1 + \Delta_2 = \alpha_1 tL - \alpha_2 tL = (\alpha_1 - \alpha_2) tL.
$$

If the cross-sectional areas of the bars are A_1 and A_2 , we get

$$
\frac{PL}{A_1E_1} + \frac{PL}{A_2E_2} = (\alpha_1 - \alpha_2) t L
$$
...(8.28)

From the above equation force *P* can be found and hence the stresses P_1 and P_2 can be determined.

Example 8.28. A bar of brass 20 mm is enclosed in a steel tube of 40 mm external diameter and 20 mm internal diameter. The bar and the tubes are initially 1.2 m long and are rigidly fastened at both ends using 20 mm diameter pins. If the temperature is raised by 60°C, find the stresses induced in the bar, tube and pins.

 $Given:$

$$
E_s = 2 \times 10^5 \text{ N/mm}^2
$$

\n
$$
E_b = 1 \times 10^5 \text{ N/mm}^2
$$

\n
$$
\alpha_s = 11.6 \times 10^{-6} / ^{\circ}C
$$

\n
$$
\alpha_b = 18.7 \times 10^{-6} / ^{\circ}C.
$$

Solution:

$$
t = 60^{\circ} \quad E_s = 2 \times 10^5 \text{ N/mm}^2 \quad E_b = 1 \times 10^5 \text{ N/mm}^2
$$

\n
$$
\alpha_s = 11.6 \times 10^{-6} / \text{°C} \quad \alpha_b = 18.7 \times 10^{-6} / \text{°C}
$$

\n
$$
A_s = \frac{\pi}{4} (40^2 - 20^2) \qquad A_b = \frac{\pi}{4} \times 20^2
$$

\n= 942.48 mm² = 314.16 mm²

Since free expansion of brass (α_b, t) is more than free expansion of steel (α_s, t) , compressive force P_b develops in brass and tensile force P_s develops in steel to keep the final position at CC (Ref: Fig. 8.47).

Horizontal equilibrium condition gives $P_b = P_s$, say *P*. From the figure, it is clear that

$$
\Delta_s + \Delta_b = \alpha_b tL - \alpha_s tL = (\alpha_b - \alpha_s)tL.
$$

where Δ _s and Δ _{*b*} are the changes in length of steels and brass bars.

Given:
\n
$$
E_x = 2 \times 10^5 \text{ N/mm}^2
$$
\n
$$
\alpha_x = 11.6 \times 10^{-6} \text{/C}
$$
\nSolution:
\nSolution:
\n
$$
\begin{array}{rcl}\n\text{Spin}(n) & \text{Spin}(n) \\
\hline\n\end{array}
$$
\n
$$
\begin{array}{rcl}\n\text{Fig. 8.47} \\
\hline\n\end{array}
$$
\n
$$
\begin{array}{rcl}\n\text{Sine (free expansion of brass } (a_k, iL) \text{ is more than free expansion of steel } (a_k, iL) \text{ compi} \\
\hline\n\end{array}
$$
\n
$$
\begin{array}{rcl}\n\text{Since free expansion of brass } (a_k, iL) \text{ is more than free expansion of steel } (a_k, iL) \text{ compi} \\
\hline\n\end{array}
$$
\n
$$
\begin{array}{rcl}\n\text{Since } P_k \text{ develops in brass and tensile force } P_s \text{ develops in steel to keep the final position (Ref: Fig. 8.47).\n\end{array}
$$
\n
$$
\begin{array}{rcl}\n\therefore & \frac{PL}{A_kE_x} + \frac
$$

The pin resist the force *P* at the two cross-sections at junction of two bars.

$$
\therefore \qquad \text{Shear stress in pin} = \frac{P}{2 \times \text{Area of pin}}
$$
\n
$$
= \frac{11471.3}{2 \times \pi/4 \times 20^2} = 18.26 \text{ N/mm}^2
$$

Example 8.29. A compound bar is made of a steel plate 50 mm wide and 10 mm thick to which copper plates of size 40 mm wide and 5 mm thick are connected rigidly on each side as shown in Fig. 8.48. The length of the bar at normal temperature is 1 m. If the temperature is raised by 80°, determine the stresses in each metal and the change in length. Given $\alpha_s = 12 \times 10^{-6}$ */°C,* $\alpha_c = 17 \times$ 10^{-6} \degree C, $E_s = 2 \times 10^5$ *N/mm*², $E_c = 1 \times 10^{5-}$ *N/mm*².

Solution: Now, $L = 1000$ mm, $A_s = 50 \times 10 = 500$ mm², A_c = Free expansion of copper $\alpha_c tL$ is greater than free expansion of steel $\alpha_s tL$. To bring them to the same position, tensile force P_s acts on steel plate and compressive force P_c acts on each copper plate.

$$
P_s = 2P_c
$$

If Δ _s and Δ _c are changes in the length of steel and copper plates. $\Delta_{s} + \Delta_{c} = (\alpha_{c} - \alpha_{s}) tL$

$$
\frac{P_s L}{A_s E_s} + \frac{P_c L}{A_c E_c} = (\alpha_c - \alpha_s) tL
$$

But $P_s = 2P_c$

∴

$$
\therefore \frac{2P_c}{500 \times 2 \times 10^5} + \frac{P_c}{200 \times 1 \times 10^5} = (17 - 12) \times 10^{-6} \times 80
$$

\n
$$
\therefore \frac{P_c}{s} = 5714.28 \text{ N}
$$

\n
$$
\therefore \frac{P_s}{s} = 11428.56 \text{ N}
$$

... Stress in copper =
$$
\frac{P_c}{A_c} = \frac{5714.28}{200} = 28.57 \text{ N/mm}^2
$$

Stress in steel =
$$
\frac{P_s}{A_s} = \frac{5714.28 \times 2}{500} = 22.86 \text{ N/mm}^2
$$

∴ **Change in length** = α*s t L* + δ*^s*

=
$$
\alpha_s t L + \frac{P_s L_s}{A_s E_s}
$$

= $12 \times 10^{-6} \times 80 \times 1000 + \frac{11428.56 \times 1000}{500 \times 2 \times 10^5}$
= **1.07 mm**

8.22 HOOP STRESSES

The internal or external pressure applied to thin cylinders is resisted by stresses developed in the circumferential direction of the cylinder. This type of stress is called hoop stress. Consider a thin cylinder of diameter subject to internal press *p* as shown in Fig. 8.49(*a*).

Consider half the section shown in Fig. 8.49(*b*). The normal pressure on the element of unit length as shown in Fig. 8.49(*b*)

$$
= p \times \frac{D}{2} d\theta
$$

Its vertical component = $\frac{pD}{2}d\theta$ cos θ

∴ Bursting force normal to the horizontal section

$$
= 2 \int_0^{\pi/2} p \frac{d}{2} \cos \theta \, d\theta = pd \left[\sin \theta \right]_0^{\pi/2}
$$

$$
= pd.
$$

This bursting force is resisted by the hoop stresses *f*. Let '*t*' be the thickness of cylinder. Then resisting force per unit length of cylinder.

$$
= 2 \times ft
$$

Equating resisting force to bursting force, we get

$$
2ft = pd
$$

or

$$
f = \frac{pd}{2t}
$$
...(8.29)
Example 8.30. A thin cylinder of internal diameter $D = 1$ *m and thickness t = 12 mm is subjected to internal pressure of 2 N/mm² . Determine the hoop stress developed.*

Solution: **Hoop** stress $f = \frac{pd}{2t} = \frac{2 \times 1}{2 \times 1}$ $\frac{2 \times 1000}{2 \times 12}$ = **83.33** N/mm²

IMPORTANT FORMULAE

1. If stress is uniform

$$
p = \frac{P}{A}
$$

- 2. (*i*) Linear strain $=$ $\frac{\text{Change in length}}{\text{Original length}}$
	- (iii) Lateral strain $=$ Change in lateral dimension Original lateral dimension
- **3.** Poisson's ratio $=$ $\frac{\text{Lateral strain}}{\text{Linear strain}}$, within elastic limit.

4. Percentage elongation =
$$
\frac{L'-L}{L} \times 100.
$$

5. Percentage reduction in area =
$$
\frac{A - A'}{A} \times 100
$$
.

- **6.** Nominal stress $=$ $\frac{\text{Load}}{\text{Original cross-sectional area}}$.
- 7. True stress $=$ $\frac{\text{Load}}{\text{Actual cross-sectional area}}$.
- **8.** Factor of safety $=$ $\frac{\text{Ultimate stress}}{\text{Working stress}}$

However in case of steel, $=$ $\frac{\text{Yield stress}}{\text{Working stress}}$.

9. Hooke's Law, *p* = *Ee.*

10. Extension/shortening of bar = $\frac{PL}{AE}$.

11. Extension of flat bar with linearly varying width and constant thickness = $\frac{PL}{tE(b_1 - b_2)} \log \frac{b}{b}$ 1 2 . **12.** Extension of linearly tapering rod = $\frac{4}{\pi}$ $d_1 d_2$ $(\pi/4 \, d_1 d_2)$ *PL* $E d_1 d$ $\frac{4PL}{\pi E d_1 d_2} = \frac{PL}{(\pi/4 d_1 d_2) E}.$ 2*i* 2 × 12
 IMPORTANT FORMULAE

1. If stress is uniform
 $p = \frac{P}{A}$

2. *(i)* Linear strain = Change in lateral dimension
 (ii) Lateral strain = Change in lateral dimension

3. Poisson's ratio = Lateral strain $\frac{P$

- **13.** Direct shear stress = $\frac{Q}{A}$.
- **14.** Volumetric strain $e_y = \frac{\delta V}{V} = e_x + e_y + e_z$.
- **15.** $E = 2G(1 + \mu) = 3K(1 2\mu)$
	- or

$$
\frac{9}{E} = \frac{3}{G} + \frac{1}{K}.
$$

16. Extension due to rise in temperature:

$$
\Delta = \alpha tL
$$

17. Thermal force, *P* is given by

= extension prevented.

THEORY QUESTIONS

- **1.** Explain the terms: stress, strain, modulus of elasticity.
- **2.** Draw a typical stress-strain curve for mild steel, indicate salient point and define them.
- **3.** Distinguish between
	- (*i*) Elastic limit and yield point. (*ii*) Nominal stress and true stress.

PL

- **4.** Explain the term factor of safety.
- **5.** A bar of uniform thickness *t* tapers uniformly from a width $b₁$ at one end to $b₂$ at the other end in a length *L*. Find the expression for its extension under an axial pull *P*.
- **6.** A tapering rod has diameter d_1 at one end and it tapers uniformly to a diameter d_2 at the other end in a length *L*. If the modulus of elasticity is *E*, find the change in length when subjected to an axial force *P*. 14. Coulmete strain $e_i = \frac{\gamma}{V} = e_s + e_s + e_s$

15. $E = 2G (1 + \mu) = 3K (1 - 2\mu)$

or

cor

civilgation due to rise in temperature:

16. Extension due to rise in temperature:

17. Thermal force, P is given by
 $\frac{P L}{AE} =$ extensio
	- **7.** Show that pure shear gives rise to tensile and compressive stresses across the planes inclined at 45° to the shearing planes, the intensity of the direct stresses being of the same magnitude as shearing stress.
	- **8.** Show that volumetric strain in a bar is equal to the sum of strains in three mutually perpendicular directions.
	- **9.** Explain the terms:
		- (*i*) Modulus of elasticity (*ii*) Modulus of rigidity and
		- (*iii*) Bulk modulus.
	- **10.** Derive the relationship between
		- (*i*) Modulus of elasticity and modulus of rigidity
		- (*ii*) Modulus of elasticity and bulk modulus.
	- **11.** Show that in a compound bar of length *L*, when temperature increases by *t*, the force *'P'* developed is given by

$$
\frac{PL}{A_1E_1} + \frac{PL}{A_2E_2} = (\alpha_1 - \alpha_2) tL
$$

where A_1 , A_2 -Cross-sectional areas of bar 1 and bar 2 respectively E_1 , E_2 –Young's moduli of bar 1 and bar 2 respectively and

 α_1 , α_2 –are coefficient of thermal expansion of bars 1 and 2 respectively.

12. Show that in a thin cylinder subject to internal pressure '*p*' is given by

$$
f = \frac{pd}{2t}
$$

where *d*–diameter of the cylinder

t–thickness of cylinder.

PROBLEMS FOR EXERCISE

1. Tension test was conducted on a specimen and the following readings were recorded.

Diameter $= 25$ mm

Gauge length of extensometer = 200 mm

Least count of extensometer = 0.001 mm

At a load of 30 kN, extensometer reading $= 60$

At a load of 50 kN, extensometer reading = 100

Yield load $= 160$ kN

Maximum load $= 205$ kN

Diameter neck = 17 mm

Final extension over 125 mm original length = 150 mm

Find Young's Modulus, yield stress, ultimate stress, percentage elongation and percentage reduction in area.

[Ans. $E = 2.0372 \times 10^5$ N/mm², $p_y = 325.95$ N/mm². Ultimate stress = 417.62 N/mm², % elongation = 20% reduction in area = 53.76]

2. The composite bar shown in Fig. 8.50 is subjected to a tensile force of 30 kN. The extension observed is 0.44. Find the Young's modulus of brass, if Young's modulus of steel is 2×10^5 N/mm². **[Ans.** 99745 N/mm²]

3. The steel flat shown in Fig. 8.51 has uniform thickness of 20 mm. Under an axial load of 100 kN, its extension is found to be 0.21 mm. Determine the Young's modulus of the material.

 $[Ans. E = 2.06293 \times 10^5 \text{ N/mm}^2]$

4. Find the extension of the bar shown in Fig. 8.52 under an axial load of 25 kN.

5. The compound bar shown in Fig. 8.54 consists of three materials and supports a rigid horizontal platform. Find the stresses developed in each bar, if the platform remains horizontal even after the loading with 100 kN. Given

[Hint: $\Delta_1 = \Delta_2 = \Delta_3 = \Delta$, Find $P_2 - P_3$ in terms of $P_1 \cdot P_1 + P_2 + P_3 = 100 \times 10^3$] [Ans. $p_1 = 59.52$ N/mm², $p_2 = 28.57$ N/mm², $p_3 = 19.84$ N/mm²]

6. The composite bar shown in Fig. 8.53 is 0.20 mm short of distance between the rigid supports at room temperature. What is the maximum temperature rise which will not produce stresses in the bar ? Find the stresses induced in the bar when temperature rise is 60°. Given:

$$
\alpha_s = 12 \times 10^{-6} / ^{\circ}C \qquad \alpha_c = 17.5 \times 10^{-6} / ^{\circ}C
$$

\n
$$
E_s = 2 \times 10^5 \text{ N/mm}^2 \qquad E_c = 1.2 \times 10^5 \text{ N/mm}^2
$$

\n
$$
A_s = 400 \text{ mm}^2 \qquad A_c = 300 \text{ mm}^2
$$

\n[Ans. (a) 42.10 °C, (b) $p_s = 36 \text{ N/mm}^2$, $p_c = 48 \text{ N/mm}^2$]

9 Beams

9.1 INTRODUCTION

A beam may be defined as a structural element which has one dimension considerably larger than the other two dimensions, namely breadth and depth, and is supported at few points. The distance between two adjacent supports is called span. It is usually loaded normal to its axis. The applied loads make every cross-section to face bending and shearing. The load finally get transferred to supports. The system of forces consisting of applied loads and reactions keep the beam in equilibrium. The reactions depend upon the type of supports and type of loading. In this chapter type of supports, types of beams and types of loading are first explained and then the methods of finding reactions, bending moment and shear forces are illustrated for the following determinate beams: **BE:**
 CALC
 CALC

- (*a*) Simply supported beams
- (*b*) Cantilever beams and
- (*c*) Overhanging beams.

9.2 TYPES OF SUPPORTS

Various types of supports and reactions developed are listed below:

Simple Support: If the beam rests simply on a support it is called a simple support. In such case the reaction at the support is at right angles to the support and the beam is free to move in the direction of its axis and also it is free to rotate about the support (Fig. 9.1).

Roller Support: In this case, beam end is supported on rollers. In such cases, reaction is normal to the support since rollers can be treated as frictionless. [Fig. 9.2 (*a*)]. Many mechanical components are having roller supports which roll between guides. In such cases, reaction will be normal to the guides, in both the direction (Fig. 9.2). At roller support beam is free to move along the support. It can rotate about the support also.

Hinged Support: At a hinged end, a beam cannot move in any direction. However, it can rotate about the support (Fig. 9.3). Hence the support will not develop any resisting moment, but it can develop reaction in any direction to keep the end stationary. The reaction *R* can be split into its horizontal and vertical components for the purpose of analysis.

Thus, in Fig. 9.3,

$$
V_A = R \sin \theta
$$

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$H_A = R \cos \theta$

Fixed Support: At such supports, the beam end is not free to translate or rotate. Translation is prevented by developing support reaction in any required direction.

Referring to Fig. 9.4 the support reaction *R* which is at an angle θ to *x* axis may be represented by H_A and V_A , where

 $V_A = R \sin θ$

H_A = *R* cos θ

Rotation is prevented by developing support moment M_A as shown in Fig. 9.4. Thus at fixed support *A*, there are three reactions H_A , V_A and M_A .

 M_A

 R V_A

 θ

9.3 TYPES OF BEAMS

Simply Supported Beam: When both end of a beam are simply supported it is called simply supported beam (Fig. 9.5). Such a beam can support load in the direction normal to its axis.

Beam with One End Hinged and the Other on Rollers: If one end of a beam is hinged and other end is on rollers, the beam can resist load in any direction (*see* Fig. 9.6).

Over-hanging Beam: If a beam is projecting beyond the support. It is called an over-hanging beam (Fig. 9.7). The overhang may be only on one side as in Fig. 9.7(*a*) or may be on both sides as in Fig. 9.7(*b*).

Cantilever Beam: If a beam is fixed at one end and is free at the other end, it is called cantilever beam (Fig. 9.8).

Propped Cantilever: It is a beam with one end fixed and the other end simply supported (Fig. 9.9).

Both Ends Hinged: In these beams both ends will be having hinged supports (Fig. 9.10).

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Continuous Beam: A beam is said to be continuous, if it is supported at more than two points (Fig. 9.11).

In the case of simply supported beams, beams with one end hinged and the other on rollers, cantilever and over-hanging beams, it is possible to determine the reactions for given loadings by using the equations of equilibrium only. In the other cases, the number of independent equilibrium equations are less than the number of unknown reactions and hence it is not possible to analyse them by using equilibrium equations alone. The beams which can be analysed using only equilibrium equations are known as *Statically Determinate* beams and those which cannot be analysed are known as *Statically Indeterminate* beams. The latter beams can be analysed using the conditions of continuity in deformations in addition to equilibrium equations. Such cases will not be treated in this book. In the case of simply supported beams, because this touse end hinged and the other
in the case of simply any is it is possible to determine the reactions for given level
using the equations of equations of consideration o

9.4 TYPES OF LOADING

Usual types of loadings on the beams are discussed here.

Concentrated Loads: If a load is acting on a beam over a very small length, it is approximated as acting at the mid point of that length and is represented by an arrow as shown in Fig. 9.12.

Uniformly Distributed Load (UDL): Over considerably long distance such load has got uniform intensity. It is represented as shown in Fig. 9.13 (*a*) or as in (*b*). For finding reaction, this load may be assumed as total load acting at the centre of gravity of the loading (middle of the loaded length). For example, in the beam shown in Fig. 9.13, the given load may be replaced by a $20 \times 4 = 80$ kN concentrated load acting at a distance 2 m from the left support.

Uniformly Varying Load: The load shown in Fig. 9.14 varies uniformly from *C* to *D.* Its intensity is zero at *C* and is 20 kN/m at *D.* In the load diagram, the ordinate represents the load intensity and the abscissa represents the position of load on the beam.

Hence the area of the triangle represents the total load and the centroid of the triangle represents the centre of gravity of the load. Thus, total load in this case is $\frac{1}{2} \times 3 \times 20 = 30$ kN and the centre of gravity of this loading is at $\frac{1}{2} \times 3 = 1$ m from *D*, *i.e.*, $1 + 3 - 1 = 3$ m from *A*. For 3 finding the reactions, we can assume that the given load is equivalent to 30 kN acting at 3 m from *A*.

General Loadings: Figure 9.15 shows a general loading. Here the ordinate represents the intensity of loading and abscissa represents position of the load on the beam. For simplicity in analysis such loadings are replaced by a set of equivalent concentrated loads.

Fig. 9.15

External Moment: A beam may be subjected to external moment at certain points. In Fig. 9.16, the beam is subjected to clockwise moment of 30 kN-m at a distance of 2 m from the left support.

In this chapter the beams subjected to concentrated loads, *udl* and external moments are dealt with.

9.5 REACTIONS FROM SUPPORTS OF BEAMS

A beam is in equilibrium under the action of the loads and the reactions. Hence the equilibrium may be written for the system of forces consisting of reactions and the loads. Solutions of these equations give the unknown reactions.

Example 9.1. The beam AB of span 12 m shown in Fig. 9.17 (a) is hinged at A and is on rollers at B. Determine the reactions at A and B for the loading shown in the Figure.

Solution: At A the reaction can be in any direction. Let this reaction be represented by its components V_A and H_A as shown in Fig. 9.17 (*b*). At B the reaction is in vertical direction only. The beam is in equilibrium under the action of system of forces shown in Fig. 9.17 (*b*).

[**Note:** For finding moments, inclined loads are split into their vertical and horizontal components. Horizontal components do not produce moment about *A.*]

Solution: The reaction at *B* will be at right angles to the support, *i.e.,* at 60° to horizontal as shown in the figure. Let the components of the reactions at *A* be H_A and V_A . Then

$$
\frac{1}{x} + \frac{1}{x} = \frac{1}{x
$$

i.e., $\alpha = 79.45^{\circ}$, as shown in Fig. 9.18(*b*).

*Example 9.3***.** *Find the reactions at supports A and B of the loaded beam shown in Fig. 9.19(a).*

Fig. 9.19

Solution: The reaction at *A* is vertical. Let H_B and V_B be the components of the reaction at *B*. $\sum M_B = 0$, gives $R_A \times 9 - 20 \times 7 - 30 \times 4 \times 5 - 60 \sin 45^\circ \times 2 = 0$ ∴ $R_A = 91.6503$ kN. $\sum H_A = 0$, gives $H_B - 60 \cos 45^\circ = 0$ ∴ *H_B* = 42.4264 kN. $\sum V_A = 0$ $91.6503 + V_B - 20 - 30 \times 4 - 60 \sin 45^\circ = 0$ V_B = 90.7761 kN. ∴ $R_B = \sqrt{42.4264^2 + 90.7761^2}$ ∴ $R_B = 100.2013$ kN. $\alpha = \tan^{-1} \frac{90}{42}$. $_{-1}$ 90.7761 42 4264 $\alpha = 64.95^{\circ}$, as shown in Fig. 9.19(*b*). Consists $H_B = 60$ and $85^\circ = 0$
 $\therefore H_B = 42.4264$ kN.
 $\therefore H_B = 42.4264$ kN.
 $\therefore H_B = 42.4264$ kN.
 $\therefore H_B = 90.7761$ kN.
 $\therefore R_B = 142.4264 + 90.776F$
 $\therefore R_B = 100.2013$ kN.
 $\therefore R_B = 100.2013$ kN.
 $\therefore R_B = 100.2013$ kN.
 \there

Example 9.4. The cantilever shown in Fig. 9.20 is fixed at A and is free at B. Determine the reactions when it is loaded as shown in the Figure.

Fig. 9.20

Solution: Let the reactions at *A* be H_A , V_A and M_A as shown in the figure

Now $\Sigma H = 0$, gives $H_A = 0.$ $\Sigma V = 0$, gives $V_A - 16 \times 2 - 20 - 12 - 10 = 0$ $V_A = 74$ kN. $\Sigma M = 0$, gives $M_A - 16 \times 2 \times 1 - 20 \times 2 - 12 \times 3 - 10 \times 4 = 0$ ∴ *M_A* = 148 kN-m.

Example 9.5. Compute the reaction developed at support in the cantilever beam shown in Fig. 9.21. **Solution:** Let the vertical reaction be V_A and moment be M_A . There is no horizontal component of reactions, since no load is having horizontal component

$$
\Sigma V = 0, \text{ gives}
$$

$$
V_A - 20 \times 2 - 15 \times 10 = 0
$$

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Solution: As supports *A* and *B* are simple supports and loading is only in vertical direction, the reactions R_A and R_B are in vertical directions only.

$$
\Sigma M_A = 0, \text{ gives}
$$

\n $R_B \times 5 - 30 \times 1 - 20 \times 3 \times (2 + 1.5) - 40 \times 6.5 = 0$
\n $R_B = 100 \text{ kN.}$
\n $\Sigma V = 0, \text{ gives}$
\n $R_A + R_B - 30 - 20 \times 3 - 40 = 0$
\n \therefore
\n $R_A = 130 - R_B = 130 - 100 = 30 \text{ kN.}$

Example 9.7. Find the reactions at supports A and B of the beam shown in Fig. 9.23.

Fig. 9.23

Solution: Let V_A and H_A be the vertical and the horizontal reactions at *A* and V_B be vertical reaction at *B*.

> $\Sigma H = 0$, gives $H_A = 0$.

$$
\Sigma M_A = 0, \text{ gives}
$$

\n
$$
-20 \times 2 \times 1 + 60 \times 4 + 30 + 20 \times 11 - V_B \times 9 = 0
$$

\n
$$
\Sigma V = 50 \text{ kN.}
$$

\n
$$
\Sigma V = 0, \text{ gives}
$$

\n
$$
-20 \times 2 + V_A - 60 + V_B - 20 = 0
$$

\n
$$
V_A = 120 - V_B = 120 - 50
$$

\n
$$
V_A = 70 \text{ kN.}
$$

Example 9.8. Determine the reactions at A and B of the overhanging beam shown in Fig. 9.24(a).

(Negative sign show that the assumed direction of V_A is wrong. In other words, V_A is acting vertically downwards).

$$
R_{A} = \sqrt{V_{A}^{2} + H_{A}^{2}}
$$

\n
$$
R_{A} = 23.3666
$$
 kN.
\n
$$
\alpha = \tan^{-1} \frac{V_{A}}{H_{A}}
$$

\n
$$
\alpha = 24.79^{\circ}, \text{ as shown in Fig. 9.24(b)}.
$$

Example 9.9. A beam AB 20 m long supported on two intermediate supports 12 m apart, carries a uniformly distributed load of 6 kN/m and two concentrated loads of 30 kN at left end A and 50 kN at the right end B as shown in Fig. 9.25. How far away should the first support C be located from the end A so that the reactions at both the supports are equal ?

Solution: Let the left support *C* be at a distance *x* metres from *A*.

Now,
\n
$$
R_C = R_D
$$
 (given)
\n $\Sigma V = 0$, gives
\n $R_C + R_D - 30 - 6 \times 20 - 50 = 0$

or $2R_C = 30 + 120 + 50$ since $R_C = R_D$ or $R_C = 100 \text{ kN}$ $R_D = 100 \text{ kN}$ $\sum M_A = 0$, gives $100x + 100 (12 + x) - 6 \times 20 \times 10 - 50 \times 20 = 0$ $200x = 1000$ $x = 5$ m. R_D 20 m x 12 m 30 kN k N $/m$ A¹ 50 kN C _D≰ D B **Fig. 9.25** $\sum M_A = 0$, gives
 $100x + 100 (12 + x) - 6 \times 20 \times 10 - 50 \times 20 = 0$
 $200x = 1000$
 $x = 5$ m.
 $x = 6$ m.
 x

9.6 SHEAR FORCE AND BENDING MOMENT

The load applied on a beam gets transferred to supports. To see how this transfer takes place, consider a simply supported beam subject to the loads as shown in Fig. 9.26.

$$
\sum M_B = 0, \text{ gives}
$$

\n $R_A \times 7 = 20 \times 5 + 40 \times 3 + 60 \times 1$
\n $R_A = 40 \text{ kN}$
\n $R_B = (20 + 40 + 60) - 40$
\n $= 80 \text{ kN}$
\n 20 kN
\n 40 kN
\n60 kN
\nA
\nFig. 9.26

Now to find what is happening at a section, consider the section at *C* which is at a distance of 3 m from *A*. Imagining a cut at this section, left hand side portion and right hand side portions are drawn separately in Fig. 9.27.

Considering the algebraic sum of vertical forces acting on left hand side portion, it is found that a net vertical force of $40 - 20 = 20$ kN is experienced at the section. This effect is shown in Fig. 9.27 by dotted line. Again considering the portion on left hand side, the moment experienced at *C* is given by

 $M_C = 40 \times 3 - 20 \times 1$

 $= 100$ kN-m clockwise.

This moment is also shown on left hand side portion of the beam at C by dotted line. Now, considering the right hand side portion:

Force experienced at *C*

 $= 80 - 60 - 40 = -20$ kN

$$
= 20
$$
 kN. downward

and the moment experienced is $M = 80 \times 4 - 60 \times 3 - 40 \times 1$

= 100 kN-m. anticlockwise

These forces and moments are also shown in Fig. 9.27 on right hand side portion of the beam at *C*.

Thus the section *C* is subjected to a force of 20 kN, which is trying to shear off the beam as shown in Fig. $9.28(a)$, and is also subjected to a moment of 100 kN-m which is trying to bend the beam as shown in Fig. 9.28(*b*). Since this force of 20 kN is trying to shear off the section, it is called as shear force at section *C*. The moment is trying to bend the beam at *C* and hence it is called as bending moment at that section. The shear force and bending moment at a section in a beam may be defined as follows:

Fig. 9.28

"*Shear Force* at a section in a beam (or any structural member) is the force that is trying to shear off the section and is obtained as algebraic sum of all the forces acting normal to the axis of beam either to the left or to the right of the section".

"*Bending Moment* at a section in a beam is the moment that is trying to bend the beam and is obtained as algebraic sum of moment of all the forces about the section, acting either to the left or to the right of the section".

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Hence to find shear force or bending moment at a section, a cut at the section is to be imagined and any one portion with all the forces acting on that portion, is to be considered. It may be noted that for finding bending moment at a section, the moment of the forces are to be found about the section considered.

9.7 SIGN CONVENTION

Although different sign conventions may be used, most of the engineers use the following sign conventions for shear forces and bending moment.

(*a*) The shear force that tends to move left portion upward relative to the right portion shall be called as positive shear force (Fig. 9.29).

(*b*) The bending moment that is trying to sag the beam shall be taken as positive bending moment. If left portion is considered positive bending moment comes out to be clockwise moment (Fig. 9.30).

To decide the sign of moment due to a force about a section, assume the beam is held tightly at that section and observe the deflected shape. Then looking at the shape sign can be assigned. Thus in the problem shown in Fig. 9.26 and 9.27, 40 kN reaction at *A* produces positive moment at *C* and 20 kN load produces negative moment.

9.8 RELATIONSHIP BETWEEN LOAD INTENSITY, SHEAR FORCE AND BENDING MOMENT

Consider the beam *AB* subject to a general loading as shown in Fig. 9.31(*a*). The free body diagram of a segment of beam at a distance *x* from *A* and of length δ*x* is shown in Fig. 9.31(*b*). The intensity of loading on this elemental length may be taken as constant. Let it be *w*/unit length.

Let shear force and bending moment acting on the section at a distance *x* be *F* and *M* respectively. At section at a distance $x + \delta x$, these values be $F + \delta F$ and $M + \delta M$ respectively. Now from the equilibrium of the element.

9.9 SHEAR FORCE AND BENDING MOMENT DIAGRAMS

Shear force and bending moment in a beam vary from section to section. The graphical representation of shear force in which ordinate represents shear force and the abscissa represents the position of the section is called *Shear Force Diagram* (*SFD*)*.* The diagram in which the ordinate represent bending moment the abscissa represent the position of the section is called *Bending Moment Diagram (BMD).* In drawing *SFD* and *BMD*, the sign conventions explained earlier are used. These diagrams are usually located below the load diagram.

From equations 1 and 2, it may be concluded that the rate of change of shear force (slope of shear force diagram curve) at any section is equal to the intensity of loading at that section and the rate

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of change of bending moment (*i.e.,* shape of bending moment diagram curve) is equal to the shear force at that section. From equation 2, it can also be concluded that the bending moment will be maximum/minimum where shear force (*dM/dx*) is zero.

At any section, if moment changes its sign the point representing that section is called the point of contraflexure. Obviously, the moment at that section is zero.

9. 10 SFD AND BMD FOR A FEW STANDARD CASES

The methods of drawing shear force and bending moment diagrams have been explained here in case of the following beams subjected to standard loading conditions.

- (*a*) Cantilever beams
- (*b*) Simply supported beams and
- (*c*) Overhanging beams

9.10.1 Cantilever Subject to a Concentrated Load at Free End

Consider the section *X–X* at a distance *x* from free end in a cantilever beam shown in Fig. 9.32(*a*).

From left hand side segment of beam,

 $F = -W$

Thus shear force is constant i.e., it will not vary with *x*. Hence the *SFD* is as shown in Fig. 9.32(*b*).

A
\nof contralexure. Obviously, the moment at that section is zero.
\n9.10 SFD AND BMD FOR A FEW STANDARD CASES
\nThe methods of drawing shear force and bending moment diagrams have been explained the
\nof the following beams subjected to standard loading conditions.
\n(a) Canilever beams
\n(b) Simply supported beams and
\n(c) Overhanging beams
\n9.10.1 Cantilever Subject to a concentrated Load at Free End
\nConsider the section X-X at a distance x from free end in a cantilever beam
\nFig. 9.32(a).
\nFrom left hand side segment of beam,
\n
$$
F = -W
$$

\nThus shear force is constant i.e., it will not vary with x. Hence the SFD is as
\nFig. 9.32(b).
\n $M = -Wx$, linear variation.
\nAt
\n $x = 0$, $M_A = 0$
\n $M_A = 0$
\n $W = -Wx$, linear variation.
\nAt
\n $x = 0$, $M_A = 0$
\n $W = 0$
\

 $\alpha = l, \quad M_B = -Wl$

Hence *BMD* is as shown in Fig. 9.32 (c) .

9.10.2 A Cantilever Subject to UDL Over its Entire Span

Consider the beam shown in Fig. 9.33(*a*).

Considering the left hand side portion of the beam from the section *X–X* which is at a distance *x* from the free end *A,*

 $F = -Wx$, linear variation At $x = 0, F_A = 0$

Fig. 9.33

Hence *SFD* is as shown in Fig. 9.33(b).

$$
M=-wx\frac{x}{2}=-\frac{wx^2}{2}
$$

This is parabolic variation. As magnitude increases at a faster rate with *x*, it is concave parabola as shown in Fig. $9.33(c)$, with extreme values as given below:

At
$$
x = 0
$$
, $M_A = 0$,
At $x = l$, $M_B = -\frac{wl^2}{2}$

9.10.3 Simply Supported Beam Subjected to a Concentrated Load

Let *W* be the concentrated load acting on beam *AB* at a distance '*a*' from the end *A* as shown in Fig. 9.34(*a*).

Now
$$
R_A = \frac{Wb}{l}
$$
 and $R_B = \frac{Wa}{l}$

Consider the portion *AC.* At any distance *x* from *A*,

$$
F = R_A = \frac{Wb}{l}, \text{ constant}
$$

\nAt
\n
$$
M = R_A x = \frac{Wb}{l}x, \text{ linear variation.}
$$

\nAt
\n
$$
x = 0, M_A = 0
$$

\nAt
\n
$$
x = a, M_C = \frac{Wab}{l}
$$

For portion *AC SFD* and *BMD* can be drawn. Consider portion *CB.* The expression derived for portion *AC* will not hold good for this portion. Taking a section at a distance *x* from *B* and considering the right hand side segment of the beam, $F = -R_B = -\frac{Wa}{l}$, constant $M = R_B x = \frac{Wa}{l} x$, linear variation. At $x = 0$, $M_B = 0$ At $x = b$, $M_B = \frac{Wab}{l}$ W $A \rightarrow B$ a _bC b $F = -R_R = \frac{M}{I}$, constant
 $M = R_B x = \frac{W_0}{I} x$, linear variation.

At $x = 0$, $M_B = 0$

At $x = b$, $M_B = \frac{W_0 b}{I}$
 $\frac{W_0}{B} = \frac{W_0}{I}$

(a) Lead Dangman

(a) Lead Dangman

(a) Lead Dangman

(b) SFD

(a) User Civil (a) Use

(a) Load Diagram

<u>w a</u> *l*

l

Fig. 9.34

SF and *BM* diagrams for this portion can now be drawn. Hence *SFD* and *BMD* for the beam is as shown in Fig. $9.34(b)$ and $9.34(c)$ respectively.

Particular case:

When $a = b = l/2$

$$
F = \frac{Wb}{l} = \frac{W(l/2)}{l} = \frac{W}{2}
$$

and moment under the load (centre of span since $a = b = l/2$)

$$
M = \frac{Wab}{l} = \frac{W\,l/2 \times l/2}{l} = \frac{Wl}{4}
$$

9.10.4 A Simply Supported Beam Subjected to UDL

Let the beam *AB* of span *l* be subjected to uniformly distributed load of intensity w/unit length as shown in Fig. 9.35(*a*).

$$
R_A = \frac{wl \cdot l/2}{l} = \frac{wl}{2}
$$

$$
R_B = \frac{wl}{2}
$$

At a section *X–X* which is at a distance *x* from *A*,

 $F = R_A - wx = \frac{wl}{2} - wx$, linear variation

At
$$
x = 0
$$
, $F_A = \frac{wl}{2}$
At $x = l$, $F_B = -\frac{wl}{2}$

∴ *SFD* is as shown in Fig. 9.35(*b*).

Shear force is zero at *x* given by

$$
0 = \frac{wl}{2} - wx
$$

$$
x = \frac{l}{2}
$$

∴ Maximum moment occurs at this points.

At section *X–X*

$$
M = R_A x - wx \cdot \frac{x}{2} = \frac{wl}{2}x - \frac{wx^2}{2}
$$
, parabolic variation.

As *x* increases rate of reduction in the value of *M* is faster. Hence it is convex parabola.

At
\nAt
\n
$$
x = 0
$$
, $M_A = 0$
\nAt
\n $x = l$, $M_B = 0$
\nMaximum moment occurs at $x = \frac{l}{2}$ where shear force $(i.e., \frac{dM}{dx}) = 0$
\n
$$
M_{\text{max}} = \frac{wl}{2} \cdot l/2 - \frac{w(l/2)^2}{2} = \frac{wl^2}{8}
$$

Hence *BMD* is as shown in Fig. 9.35(*c*).

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9.10.5 Overhanging Beam Subjected to a Concentrated Load at Free End

Consider the overhanging beam *ABC* of span $AB = l$ and overhang $BC = a$, subjected to a concentrated load *W* at free end as shown in Fig. 9.36(*a*).

SFD and *BMD* for the entire beam is shown in Fig. 9.36(*b*) and 9.36(*c*) respectively.

Example 9.10. Draw shear force and bending moment diagram for the cantilever beam shown in Fig. 9.37(a).

Solution: Portion *AB*: At distance *x*, from *A*, $F = -20 - 20 x$, linear variation. At $x = 0$, $F_A = -20$ kN At $x = 1$, $F_B = -20 - 20 \times 1 = -40$ kN. $M = -20x - 20x \cdot \frac{x}{2}$, parabolic variation At $x = 0$, $M_A = 0$ At $x = 1$ m, $M_B = -20 - 20 \times 1 \times \frac{1}{2} = -30$ kN-m. Portion *BC*: Measuring *x* from *A*, $F = -20 - 40 - 20x$, linear variation. At $x = 1$ m, $F_B = -80$ kN At $x = 3$ m, $F_C = -120$ kN. $M = -20x - 40(x - 1) -20x$. $\frac{\pi}{2}$, parabolic variation; At *x* = 1 m, *M* = –30 kN-m At $x = 3$ m, $M = -60 - 40 \times 2 - 20 \times 3 \times \frac{3}{2}$ 2 $= -230$ kN-m Hence *SFD* and *BMD* are shown in Fig. 9.37(*b*) and 9.37(*c*) respectively. $A \perp m$ BL $2m$ SC (a) Load D iagram 1 m BL 2 m 20 kN/m B 20 kN 40 kN 40*///////*
///////- ve 120 80 k 20 (b) SFD 230 30 (c) BMD At $x = 0$. $F_A = -20$ and $F_A = -20$ and $F_A = -20$ and $F_B = -20 - 20 \times 1 = -40$ kN.
 $M = -20x - 20x + \frac{x}{2}$, parabolic variation
 $x = 0$. $M_A = 0$

At $x = 1$ m, $M_B = -20 - 20 \times 1 \times \frac{1}{2} = -30$ kNem.

Purtion *BC*:

Messuring *x* from A

Example 9.11. Draw the SF and BM diagrams for the beam shown in Fig. 9.38(a) and find out the position and the magnitude of maximum moment.

Solution: $\sum M_A = 0 \rightarrow$ $R_B \times 10 = 20 \times 5 \times 2.5 + 20 \times 5 + 40 \times 7.5 + 20 \times 8.5$ ∴ $R_B = 82$ kN. $\Sigma V = 0 \rightarrow$ ∴ $R_A = 20 \times 5 + 20 + 40 + 20 - 82 = 98$ kN. Portion *AC:* Measuring *x* from *A*, $F = 98 - 20x$, linear variation At $x = 0, F_A = 98$ kN At $x = 5$ m, $F_B = 98 - 100 = -2$ kN Points where shear force is zero is given by, $0 = 98 - 20x$ or $x = 4.9 \text{ m}$ (c) BMD A C D E (a) Load D iagram B 20 kN/m 20 kN 40 kN 20 kN 20 kN/m 2 98 (b) SFD 4.9 m $\frac{1}{5}$ m $\frac{1}{2.5}$ m $\frac{1}{1}$ m $\frac{1}{1.5}$ m 22 $62 \frac{22}{82} \frac{1}{100}$ – ve 3 240.1 240 185 123 + ve **Fig. 9.38** Eq. 10 = 20 × 5 × 2.5 + 20 × 5 + 40 × 7.5 + 20 × 8.5
 $R_9 \times 10 = 20 \times 5 + 20 + 40 + 20 - 82 = 98 \text{ kN}$.
 $\sum V = 0 \rightarrow$
 $R_A = 20 \times 5 + 20 + 40 + 20 - 82 = 98 \text{ kN}$.

Portion AC:

Measuring x from A,
 $K = 98 - 20x$, linear variation

A

Moment is given by

At
\n
$$
M = 98x - 20x \frac{x}{2}, \text{ parabolic variation.}
$$
\n
$$
x = 0, \quad M_A = 0
$$
\nAt
\n
$$
x = 5 \text{ m}, \quad M_B = 98 \times 5 - \frac{20 \times 5^2}{2} = 240 \text{ kN-m}
$$

Maximum moment occurs at $x = 4.9$ m where $F = \frac{dM}{dx} = 0$ $M_{\text{max}} = 98 \times 4.9 - 20 \times \frac{4.9}{2}$ $\frac{.9^2}{2}$ = 240.1 kN-m Portion *CD:* Measuring *x* from *B* and considering right hand side segment of the beam, $F = -82 + 20 + 40 = -22$ kN, constant $M = 82x - 20(x - 1.5) - 40(x - 2.5)$ $= 22x + 130$, linear variation. At $x = 2.5$ m, $M_D = 22 \times 2.5 + 130 = 185$ kN-m At $x = 5$ m, $M_C = 22 \times 5 + 130 = 240$ kN-m. Portion *DE:* Measuring *x* from *B* and considering the portion of the beam on the right side of the section, $F = -82 + 20 = -62$ kN, constant $M = 82x - 20(x - 1.5)$ linear variation At $x = 1.5$ m, $M_F = 82 \times 1.5 = 123$ kN-m. At $x = 2.5$, $M = 82 \times 2.5 - 20 \times 1 = 185$ kN-m. Portion *EB:* Measuring x from B and considering the right side segment, $F = -82$ kN, constant *M* = 82*x,* linear variation. At $x = 0, M_B = 0$ At $x = 1.5$ m, $M_E = 82 \times 1.5 = 123$ kN-m. *SFD* and *BMD* are shown in Fig. 9.38(*b*) and 9.38(*c*) respectively, for the entire beam. *Example 9.12. A beam of span 8 m has roller support at A and hinge support at B as shown in Fig. 9.39(a). Draw SF and BM diagrams when the beam is subjected to udl, a concentrated load and an externally applied moment as shown in the Figure. Solution:* $\sum M_A = 0 \rightarrow$ Portion CD:

Weave $m_{\text{max}} = 26 \times 4.5t = 2t \times 20 + 40 = -22$ kN-1

Neasuring *x* from *B* and considering right hand side segment of the beam,
 $F = 82x + 20k + 130 = -22 \times 140$, constant
 $x = 22x + 130 = 188$ kN-m

At
 $x = 25$ m,

 $R_B \times 8 - 10 \times 4 \times 2 - 20 \times 4 + 240 = 0$ $R_B = -10$ kN (upwards) = 10 kN. (downwards) $\overline{\Sigma}V = 0 \rightarrow$ $R_A = 10 \times 4 + 20 + 10 = 70$ kN.

Portion *AC:*

Measuring *x* from *A* and considering left hand side segment of the beam,

At
\nAt
\n
$$
F = 70 - 10x
$$
, linear variation
\n $x = 0$, $F_A = 70$ kN
\nAt
\n $x = 4$ m, $F_C = 70 - 40 = 30$ kN

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SFD and *BMD* are shown in Fig. 9.39(*b*) and 9.39(*c*).

Note: The bending moment value will be the same at a point, whether calculated by considering left hand side or the right hand side segment of the beam, if there is no external moment acting at that point. If the external moment is acting at a point, there will be sudden change in *BMD* at that point to the extent equal to the magnitude of external bending moment.

Example 9.13. The overhanging beam ABC is supported at A and B, the span AB being 6 m. The overhang BC is 2 m (see Fig. 3.42 a). It carries a uniformly distributed load of 30 kN/m over a length of 3 m from A and concentrated load of 20 kN at free end. Draw SF and BM diagrams.

$$
\text{BEAMS} \qquad \qquad 305
$$

In this portion the bending moment changes the sign. The point of contraflexure is given by the expression

 $0 = -20x + 49.167 (x - 2)$ *i.e.*, $x = 3.371$ m from free end.

Portion *BC:*

Measuring *x* from free end,

Hence *SF* and *BM* diagrams are as shown in Fig. 9.40(*b*) and 9.40(*c*) respectively. *Example 9.14. Draw BM and SF diagrams for the beam shown in Fig. 9.41(a), indicating the values at all salient points.*

Fig. 9.41

 $R_E = 68.125$ kN. $\Sigma V = 0 \rightarrow$

i.e., the

$$
R_B = 20 + 30 \times 2 + 40 + 25 \times 1 - 68.125
$$

\n
$$
= 76.875 \text{ kN.}
$$

\n
$$
M = -20 \text{ kN, constant}
$$

\n
$$
M = -20 \text{ kN, constant}
$$

\n
$$
M = 0
$$

\nAt
\n
$$
M = 0
$$

\nAt
\n
$$
x = 0
$$
, $H = 4$
\n
$$
M = 20 \text{ kN-m}
$$

\n
$$
= 20 + 76.875 - 30x
$$
, linear variation
\nAt
\n
$$
x = 0
$$
, $F = 56.875 \text{ kN.}$
\nAt
\n
$$
x = 2 \text{ m}, F = -3.125 \text{ kN-m.}
$$

\nThe point of zero shear force is given by
\n
$$
M = -20(x + 1) + 76.875x - 30x
$$

\n
$$
x = 1.896 \text{ m from } B.
$$

\nAt distance *x* from *B* the moment is given by
\n
$$
M = -20(x + 1) + 76.875x - 30x
$$

\n
$$
x = 2 \text{ m}, M_0 = +20 \text{ kN-m.}
$$

\nAt
\n
$$
x = 0
$$
, $M_B = +20 \text{ kN-m.}$
\nAt
\n
$$
x = 2 \text{ m}, M_0 = +20 \text{ kN-m.}
$$

\nAt
\n
$$
x = 2 \text{ m}, M_0 = +20 \text{ kN-m.}
$$

\nMaximum moment occurs where $SF = 0$. *i.e.*, at *x* = 1.896 m.
\n
$$
\therefore M_{\text{max}} = -20 + 56.875 \times 1.896 - 15 \times 1.896^2
$$

\nThe bending moment is changing its sign in this portion. Hence the point of contrall
\n
$$
M = 23.913 \text{ kN-m.}
$$

\nThe second for contrall
\n
$$
M = 23.
$$

 Mea Shear force = $25 \times 1 - 68.125 - 40$ $= -3.125$ kN, constant *M* = $-25 \times 1 \times (x - 0.5) + 68.125 (x - 1) - 40 (x - 2)$, linear variation At $x = 3$ m, $M_C = 33.75$ kN-m. At $x = 2$ m, $M_D = -25 \times 1.5 + 68.125$ $= 30.625$ kN-m.

$$
\text{BEAMS} \qquad \qquad \textbf{307}
$$

SF and *BM* diagrams are as shown in Fig. 9.41(*b*) and 9.41(*c*) respectively.

9.11 SHORT-CUT PROCEDURE

To save time, one can calculate the values only at salient points and draw shear force and bending moment diagrams, noting the following points:

(*a*) Wherever concentrated load is acting (including support reaction), there will be sudden change in *SF* value to the extent of that force at that point, in the direction of that force (while proceeding from left to right).

(*b*) If an external moment is acting at a point on the beam, there will be sudden change in the value of bending moment to the extent of that external bending moment. While proceeding from left, it will be increase in value, if the moment is clockwise. There will be drop in bending moment value if the external bending moment is anticlockwise. Exactly opposite phenomenon will be observed while proceeding from right to left.

(*c*) Since $\frac{dF}{dx} = w$ and $\frac{dM}{dx}$ = *F*, the shear force curve will be one degree higher than that of load

curve and bending moment curve will be one degree higher than that of shear force curve. The nature of variation, for various loads is listed in Table 9.1.

Load	SF	BM
No load	Constant	Linear
Uniformly distributed load	Linear	Parabolic
Uniformly varying load	Parabolic	Cubic

Table 9.1 Nature of SF and BM Variation

The following points also may be noted while drawing shear force and bending moment diagrams:

(*a*) The bending moment is maximum when shear force is zero. The location and the value of maximum moment should always be indicated in bending moment diagram.

(*b*) The point of contraflexure is an important point in a *BMD*, hence if it exists, its location should be indicated.

This method of drawing *SFD* and *BMD is illustrated with two problems below: Example 9.15. Determine the reactions and construct the shear force and bending moment diagrams for the beam shown in Fig. 9.42(a). Mark the salient points and the values at those points. Solution:* $\sum M_A = 0 \rightarrow$

 $R_B \times 6 + 120 - 60 \times 4 - 60 \times 7 = 0$ ∴ $R_B = 90$ kN. $\Sigma V = 0 \rightarrow R_A = 60 + 60 - 90 = 30$ kN. *SFD:* In portion ADE , $F = 30$ kN. In portion *EB*, $F = 30 - 60 = -30$ kN. In portion *BC*, $F = 60$ kN. *SFD* is as shown in Fig. 3.45(*b*). *BMD:* At *A*, $M_A = 0$ Just to the left of *D*, $M = 30 \times 2 = 60$ kN-m. Just to the right of *D*, $M = 30 \times 2 - 120 = -60$ kN-m. At *E*, $M = -60 \times 3 + 90 \times 2 = 0$ At *B*, $M = -60 \times 1 = -60$ kN-m. At *C*, $M = 0$ *BMD* is as shown in Fig. $9.42(c)$. Show the musical content of the reactions are continued with two problems below.

Example 9.15. Determine the reactions and construct the shear force and beneding moment

Example 9.15. Determine the reactions and the valu

Fig. 9.42

Example 9.16. Draw the bending moment and shear force diagram for the beam loaded as shown in Fig. 9.43. (a). Mark the values at the salient points Determine the point of contraflexure also.

At *D*, calculating from right hand side $M_D = -50 + 38 \times 2 = 26$ kN-m At *E*, $M_E = -50 \text{ kN-m}$ Point of contraflexure is between *D* and *E*. Let it be at distance *x* from support *E*. Then $-50 + 38x = 0$ \therefore *x* = 1.32 m from E as shown in Fig. 9.43 (c) BMD is as shown in Fig. 9.43 (*c*) Control is as shown in Fig. 9.43 (c)

EMD is as shown in Fig. 9.43 (c)

EMD is as shown in Fig. 9.43 (c)

EMPORTANT FORMULAE

1. Shear force at a section of the beam = Σ and forces to the left or right of the section.

IMPORTANT FORMULAE

- **1.** Shear force at a section of the beam = Σ all forces to the left or right of the section.
- **2.** Bending moment at a section of the beam = Σ moments of all forces to the left or to the right of the section, moment centre being the section.

3.
$$
\frac{dF}{dx} = w \text{ and } \frac{dM}{dx} = -F
$$

THEORY QUESTIONS

- **1.** Explain the following types of supports to beams: (*i*) Simple support (*ii*) Hinged support
	- (*iii*) Fixed support
- **2.** Bring out the difference between statically determinate and indeterminate beams.
- **3.** Derive the relationship between, (*i*) Shear force and load intensity (*ii*) Bending moment and shear force

PROBLEMS FOR EXERCISE

1. Determine the reaction at A and B on the overhanging beam shown in Fig. 9.44.

Fig. 9.44

[Ans. $R_A = 23.3666 \text{ kN}; \text{ a } = 24.79^{\circ}; R_B = 71.011 \text{ kN}$] **2.** An overhanging beam is on rollers at *A* and is hinged at *B* and is loaded as shown in Fig. 9.45. Determine the reactions at *A* and *B*.

 $[Ans. R_A = 45 kN; R_B = 15 kN]$

[Ans. $F_B = 4.25$, $F_E = 3.75$, $M_{\text{max}} = 4.5$ kN-m at *D* Point of contraflexure at 1.47 m from *A*] 50 kN

[Ans. $F_B = 70.83$, $F_E = 29.17$, $M_{\text{max}} = 45$ kN-m at *B* Point of contraflexure at 4.313 m from *A*]

A BCD 10 kN/m 10 kN/m 2 m 2 m 2 m E 40 kN 20 kN F 2 m 2 m

Fig. 9.48

[Ans. $R_B = 46.67 \text{ kN}, R_E = 53.33 \text{ kN}, M_B = -20 \text{ kN-m}$ M_E = – 40 kN-m, M_C = 33.33 kN-m, Point of contraflexure : 2.75 m from *A* and 3.57 m from *F*]

[Ans. $R_B = 95$, $R_C = 45$, $M_{\text{max}} = -4.375$ at 4.75 m from *A* No +ve moment anywhere]

[Ans. $R_A = 47$, $R_D = 53$, $M_B = 47$, $M_C = 44$ to 24, $M_D = 30$ Point of contraflexure 2.59 m from *F*]

4. Determine load *P* such that reactions at supports *A* and *B* are equal in the beam shown in Fig. 9.51. Draw *SF* and *BM* diagram marking the values at salient point.

> [Ans. $P = 20$ kN, $R_A = R_B = 50$ kN, $M_{\text{max}} = 62.5$ kN-m at 2.5 m from *A*, Point of contraflexure 1.667 m from *C*]

5. A bar of length '*l*' is supported at *A* and *B* which are at distances '*a*' from the ends as shown in Fig. 9.52. Find the distance '*a*' such that maximum moment is least. Using the above results find the most economical length of railway sleeper if the rails are 1.6 m apart.

Fig. 9.52

[**Hint:** Moment at *B* = Moment at central point *E*] [Ans. $a = 0.2071$ Economical length of railway sleeper $= 2.73$ m]

10 Stresses in Beams

As seen in the last chapter beams are subjected to bending moment and shear forces which vary from section to section. To resist them stresses will develop in the materials of the beam. For the simplicity in analysis, we consider the stresses due to bending and stresses due to shear separately.

Fig. 10.1. Nature of Stresses in Beams

Due to pure bending, beams sag or hog depending upon the nature of bending moment as shown in Fig. 10.1. It can be easily observed that when beams sag, fibres in the bottom side get stretched while fibres on the top side are compressed. In other words, the material of the beam is subjected to tensile stresses in the bottom side and to compressive stresses in the upper side. In case of hogging the nature of bending stress is exactly opposite, *i.e*., tension at top and compression at bottom. Thus bending stress varies from compression at one edge to tension at the other edge. Hence somewhere in between the two edges the bending stress should be zero. The layer of zero stress due to bending is called **neutral layer** and the trace of neutral layer in the cross-section is called **neutral axis** [Refer Fig. 10.1].
10.1 ASSUMPTIONS

Theory of simple bending is developed with the following assumptions which are reasonably acceptable:

- (*i*) The material is homogeneous and isotropic.
- (*ii*) Modulus of elasticity is the same in tension and in compression.
- (*iii*) Stresses are within the elastic limit.
- (*iv*) Plane section remains plane even after deformations.
- (*v*) The beam is initially straight and every layer of it is free to expand or contract.
- (*vi*) The radius of curvature of bent beam is very large compared to depth of the beam.

10.2 BENDING EQUATION

There exists a define relationship among applied moment, bending stresses and bending deformation (radius of curvature). This relationship can be derived in two steps:

- (*i*) Relationship between bending stresses and radius of curvature.
- (*ii*) Relationship between applied bending moment and radius of curvature.
- (*i*) *Relationship between bending stresses and radius of curvature:* Consider an elemental length *AB* of the beam as shown in Fig. 10.2(*a*). Let *EF* be the neutral layer and *CD* the bottom most layer. If *GH* is a layer at distance *y* from neutral layer *EF*, initially $AB = EF = GH =$ *CD.*

 Let after bending *A*, *B*, *C*, *D*, *E*, *F*, *G* and *H* take positions *A*′, *B*′, *C*′, *D*′, *E*′, *F*′, *G*′ and *H*′ respectively as shown in Fig. 10.2(*b*). Let *R* be the radius of curvature and φ be the angle subtended by *C*′*A*′ and *D*′*B*′ at centre of radius of curvature. Then,

$$
EF = E'F', \text{ since } EF \text{ is neutral axis}
$$

= $R\phi$...(i)
Strain in $GH = \frac{\text{Final length} - \text{Initial length}}{\text{Initial length}}$

$$
= \frac{G'H' - GH}{GH}
$$

But $GH = EF$ (The initial length)
and $GH' = (R + y) \phi$
 $= \frac{R}{R} \phi$

∴ Strain in layer GH = $\frac{(R + y) \phi - R\phi}{R\phi}$
 $= \frac{y}{R}$...(ii)
Since strain in GH is due to tensile forces, strain in GH = f/E ...(iii)
where f is tensile stress and E is modulus of elasticity.
From eqns. (ii) and (iii), we get
 $\frac{f}{E} = \frac{y}{R}$...(10.1)
(ii) Relationship between bending moment and radius of curvature: Consider an elemental area
δa at distance y from neutral axis as shown in Fig. 10.3.
From eqn. 10.1, stress on this element is
 $f = \frac{E}{R} y$...(i)
∴ Force on this element
 $= \frac{E}{R} y$ δa
Moment of resistance of this elemental force about neutral axis
 $= \frac{E}{R} y$ δa
 $= \frac{E}{R} y$ δa
 $= \frac{E}{R} y^2 \delta a$

or

(*ii*) *Relationship between bending moment and radius of curvature:* Consider an elemental area δ*a* at distance *y* from neutral axis as shown in Fig. 10.3.

 $=\frac{2}{R}$...(10.1)

Fig. 10.3

From eqn. 10.1, stress on this element is

 $f = \frac{E}{R} y$

E

f y *E*

 $\frac{E}{R}$ *y* ...(*i*)

∴ Force on this element

$$
= \frac{E}{R} y \, \delta a
$$

Moment of resistance of this elemental force about neutral axis

$$
= \frac{E}{R} y \, \delta a y
$$

$$
= \frac{E}{R} y^2 \, \delta a
$$

∴ Total moment resisted by the section *M*′ is given by

$$
M' = \sum \frac{E}{R} y^2 \, \delta a
$$

$$
= \frac{E}{R} \sum y^2 \, \delta a
$$

From the definition of moment of inertia (second moment of area) about centroidal axis, we know $=\frac{E}{R}\sum y^5 \delta a$

From the definition of moment of inertia (second moment of area) about centroidal

know
 $I = \Sigma y^2 \delta a$
 \therefore $M' = \frac{E}{R}I$

From equilibrium condition, $M = M'$ where M is applied moment.
 \therefore $M = \frac{E}{R}$

I = Σ*y*² δ*a* ∴ $M' = \frac{E}{R}$ *I*

From equilibrium condition, $M = M'$ where M is applied moment.

R I

E

or

From eqns. (10.1) and (10.2) , we get

∴ $M = \frac{E}{R}$

M I f y $=\frac{f}{y} = \frac{E}{R}$...(10.3)

 $=\frac{2}{R}$...(10.2)

where $M =$ bending moment at the section

 $I =$ moment of inertia about centroid axis

M I

 $f =$ bending stress

y = distance of the fibre from neutral axis

E = modulus of elasticity and

R = radius of curvature of bent section.

Equation (10.3) is known as bending equation.

10.3 LOCATING NEUTRAL AXIS

Consider an elemental area δ*a* at a distance *y* from neutral axis [Ref. Fig. 10.3].

If '*f*' is the stress on it, force on it = $f \, \delta a$

But
$$
f = \frac{E}{R}
$$
 y, from eqn. (10.1).

∴ Force on the element = $\frac{E}{R}$ *y* δa

Hence total horizontal force on the beam

 $=\sum \frac{E}{R} y \delta a$

 $=\frac{E}{R}$ Σ_{*y*} δ*a*

Since there is no other horizontal force, equilibrium condition of horizontal forces gives

$$
\frac{E}{R} \Sigma y \, \delta a = 0
$$

As $\frac{E}{R}$ *R* is not zero,

$$
\Sigma y \, \delta a = 0 \tag{1}
$$

If *A* is total area of cross-section, from eqn. (*i*), we get

$$
\sum \frac{y \,\delta a}{A} = 0 \qquad \qquad \dots (ii)
$$

Noting that $\Sigma y \delta a$ is the moment of area about neutral axis, $\frac{\Sigma y \delta a}{\Delta}$ $\frac{\partial \delta a}{\partial A}$ should be the distance of centroid of the area from the neutral axis. Hence $\frac{\Sigma y \delta a}{\Lambda}$ $\frac{y\delta a}{A} = 0$ means the *neutral axis coincides with*

the centroid of the cross-section.

10.4 MOMENT CARRYING CAPACITY OF A SECTION

From bending equation, we have

i.e.,
$$
\frac{M}{I} = \frac{f}{y}
$$
...(i)

$$
f = \frac{M}{I}y
$$

Hence bending stress is maximum, when *y* is maximum. In other words, maximum stress occurs in the extreme fibres. Denoting extreme fibre distance from neutral fibre as y_{max} equation (*i*) will be

$$
f_{\max} = \frac{M}{I} y_{\max} \qquad \qquad \dots (ii)
$$

In a design f_{max} is restricted to the permissible stress in the material. If f_{per} is the permissible stress, then from equation (*ii*),

$$
f_{\text{per}} = \frac{M}{I} y_{\text{max}}
$$

$$
M = \frac{I}{y_{\text{max}}} f_{\text{per}}
$$

The moment of inertia *I* and extreme fibre distance from neutral axis y_{max} are the properties of section. Hence *^I y*max is the property of the section of the beam. This term is known as **modulus of section** and is denoted by *Z*. Thus $\frac{E}{R} \Sigma y \delta a = 0$

As $\frac{E}{R}$ is not zero,
 $\Sigma y \delta a = 0$

If A is total area of cross-section, from eqn. (i), we get
 $\sum \frac{y \delta a}{A} = 0$

Noting that $\Sigma y \delta a$ is the moment of area about neutral axis, $\frac{\Sigma y \delta x}{A}$ shoul

$$
Z = \frac{I}{y_{\text{max}}} \tag{10.4}
$$

and $M = f_{per} Z$...(10.5)

Note : If moment of inertia has unit mm⁴ and y_{max} has mm, *Z* has the unit mm³.

The eqn. (10.5) gives permissible maximum moment on the section and is known as **moment carrying capacity of the section**. Since there is definite relation between bending moment and the loading given for a beam it is possible to find the load carrying capacity of the beam by equating maximum moment in the beam to moment carrying capacity of the section. Thus

$$
M_{\text{max}} = f_{\text{per}} Z \tag{10.6}
$$

If permissible stresses in tension and compressions are different for a material, moment carrying capacity in tension and compression should be found separately and equated to maximum values of moment creating tension and compression separately to find the load carrying capacity. The lower of the two values obtained should be reported as the load carrying capacity. Note : If moment of ineria has unit mm³ and y_{max} has nom, Z has the unit mm³.

The eqn. (10.5) airso permissible maximum moment on the section and is known as

earrying capacity of the section. Since there is def

10.5 SECTION MODULI OF STANDARD SECTIONS

 $y_{\text{max}} = d/2$

Section modulus expressions for some of the standard sections are presented below:

- (*i*) **Rectangular section:** Let width be '*b*' and depth be '*d*' as shown in Fig. 10.4.
	- Since *N*-*A* is in the mid depth

$$
I = \frac{1}{12}bd^{3}
$$

\n
$$
Z = \frac{I}{y_{\text{max}}} = \frac{1/12bd^{3}}{d/2}
$$

\ni.e.,
\n
$$
Z = 1/6bd^{2}
$$
...(10.7)

$$
i.e.,
$$

d/2

 $d/2$

¥ - - ∆

b

G

y I
^Imax

N

(*ii*) **Hollow rectangular section.** Figure 10.5 shows a typical hollow rectangular section with symmetric opening. For this,

$$
I = \frac{BD^3}{12} - \frac{bd^3}{12} = \frac{1}{12} (BD^3 - bd^3)
$$

$$
y_{\text{max}} = D/2
$$

$$
Z = \frac{I}{y_{\text{max}}} = \frac{1}{12} \frac{(BD^3 - bd^3)}{D/2}
$$

i.e.
$$
Z = \frac{1}{6} \frac{BD^3 - bd^3}{D} \qquad ...(10.8)
$$

d

4

i.e., $Z = \frac{\pi}{2}$

 $\frac{\pi}{32} d^3$ (*iv*) **Hollow circular tube of uniform section.** Referring to Fig. 10.7,

 \sum_{max} *d* $=\frac{\pi/64}{d/2}$

 $I = \frac{\pi}{64} d^4$

 $y_{\text{max}} = d/2$

∴ $Z = \frac{I}{y_m}$

(*v*) **Triangular section of base width** *b* **and height '***h***'.** Referring to Fig. 10.8,

Example 10.1. A simply supported beam of span 3.0 m has a cross-section 120 mm × 180 mm. If the permissible stress in the material of the beam is 10 N/mm2 , determine

(*i*) *maximum udl it can carry*

(*ii*) *maximum concentrated load at a point 1 m from support it can carry. Neglect moment due to self weight.*

Solution:

Here
$$
b = 120
$$
 mm, $d = 180$ mm, $I = \frac{1}{12} bd^3$, $y_{max} = \frac{d}{2}$
\n \therefore $Z = \frac{1}{6} bd^2$
\n $= \frac{1}{6} \times 120 \times 180^2 = 648000$ mm³
\n $f_{per} = 10 \text{ N/mm}^2$
\n \therefore Moment carrying capacity of the section
\n $= f_{per} \times Z$
\n $= 10 \times 648000$ N-mm
\n(i) Let maximum *udl* beam can carry be w/metre length as shown in Fig. 10.9.
\nIn this case, we know that maximum moment occurs at mid span and is equal to M_{max}
\nEquating it to moment carrying capacity, we get
\n $\frac{w \times 3^2}{8} \times 10^6 = 10 \times 648000$
\n \therefore $w = 5.76$ kN/m.
\n(ii) Concentrated load at distance 1 m from the sup-
\nport be *P* kN. Referring to Fig. 10.10.
\n $M_{max} = \frac{P \times a \times b}{L} = \frac{P \times 1 \times 2}{3}$
\n $= \frac{2P}{3} \times 10^6$ N-mm
\nEquating it to moment carrying capacity, we get
\n $\frac{2P}{3} \times 10^6 = 10 \times 648000$
\n \therefore $P = 9.72$ kN-m.
\nExample 10.2. A circular steel pipe of external diameter 60 mm and thickness 8 mm is
\nsimply supported beam over an effective span of 2 m. If permissible stress in steel is 15
\ndetermine the maximum concentrated load that can be carried by it at mid span.
\nSolution:
\nExtremal diameter $D = 60$ mm
\nThickness

 $=f_{\text{per}} \times Z$ $= 10 \times 648000$ N-mm

(*i*) Let maximum *udl* beam can carry be *w*/metre length as shown in Fig. 10.9.

In this case, we know that maximum moment occurs at mid span and is equal to $M_{\text{max}} = \frac{wL^2}{8}$.

Equating it to moment carrying capacity, we get,

$$
\frac{w \times 3^2}{8} \times 10^6 = 10 \times 648000
$$

$$
\therefore \qquad w = 5.76 \text{ kN/m}.
$$

(*ii*) Concentrated load at distance 1 m from the support be *P* kN. Referring to Fig. 10.10.

$$
M_{\text{max}} = \frac{P \times a \times b}{L} = \frac{P \times 1 \times 2}{3}
$$

$$
= \frac{2P}{3} \text{ kN-m}
$$

$$
= \frac{2P}{3} \times 10^6 \text{ N-mm}
$$

w/m

Fig. 10.10

Equating it to moment carrying capacity, we get

$$
\frac{2P}{3} \times 10^6 = 10 \times 648000
$$

$$
P = 9.72 \text{ kN-m.}
$$

Example 10.2. A circular steel pipe of external diameter 60 mm and thickness 8 mm is used as a simply supported beam over an effective span of 2 m. If permissible stress in steel is 150 N/mm² , determine the maximum concentrated load that can be carried by it at mid span.

Solution:

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60 min	P = ?
(a)	P1g. 10.11
∴ Internal diameter	= 60 - 2 × 8 = 44 mm.
$I = \frac{\pi}{64} (60^4 - 44^4) = 452188 \text{ mm}^4$	
$\frac{y_{max}}{y_{max}} = 30 \text{ mm}.$	
∴	$Z = \frac{I}{y_{max}} = \frac{452188}{30} = 15073 \text{ mm}^3.$
Moment carrying capacity	$M = f_{\text{per}} Z = 150 \times 15073 \text{ N-mm}.$
Let maximum load it can carry be <i>P</i> KN.	
Then maximum moment = $\frac{pE}{4}$	
$= \frac{P. X2}{4} \text{ kN-m}$	
Equating maximum bending moment to moment carrying capacity, we get	
$0.5P \times 10^6 = 150 \times 15073.$	
$P = 4.52 \text{ kN}.$	
<i>Example 10.3: Figure 10.12 (a) shows the cross-section of a cantilever beam of 2.5 m span. Material used is steel for which maximum permissible stress is 150 N/mm². What is the maximum uniformly distributed load this beam can carry?</i>	
<i>Solution:</i> Since it is a symmetric section, centroid is at mid depth.	
$I = MI$ of 3 rectangles about centroid	

Solution: Since it is a symmetric section, centroid is at mid depth.

 $I = MI$ of 3 rectangles about centroid

Fig. 10.12

$$
= \frac{1}{12} \times 180 \times 10^{3} + 180 \times 10 (200 - 5)^{2}
$$

+
$$
\frac{1}{12} \times 10 \times (400 - 20)^{3} + 10 \times (400 - 20) \times 0^{2}
$$

+
$$
\frac{1}{12} \times 180 \times 10^{3} + 180 \times 10 (200 - 5)^{2}
$$

= 182.6467 × 10⁶ mm⁴

 $= 913233$ mm³.

[Note: Moment of above section may be calculated as difference between MI of rectangle of size 180×400 and 170 × 380. *i.e*.,

$$
I = \frac{1}{12} \times 180 \times 400^3 - \frac{1}{12} \times 170 \times 380^3
$$

$$
y_{\text{max}} = 200 \text{ mm.}
$$

 $=\frac{182.6467\times 10}{200}$

[∴] *^Z* = *^I*

∴ Moment carrying capacity

$$
= f_{\text{per}} \times Z
$$

= 180 × 913233
= 136985000 N-mm.

*y*max

If *udl* is *w* kN/m, maximum moment in cantilever

$$
= wL = 2w \text{ kN-mm}
$$

$$
= 2w \times 10^6 \text{ N-mm}
$$

Equating maximum moment to movement carrying capacity of the section, we get

 $2w \times 10^6 = 136985000$

$$
w = 68.49 \text{ kN/m}
$$

Example 10.4. Compare the moment carrying capacity of the section given in example 10.3 with equivalent section of the same area but (*i*) *square section* (*ii*) *rectangular section with depth twice the width and* (*iii*) *a circular section. Solution:* Area of the section $= 180 \times 10 + 380 \times 10 + 180 \times 10$ $= 7400$ mm² (*i*) Square section If '*a*' is the size of the equivalent square section, *a* 2 $\therefore a = 86.023$ mm. Moment of inertia of this section $=\frac{1}{12} \times 86.023 \times 86.023^3$ $= 4563333$ mm⁴ $Z = \frac{I}{y_{\text{max}}} = \frac{4563333}{86.023/2} = 106095.6 \text{ mm}^3$ Moment carrying capacity $= fZ = 150 \times 106095.6$ $= 15.914 \times 10^6$ N-mm ∴ **Moment carrying capacity of I section Moment carrying capacity of I section** $=$ $\frac{136985000}{15.914 \times 10^6}$ 15.914×10^6 = **8.607.** (*ii*) Equivalent rectangular section of depth twice the width. Let *b* be the width ∴ Depth *d* = 2*b*. Equating its area to area of *I*-section, we get $b \times 2b = 7400$ $b = 60.8276$ mm $y_{\text{max}} = d/2 = b = 60.8276$ $M = f \frac{I}{I}$ *y* $b \times 2b$ \sum_{max} 12 *b* $=150\times\frac{1}{12}\times\frac{b\times(2b)}{1}$ 12 $2b)^3$ $= 150 \times \frac{8}{16}$ $\frac{8}{12}$ $b^3 = 150 \times \frac{8}{12}$ 12 $\times 60.8276^{3}$ = 22506193 N-mm. ∴ **Moment carrying capacity of I section Moment carrying capacity of this section** $\frac{136985000}{22506193} = 6.086.$ University and section

(iii) a circular section.

Solution:

Section = 180 x 10 + 380 x 10 + 180 x 10

Area of the section

(f) Square section

If 'a' is the size of the equivalent square section,

Moment of inertia of t

(*iii*) Equivalent circular section.

Let diameter be *d*.

Then,
\n
$$
\frac{\pi d^2}{4} = 7400
$$
\n
$$
I = \frac{\pi}{64} d^4
$$
\n
$$
y_{\text{max}} = d/2
$$
\n
$$
Z = \frac{I}{y_{\text{max}}} = \frac{\pi}{32} d^3.
$$
\n
$$
M = f_{\text{per}} Z = 150 \times \frac{\pi}{32} \times 97.067^3 = 13468024
$$
\n
$$
\therefore \text{ Moment carrying capacity of l section} = \frac{136985000}{13468024} = 10.17.
$$
\n[Note. I section of same area resists more bending moment compared to an equivalent square, rec circular section. Reason is obvious because in l-section most of the area of material is in heavily stream *x* 10 mm *x* 10 mm *rectangular plate on top flange as shown is Fig. 10.13. If permissible at material of 150 N/mm², determined how much concentrated load the beam of this section at centre of 4 m span. Given ends of beam are simply supported.
\nSolution: Area of section A = 240 × 10 + 180 × 8 + 384 × 8 + 180 × 8 = 8352 m
\nSolution: Area of section A = 240 × 10 + 180 × 8 + 384 × 8 + 180 × 8 = 8352 m
\n
$$
= 240 \times 10 \times 405 + 180 \times 8 \times (400 - 4) + 384 \times 8 \times 200 + 180
$$
\n
$$
= 240 \times 10 \times 405 + 180 \times 8 \times (400 - 4) + 384 \times 8 \times 200 + 180
$$*

[**Note.** *I* section of same area resists more bending moment compared to an equivalent square, rectangular or circular section. Reason is obvious because in *I*-section most of the area of material is in heavily stressed zone.]

Example 10.5. A symmetric I-section of size 180 mm × 40 mm, 8 mm thick is strengthened with 240 mm × 10 mm rectangular plate on top flange as shown is Fig. 10.13. If permissible stress in the material is 150 N/mm2 , determine how much concentrated load the beam of this section can carry at centre of 4 m span. Given ends of beam are simply supported.

Solution: Area of section *A*

 $= 240 \times 10 + 180 \times 8 + 384 \times 8 + 180 \times 8 = 8352$ mm²

Let centroid of the section be at a distance *y* from the bottom most fibre. Then *A* \bar{y} = 240 × 10 × 405 + 180 × 8 × (400 – 4) + 384 × 8 × 200 + 180 × 8 × 4

i.e., $8352 \bar{v} = 2162400$ ∴ $\bar{v} = 258.9$ mm $I = \frac{1}{12} \times 240 \times 10^3 + 240 \times 10 (405 - 258.9)^2$ $+\frac{1}{16}$ $\frac{1}{12}$ × 180 × 8³ + 180 × 8 (396 – 258.9)² + $\frac{1}{12} \times 8 \times 384^3 + 8 \times 384 (200 - 258.9)^2$ $+\frac{1}{16}$ $\frac{1}{12}$ × 180 × 8³ + 180 × 8 (4 – 258.9)² $= 220.994 \times 10^6$ mm⁴ ∴ $y_{\text{top}} = 405 - 258.9 = 146.1$ mm $y_{bottom} = 258.9$ mm. ∴ $y_{\text{max}} = 258.9 \text{ mm}$ [∴] *^Z* = *^I y*max $=\frac{220.994 \times 10}{258.9}$ 6 $= 853588.3$ ∴ Moment carrying capacity of the section $= f_{\text{per}} Z = 150 \times 853588.3$ $t = \frac{1}{12} \times 240 \times 10 + 240 \times 10 + 240 \times 10 + 0.05 - 258.9$
 $+ \frac{1}{12} \times 180 \times 8^4 + 180 \times 8 (396 - 258.9)^2$
 $+ \frac{1}{12} \times 180 \times 8^4 + 8 \times 384 (200 - 258.9)^2$
 $+ \frac{1}{12} \times 180 \times 8^3 + 180 \times 8 (4 - 258.9)^2$
 $+ \frac{1}{12} \times 180 \times 8^3 +$

 $= 128038238.7$ N-mm

$$
= 128.038
$$
 kN-m.

Let P kN be the central concentrated load the simply supported beam can carry. Then max bending movement in the beam

$$
= \frac{P \times 4}{4} = P \text{ kN-m}
$$

Equating maximum moment to moment carrying capacity, we get

P **= 128.038 kN.**

Example 10.6. The cross-section of a cast iron beam is as shown in Fig. 10.14(a). The top flange is in compression and bottom flange is in tension. Permissible stress in tension is 30 N/mm2 and its value in compression is 90 N/mm2 . What is the maximum uniformly distributed load the beam can carry over a simply supported span of 5 m?

Solution:

Cross-section area
$$
A = 75 \times 50 + 25 \times 100 + 150 \times 50
$$

= 13750 mm²

Let neutral axis lie at a distance \bar{y} from bottom most fibre. Then

$$
A\overline{y} = 75 \times 50 \times 175 + 25 \times 100 \times 100 + 150 \times 50 \times 25
$$

 $13750 \times \bar{y} = 1093750$

 \bar{y} = 79.54 mm

$$
+\frac{1}{12} \times 25 \times 100^{3} + 25 \times 100 (100 - 79.54)^{2}
$$

$$
+\frac{1}{12} \times 150 \times 50^{3} + 150 \times 50 (25 - 79.54)^{2}
$$

 $= 61.955493 \times 10^6$ mm⁴.

Extreme fibre distances are

$$
y_{\text{bottom}} = \overline{y} = 79.54 \text{ mm.}
$$

 $y_{\text{top}} = 200 - \overline{y} = 200 - 79.54 = 120.46 \text{ mm.}$

Top fibres are in compression. Hence from consideration of compression strength, moment carrying capacity of the beam is given by

$$
M_1 = f_{\text{per}} \text{ in compression} \times \frac{I}{y_{\text{top}}}
$$

= 90 \times \frac{61.955493 \times 10^6}{120.46}
= 46.289178 \times 10^6 \text{ N-mm}
= 46.289178 kN-m.

Bottom fibres are in tension. Hence from consideration of tension, moment carrying capacity of the section is given by

$$
M_2 = f_{\text{per}} \text{ in tension} \times \frac{I}{y_{\text{bottom}}}
$$

$$
= \frac{30 \times 61.955493 \times 10^6}{79.54}
$$

$$
= 21.367674 \times 10^6 \text{ N-mm}
$$

$$
= 21.367674 \text{ kN-m.}
$$

Actual moment carrying capacity is the lower value of the above two values. Hence moment carrying capacity of the section is

$$
= 21.367674
$$
 kN-m.

Maximum moment in a simply supported beam subjected to *udl* of *w*/unit length and span *L* is

$$
=\frac{wL^2}{8}
$$

Equating maximum moment to moment carrying capacity of the section, we get maximum load carrying capacity of the beam as

$$
w \times \frac{5^2}{8} = 21.367674
$$

...

$$
w = 6.838
$$
 kN/m.

Example 10.7. The diameter of a concrete flag post varies from 240 mm at base to 120 mm at top as shown in Fig. 10.15. The height of the post is 10 m. If the post is subjected to a horizontal force of 600 N at top, find the section at which stress is maximum. Find its value also.

Solution: Consider a section *y* metres from top. Diameter at this section is

∴ $f = 600000 \times 32 \frac{y}{\pi [120 + 12y]^3}$

Actual moment carrying capacity is the lower value of the above two values. Hence moment carrying capacity of the section is
21.367674 kN-m.
Maximum moment in a simply supported beam subjected to <i>udl</i> of <i>w</i> /unit length and span <i>L</i> is
$= \frac{wL^2}{8}$
Equating maximum moment in a simply supported beam subject of <i>udl</i> of <i>w</i> /unit length and span <i>L</i> is
$= \frac{wL^2}{8}$
Equating maximum moment on moment carrying capacity of the section, we get maximum load carrying capacity of the beam as
$w \times \frac{5^2}{8} = 21.367674$
$w = 6.838 \text{ kN/m}.$
<i>Example 10.7. The diameter of a concrete flag post varies from 240 mm at base to 120 mm at top as shown in Fig. 10.15. The height of the post is subject to a horizontal force of 600 N at top, find the section at which stress is maximum. Find its value also.</i>
<i>Solution:</i> Consider a section y metres from top. Diameter at this section is
$d = 120 + \frac{y}{10} (240 - 120)$
$d = 120 + 12y$ mm
$d = \frac{120}{64} d^4$
$Z = \frac{I}{d/2} = \frac{\pi}{32} d^3$
$= \frac{\pi}{32} [120 + 12y]^3$
At this section, moment is given by
$M = 6000$ N-m
Eq = 600000 y N-mm.
Eq

Fig. 10.16

b

^D ^d

 $\overline{\mathbf{t}}$

For f' to be maximum, $\frac{df}{dt}$ *dy* $= 0$ 600000 \times 32 [(120 + 12*y*)⁻³ + *y*(- 3) (120 + 12*y*)⁻⁴ \times 12] = 0 *i.e.*, $(120 + 12y)^{-3} = 36 (120 + 12y)^{-4} y$ *i.e.*, $1 = 36 (120 + 12y)^{-1} y$ *i.e*., 120 + 12*y* = 36*y* ∴ $y = 5$ m. *i.e.*, $(120 + 12y)^{-2} = 36 (120 + 12y)^{-4} y$
 i.e., $129 - 159 (120 + 12y)^{-1} y$
 i.e., $120 + 12y = 36y$

Stress at this section *f* is given by
 $f = 60000 \times 32 \times \frac{5}{\pi (120 + 12 \times 5)^{2}}$
 Example 10.8. A circular log of inde

Stress at this section *f* is given by

$$
f = 600000 \times 32 \times \frac{5}{\pi (120 + 12 \times 5)^3}
$$

$$
f = 5.24 \text{ N/mm}^2.
$$

Example 10.8. A circular log of timber has diameter D. Find the dimensions of the strongest rectangular section one can cut from this.

Solution: Let the width and depth of strongest section that can be cut from the log be '*b*' and '*d*' respectively. Then,

 $D^2 = b^2 + d^2$

 $= D^2 - b^2$

or d^2

For rectangular section

$$
I = \frac{1}{12} bd^3
$$

$$
y_{\text{max}} = dl2.
$$

$$
z = \frac{I}{y_{\text{max}}} = \frac{1}{6}
$$

$$
= \frac{1}{6} b (D^2)
$$

The beam is strongest if section modulus is maximum. Hence the condition is

 bd^2

 $-b^2 = \frac{1}{6}$

 $\frac{1}{6}$ (*bD*² – *b*³)

$$
\frac{dz}{db} = 0
$$

i.e.,

$$
\frac{1}{6} [D^2 - 3b^2] = 0
$$

$$
D^2 = 3b^2
$$

or
\n
$$
b = \frac{D}{\sqrt{3}}.
$$
\n
$$
d = \sqrt{(D^2 - b^2)} = \sqrt{D^2 - \frac{D^2}{3}} = D\sqrt{2/3}
$$

Thus the dimensions of strongest beam

$$
= \frac{D}{\sqrt{3}} \text{ wide} \times \sqrt{2/3} D \text{ deep.}
$$

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10.6 PROPORTIONING SECTIONS

In designing beams, span of the beam is known and load expected on the beam can be estimated. Hence bending moment to be resisted by the beam can be calculated. A designer has to select suitable section of the beam of desirable materials. Theoretically speaking, the section required changes along the span. Usually uniform sections are used. Hence the section selected should be capable of resisting the maximum moment. In case of circular sections we may find the diameter required, since section modulus required depends only on diameter. In case of rectangular sections, section modulus depends upon width and depth. Hence usually width is assumed and depth is calculated or else ratio of width to depth is assumed and section is selected. For steel sections, Indian Standard Hand Book may be used to identify standard section that satisfies the required section modulus value. This process of proportioning sections is known as **Design**. The design process is illustrated with the following examples: section of the beam of desiring tenderic naturals. Incorected years, the centro region of the parametic matrix in the case of circular scention section section in the matrix means of the matrix of the matrix of the matrix

Example 10.9. Design a timber beam is to carry a load of 5 kN/m over a simply supported span of 6 m. Permissible stress in timber is 10 N/mm2 . Keep depth twice the width. Solution:

$$
w = 5 \text{ kN/m}, \quad L = 6 \text{ m}.
$$

∴ Maximum bending moment

$$
= \frac{wL^2}{8} = \frac{5 \times 6^2}{8} = 22.5 \text{ kN-m}
$$

$$
= 22.5 \times 10^6
$$
 N-mm

Let *b* be the width and *d* the depth. Then in this problem $d = 2b$

$$
Z = \frac{1}{6} bd^2 = \frac{1}{6} \times b (2b)^2
$$

$$
= \frac{2}{3} b^3
$$

 $f = 10$ N/mm² (given)

Hence design requirement is

i.e.,
\n
$$
fZ = M
$$
\n*i.e.*,
\n
$$
10 \times \frac{2}{3} b^{3} = 22.5 \times 10^{6}
$$
\n
$$
b = 150 \text{ mm.}
$$
\n
$$
d = 2b = 300 \text{ mm.}
$$

Use 150 mm \times 300 mm section.

Example 10.10. A cantilever of 3 m span, carrying uniformly distributed load of 3 kN/m is to be designed using cast iron rectangular section. Permissible stresses in cast iron are f = 30 N/mm² in tension and $f_c = 90$ N/mm² in compression. Proportion the section suitably.

Solution:

Span of cantilever = 3 m *w* = 3 kN/m ∴ Maximum moment $=\frac{wL^2}{2} = \frac{3 \times 3^2}{2}$ 2 $=\frac{3\times3^2}{2}$ = 13.5 kN-m $= 13.5 \times 10^6$ N-mm

Let *b* be the width and *d* the depth.

∴ $Z = \frac{1}{6} bd^2$

Since it is rectangular section, *N-A* lies at mid-depth, and stresses at top and bottom are same. Hence, permissible tensile stress value is reached earlier and it governs the design.

$$
f_{\text{per}} = 30 \text{ N/mm}^2
$$

∴ Design condition is

$$
fZ = M
$$

$$
30 \times \frac{1}{6} bd^2 = 13.5 \times 10^6
$$

$$
bd^2 = 2700000
$$

Using $b = 100$ mm, we get

$$
d^2 = \frac{2700000}{100}
$$

∴ $d = 164.3$ mm

Use 100 mm wide and 165 mm deep section.

Example 10.11. A circular bar of simply supported span 1 m has to carry a central concentrated load of 800 N. Find the diameter of the bar required, if permissible stress is 150 N/mm² . *Solution:* Let the diameter of the bar be '*d*'. Now, $W = 800 \text{ N}$ $L = 1 \text{ m} = 1000 \text{ mm}$.

Since it is rectangular section, *N*-*A* lies at mid-depth, and stresses at top and bottom
\nHence, permissible tensile stress value is reached earlier and it governs the design.
\n
$$
f_{\text{per}} = 30 \text{ N/mm}^2
$$

\n \therefore Design condition is
\n $fZ = M$
\n $30 \times \frac{1}{6} bd^2 = 13.5 \times 10^6$
\n $bd^2 = 2700000$
\nUsing $b = 100 \text{ mm}$, we get
\n $d^2 = \frac{2700000}{100}$
\n \therefore
\n $d = 164.3 \text{ mm}$
\nUse 100 mm wide and 165 mm deep section.
\nExample 10.11. A circular bar of simply supported span 1 m has to carry a central con
\nload of 800 N. Find the diameter of the bar required, if permissible stress is 150 N/mm
\nSolution: Let the diameter of the bar be 'd'. Now, $W = 800 \text{ N } L = 1 \text{ m} = 1000 \text{ mm}$
\n \therefore Maximum moment $= \frac{WL}{4}$
\n $= \frac{800 \times 1000}{4} = 200000 \text{ N-mm}$
\n $f = 150 \text{ N/mm}^2$
\n $Z = \frac{\pi}{32} d^3$
\n \therefore Design condition is,
\n $150 \times \frac{\pi}{32} d^3 = 200000$
\n \therefore
\n $d = 23.8 \text{ mm}$
\nSelect 25 mm bar (which is available in market)
\n10.7 SHEAR STRESS DISTRIBUTION
\nIn the 9th chapter we have seen that in a beam bending moment as well as shearing forces
\nforce gives rise to shearing stresses in the beam. In this article expression for shearing
\nderived and its variation across the section is discussed. A designer has to see that the be
\nnot only in bending but in shear also.

∴ Design condition is,

$$
150 \times \frac{\pi}{32} d^3 = 200000
$$

∴ *d* = 23.8 mm

Select 25 mm bar (which is available in market)

10.7 SHEAR STRESS DISTRIBUTION

In the $9th$ chapter we have seen that in a beam bending moment as well as shearing forces act. Shear force gives rise to shearing stresses in the beam. In this article expression for shearing stress is derived and its variation across the section is discussed. A designer has to see that the beam is safe not only in bending but in shear also.

10.7.1 Expression for Shear Stress

Consider an elemental length 'δ*x*' of beam shown in Fig. 10.17 (*a*). Let bending moment at section *A-A* be *M* and that at section *B-B* be *M* + δ*M*. Let *CD* be an elemental fibre at distance *y* from neutral axis and its thickness be δ*y*. Then,

Bending stress on left side of elemental fibre

$$
= \frac{M}{I} y b \delta y
$$

Similarly, force on right side on elemental fibre

$$
= \frac{M + \delta M}{I} \, y \, b \, dy
$$

∴ Unbalanced horizontal force on right side of elemental fibre

$$
= \frac{M + \delta M}{I} \, y \, b\delta y - \frac{M}{I} \, y \, b\delta y
$$
\n
$$
= \frac{\delta M}{I} \, yb \, \delta y
$$

There are a number of such elemental fibres above *CD*. Hence unbalanced horizontal force on section *CD*

$$
= \int_{y}^{y_{i}} \frac{dM}{I} y b \, \delta y
$$

$$
= \int_{y}^{y_{i}} \frac{dM}{I} y b \, dy = \frac{\delta M}{I} \int_{y}^{y_{i}} y b \, dy
$$

Let intensity of shearing stress on element *CD* be *q*. [Refer Fig. 10.17 (*c*)]. Then equating resisting shearing force to unbalanced horizontal force, we get

$$
q b \, \delta x = \frac{\delta M}{I} \int_{y}^{y_{i}} yb \, dy
$$

$$
\therefore \qquad q = \frac{\delta M}{\delta x} \frac{1}{bI} \int_{y}^{y_{i}} yb \, dy
$$

As $\delta x \to 0$,
$$
q = \frac{dM}{dx} \frac{1}{bI} (a\overline{y})
$$

where $\overline{a} \overline{y}$ = Moment of area above the section under consideration about neutral axis.

But we know

∴ $q = \frac{F}{bI} (a\bar{y})$...(10.11)

The above expression gives shear stress at any fibre *y* distance above neutral axis.

10.7.2 Variation of Shear Stresses Across Standard Sections

 $\frac{dM}{dx} = F$

Variation of shear stresses across the following three cases are discussed below:

- (*i*) Rectangular
- (*ii*) Circular and
- (*iii*) Isosceles triangle.
- (*i*) **Rectangular section.** Consider the rectangular section of width '*b*' and depth shown in Fig. 10.18(*a*). Let *A-A* be the fibre at a distance *y* from neutral axis. Let the shear force on the section be *F*.

Fig. 10.18

From equation (10.11), shear stress at this section is

$$
q = \frac{F}{bI} \ (a\overline{y})
$$

where $(a\bar{y})$ is the moment of area above the section about the neutral axis. Now,

where
$$
(a\overline{y})
$$
 is the moment of area above the section about the neutral axis. Now,
\n $a = b(d/2 - y)$
\n $\overline{y} = y + \frac{1}{2} (d/2 - y) = \frac{1}{2} (d/2 + y)$
\n $\therefore a\overline{y} = \frac{b}{2} (d/2 - y) \times \frac{1}{2} (d/2 + y)$
\n $= \frac{b}{2} (d^2/4 - y^2)$
\n $I = \frac{1}{12} bd^3$
\n $\therefore a = \frac{F}{b\frac{1}{12}bd^3} \frac{b}{2} (d^2/4 - y^2)$
\n $= \frac{6F}{b\frac{3}{12}} (d^2/4 - y^2)$
\nThis shows shear stress varies parabolically.
\nWhen $y = \pm d/2$, $q = 0$
\nAt $y = 0$, $q_{max} = \frac{6F}{bd^3} \frac{d^2}{4} = 1.5 \frac{F}{bd}$
\n $= 1.5 q_{av}$
\nwhere $q_{av} = \frac{F}{bd}$ is average shear stress.
\nThus in rectangular section maximum shear stress is at neutral axis and it is 1.5 time
\nshear stress. It varies parabolically from zero at extreme fibres to 1.5 q_{av} at mid depth as
\nFig. 10.18(b).
\n(ii) Circular section. Consider a circular section of diameter 'd' as shown in Fig. 10
\nwhich a shear forces *F* is to the plot. To find moment of area of the portion above
\nneutral axis, let us consider an element at distance 'z' from neutral axis. Let its the
\ndz. Let it be at an angular distance ϕ and $A \rightarrow A$ be at angular distance θ as shown in

$$
\mathcal{L}_{\mathcal{I}}
$$

$$
= \frac{6F}{bd^3} (d^2/4 - y^2)
$$

This shows shear stress varies parabolically. When $y = \pm \frac{d}{2}$, $q = 0$

At
$$
y = 0
$$
, $q_{\text{max}} = \frac{6F}{bd^3} \frac{d^2}{4} = 1.5 \frac{F}{bd}$
= 1.5 q_{av}

where $q_{\text{av}} = \frac{F}{bd}$ is average shear stress.

Thus in rectangular section maximum shear stress is at neutral axis and it is 1.5 times average shear stress. It varies parabolically from zero at extreme fibres to 1.5 q_{av} at mid depth as shown in Fig. 10.18(*b*).

(*ii*) **Circular section.** Consider a circular section of diameter '*d*' as shown in Fig. 10.19 (*a*) on which a shear force F is acting. Let $A-A$ be the section at distance 'y' from neutral axis at which shear stress is to be found. To find moment of area of the portion above *A-A* about neutral axis, let us consider an element at distance '*z*' from neutral axis. Let its thickness be *dz*. Let it be at an angular distance φ and *A-A* be at angular distance θ as shown in figure.

[Since if cos $\phi = t$, $dt = -\sin \phi \, d\phi$ and $- t^3/3$ is integration]

$$
(a\overline{y}) = \frac{d^3}{4 \times 3} \left[-\cos^2 \frac{\pi}{2} + \cos^3 \theta \right]
$$

\n
$$
= \frac{d}{12} \cos^3 \theta
$$

\nNow
\n
$$
I = \frac{\pi d^4}{64}
$$

\n
$$
\therefore \qquad q = \frac{F}{bI} (a\overline{y})
$$

\n
$$
= \frac{F}{12 \pi d^2} \cos^3 \theta
$$

\n
$$
= \frac{64}{12 \pi d^2} \cos^2 \theta
$$

\n
$$
= \frac{16}{12 \pi d^2} \left[1 - \sin^2 \theta \right]
$$

\n
$$
= \frac{16}{3} \frac{F}{\pi d^2} \left[1 - \left(\frac{y}{d/2} \right)^2 \right]
$$

\n
$$
= \frac{16}{3} \frac{F}{\pi d^2} \left[1 - \left(\frac{y}{d/2} \right)^2 \right]
$$

\nHence shear stress varies parabolically.
\nAt $y = \pm d/2$, $q = q$
\n
$$
y = 0, \qquad q = q_{\text{max}} = \frac{16}{3} \frac{F}{\pi d^2}
$$

\n
$$
= \frac{4}{3} q_{\text{av}}
$$

\nwhere q_{av} = average shear stress.
\nThus in circular sections also shear stress varies parabolically from zero at extreme edges to the maximum value of $\frac{4}{3} q_{\text{av}}$ at mid depth as shown in Fig. 10.19(b).

where q_{av} = average shear stress.

Thus in circular sections also shear stress varies parabolically from zero at extreme edges to the maximum value of $\frac{4}{3}$ q_{av} at mid depth as shown in Fig. 10.19(*b*).

(*iii*) **Isosceles triangular section.** Consider the isosceles triangular section of width '*b*' and height '*h*' as shown in Fig. 10.20(*a*). Its centroid and hence neutral axis is at $\frac{2}{3}$ 3 $\frac{h}{\sqrt{2}}$ from top fibre. Now shear stress is to be found at section *A-A* which is at a depth '*y*' from top fibre.

$$
a\overline{y} = \frac{1}{2} \frac{b}{h} y^2 \left(\frac{2h}{3} - \frac{2y}{3} \right)
$$

$$
= \frac{1}{3} \frac{b}{h} y^2 (h - y)
$$

Moment of inertia of the section

$$
I=\frac{bh^3}{36}.
$$

∴ Shear stress at *A-A* $q = \frac{F}{bl} a \overline{y}$ $=\frac{F}{\sqrt{F}}$ *y* $\frac{y}{h}b \times \frac{bh}{36}$ *b* $\times \frac{bh^3}{26}$ 3 h $\frac{1}{3}$ \times 36 $\frac{1}{3} \frac{b}{h} y^2 (h - y)$ $=\frac{12 F}{h h^3}$ $\frac{12 F}{bh^3}$ *y*(*h* – *y*) Hence at $y = 0$, $q = 0$ At $y = h$, $q = 0$ At centroid, $y = \frac{2}{3}$ 3 *h* $q = \frac{12 F}{\sqrt{3}} \frac{2}{7}$ $3\overline{3}$ *F bh* $\frac{h}{h}$ (*h* – 2*h*/3) $=\frac{8}{3}$ $\frac{F}{bh} = \frac{4}{3}$ *F* 1 2/ *bh* $=\frac{4}{3} q_{\text{av}}$ where q_{av} is average shear stress. For q_{max} , $\frac{dq}{dy} = 0$ *i.e.*, $\frac{12 \, I}{h b^3}$ $\frac{12 F}{bh^3}$ (*h* – 2*y*) = 0 *i.e*., at *y* = *h*/2 and hence $q_{\text{max}} = \frac{12 F}{bh^3} \cdot \frac{h}{2}$ *F bh* $\frac{h}{2}$ (*h* – *h*/2) $=\frac{12}{1}$ 4 *F* 3 *bh* $=\frac{3F}{bh}$ $=\frac{1.5}{1/2}$. *F* / *bh* $= 1.5 q_{\text{av}}$. $\frac{12F}{h}b \times \frac{bh}{3h} \times \frac{1}{3} \frac{b}{h} y^2 (h - y)$

Hence at $y = 0$, $q = 0$

At ventroid, $y = \frac{2h}{3}$
 $q = \frac{12F}{bh^3} \frac{2h}{3} (h - 2h/3)$
 $= \frac{8F}{3bh} - \frac{4}{3} \frac{4F}{32bh}$
 $= \frac{4}{3}\frac{F}{bh} - \frac{4}{3} \frac{4F}{32bh}$

where q_{xx} is ave

Thus in isosceles triangular section shear stress is zero at extreme fibres, it is maximum of 1.5 q_{av} at mid depth and has a value $\frac{4}{3}$ q_{av} at neutral axis. The variation of shear stress is as shown in Fig. 10.20(*b*).

10.8 SHEAR STRESSES IN BUILT-UP SECTIONS

In sections like *I*, *T* and channel, shear stresses at various salient points are calculated and the shear stress variation diagram across depth is plotted. It may be noted that at extreme fibres shear stress is zero since $(a\bar{y})$ term works out to be zero. However it may be noted that the procedure explained below is for built up section with at least one symmetric axis. If there is no symmetric axis along the depth analysis for shear stress is complex, and that is treated beyond the scope to this book.

Example 10.12. Draw the shear stress variation diagram for the I-section shown in Fig. 10.21(a) if it is subjected to a shear force of 100 kN.

Fig. 10.21

Solution: Due to symmetry neutral axis is at mid depth.

$$
I = \frac{1}{12} \times 180 \times 10^3 + 180 \times 10 \times (200 - 5)^2
$$

+
$$
\frac{1}{12} \times 10 \times 380^2 + 10 \times 380 \times (200 - 200)^2
$$

+
$$
\frac{1}{12} \times 180 \times 10^3 + 180 \times 10 \times (200 - 5)^2
$$

 $= 182.646666 \times 10^6$ mm⁴

Shear stress at $y = 200$ mm is zero since $a\overline{y} = 0$.

Shear stress at bottom of top flange

 $=\frac{F}{bl}$ ($a\overline{y}$)

$$
= \frac{100 \times 1000}{180 \times 182.646666 \times 10^6} \times (180 \times 10 \times 195)
$$

 $= 1.068$ N/mm²

Shear stress in the web at the junction with flange

$$
= \frac{100 \times 1000}{10 \times 182.646666 \times 10^6} (180 \times 10 \times 195)
$$

$$
= 19.217 \text{ N/mm}^2
$$

Shear stress at *N-A*

$$
= \frac{100 \times 1000}{10 \times 182.646666} \times \left[180 \times 10 \times 195 + 10 \times (200 - 10) \times \frac{190}{2} \right]
$$

= 29.10 N/mm².

Symmetric values will be there on lower side. Hence shear stress variation is as shown in Fig. 10.21(*b*).

Example 10.13. A beam has cross-section as shown in Fig. 10.22(a). If the shear force acting on this is 25 kN, draw the shear stress distribution diagram across the depth.

Solution: Let \bar{y} be the distance of centroid of the section from its top fibre. Then

$$
y' = \frac{\text{Moment of area about top fibre}}{\text{Total area}}
$$

=
$$
\frac{120 \times 12 \times 6 + (120 - 12) \times 12 \times \left(12 + \frac{120 - 12}{2}\right)}{120 \times 12 + (120 - 12) \times 12}
$$

= 34.42 mm

 $\frac{1}{2}$

2

∴ Moment of inertia about centroid

$$
I = \frac{1}{12} \times 120 \times 12^{3} + 120 \times 12 (34.42 - 6)^{2}
$$

$$
+ \frac{1}{12} \times 12 \times 108^{3} + 12 \times 108 \left(34.42 - \frac{108}{2}\right)^{2}
$$

 $= 2936930$ mm⁴

Shear stresses are zero at extreme fibres.

Shear stress at bottom of flange:

Area above this level, $a = 120 \times 12 = 1440$ mm²

Centroid of this area above *N-A*

$$
\bar{y}
$$
 = 34.42 - 6 = 28.42 mm

Width at this level $b = 120$ mm.

$$
\therefore \t q_{\text{bottom of flange}} = \frac{25 \times 1000}{120 \times 2936930} \times 1440 \times 28.42
$$

$$
= 2.90 \text{ N/mm}^2
$$

Shear stress at the same level but in web, where width $b = 12$ mm

$$
= \frac{25 \times 1000}{12 \times 2936930} \times 1440 \times 28.42
$$

$$
= 29.0 \text{ N/mm}^2
$$

Shear stress at neutral axis:

For this we can consider $a\overline{y}$ term above this section or below this section. It is convenient to consider the term below this level. $+\frac{1}{12} \times 12 \times 108^3 + 12 \times 108 \left(34.42 - \frac{108}{2}\right)^9$
 $= 2936930$ mm⁴

Shear stresses are zero at extreme fibres.

Shear stresses at bottom of flange:

Area above bis level, $a = 120 \times 12 = 1440$ mm²

Centroid of this

 $a = 12 \times (120 - 34.42) = 1026.96$ mm²

The distance of its centroid from *N-A*

$$
= \frac{120 - 34.42}{2} = 42.79
$$
 mm.

Width at this section $b = 12$ mm.

$$
q = \frac{25 \times 1000}{12 \times 2936930} \times 1026.96 \times 42.79
$$

$$
= 31.17 \text{ N/mm}^2
$$

Hence variation of shear stress across the depth is as shown in Fig. 10.22(*b*).

Example 10.14. The unsymmetric I-section shown in Fig. 10.23(a) is the cross-section of a beam, which is subjected to a shear force of 60 kN. Draw the shear stress variation diagram across the depth.

Fig. 10.23

Solution: Distance of neutral axis (centroid) of the section from top fibre be y_t . Then

$$
100 \times 20 \times 10 + (200 - 20 - 20) \times 20 \times \left(20 + \frac{160}{2}\right)
$$

\n
$$
y_t = \frac{100 \times 20 + 160 \times 20 + 150 \times 20}{100 \times 20 + 160 \times 20 + 150 \times 20}
$$

\n
$$
I = \frac{1}{12} \times 100 \times 20^3 + 100 \times 20 (111 - 10)^2
$$

\n
$$
+ \frac{1}{12} \times 20 \times 160^3 + 160 \times 20 (111 - 100)^2
$$

\n
$$
+ \frac{1}{12} \times 150 \times 20^3 + 150 \times 20 (111 - 190)^2
$$

\n= 46505533 mm⁴

Shear stress at bottom of top flange

∴ Shear stress

$$
= \frac{F}{bI} a\overline{y}
$$

= $\frac{60 \times 1000}{100 \times 46505533} \times 100 \times 20 \times (111 - 10)$
= 2.61 N/mm²
at the same level, but in web
= $\frac{60 \times 1000}{20 \times 46505533} \times 100 \times 20 (111 - 10)$
= 13.03 N/mm²

 $\frac{1}{2}$

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Shear stress at neutral axis:

$$
a\overline{y} = a\overline{y}
$$
 of top flange + $a\overline{y}$ of web above *N-A*

$$
= 100 \times 20 \times (111 - 10) + 20 \times (111 - 20) \times \frac{111 - 20}{2}
$$

 $= 284810$ mm³.

∴ Shear stress at neutral axis

$$
= \frac{F}{bl} (a\overline{y})
$$

= $\frac{60 \times 1000}{20 \times 46505533} \times 284810$
= 18.37 N/mm².

Shear stress at junction of web and lower flange:

Considering the lower side of the section for finding $a\overline{y}$, we get

$$
a\overline{y}
$$
 = 150 × 20 × (190 – 111) = 237000 mm³

$$
q = \frac{60 \times 1000}{20 \times 46505533} \times 237000
$$

$$
= 15.28 \text{ N/mm}^2
$$

At the above level but in web, shear stress

$$
= \frac{60 \times 1000}{150 \times 46505533} \times 237000
$$

$$
= 2.04 \text{ N/mm}^2
$$

At extreme fibres shear stress is zero. Hence variation of shear across the depth of the section is as shown in Fig. 10.23.

IMPORTANT FORMULAE

- **1.** Bending equation: $\frac{M}{I}$ *I f y* $=\frac{f}{y}=\frac{E}{R}.$
- **2.** Modulus of section $Z = \frac{I}{I}$ $\frac{1}{y_{\text{max}}}$.

3. Moment carrying capacity of section $=f_{\text{per}} Z$.

4. Section modulus of various sections:

$$
= 100 \times 20 \times (111 - 10) + 20 \times (111 - 20) \times 2
$$

\n
$$
= 284810 \text{ mm}^3.
$$

\n∴ Shear stress at neutral axis
\n
$$
= \frac{F}{bl} (a\bar{y})
$$

\n
$$
= 18.37 \text{ Nmm}^2.
$$

\nShear stress at junction of web and lower fange:
\nConsidering the lower side of the section for finding a\bar{y}, we get
\n
$$
a\bar{y} = 150 \times 20 \times (190 - 111) = 237000 \text{ mm}^3
$$

\n∴
$$
q = \frac{60 \times 1000}{20 \times 46505533} \times 237000
$$

\n
$$
= 15.28 \text{ N/mm}^2
$$

\nAt the above level but in web, shear stress
\n
$$
= \frac{60 \times 1000}{150 \times 4650533} \times 237000
$$

\nAt extreme fibres shear stress is zero. Hence variation of shear across the depth of the
\nis as shown in Fig. 10.23.
\n
\n1. Bending equation: $\frac{M}{I} = \frac{f}{\bar{y}} = \frac{E}{R}.$
\n2. Modulus of section $Z = \frac{I}{y_{max}}$.
\n3. Moment carrying capacity of section $= f_{\text{per}} Z$.
\n4. Section modulus of various sections:
\n(i) Rectangular: $\frac{1}{6} bd^2$ (ii) Hollow rectangular: $\frac{1}{6} \frac{BD^3-bd^3}{D}$
\n(ii) Solid circular section: $\frac{\pi}{32} d^3$ (iv) Hollow circular section: $\frac{\pi}{32} \frac{D^4-d^4}{D}$
\n(v) Solid triangular section: $\frac{bh^2}{24}$

5. Shear stress in a beam $q = \frac{F}{bI}(a\bar{y})$

6. In rectangular sections,

$$
q_{\text{max}} = 1.5 \, q_{\text{av}}, \text{ at } y = d/2
$$

In circular sections $q_{\text{max}} = \frac{4}{3} q_{\text{av}}$, at centre

In triangular section, $q_{\text{max}} = 1.5 q_{\text{av}}$, at $y = \frac{h}{2}$.

THEORY QUESTIONS

1. State the assumptions made in deriving bending equation.

2. Derive the bending equation
$$
\frac{M}{I} = \frac{F}{Y} = \frac{E}{R}
$$
 with usual meaning.

- **3.** Derive the expression for finding shear stress in a beam in the form $q = \frac{F}{bI}(a\bar{y})$ with usual notations.
- **4.** Show that maximum shear stress in a beam of rectangular section is 1.5 q_{average} .
- **5.** Show that in a beam of solid circular section maximum shear stress is $\frac{4}{3}$ $\frac{1}{3}$ × *q* average.
- **6.** Draw the variation of shear stresses across the depth of a beam of triangular section, finding the values at critical depths.

PROBLEMS FOR EXERCISE

1. A *I*-section has flanges of size 200×12 mm and its overall depth is 360 mm. Thickness of web is also 12 mm. It is used as a simply supported beam over a span of 4 m to carry a load of 60 kN/m over its entire span. Draw the variation of bending and shearing stresses across the depth.

Fig. 10.24

2. Unsymmetric *I*-section shown in Fig. 10.25 is used as a cantilever of span 2 m to carry uniformly distributed load of 6 kN/m over entire span. Draw the variation of bending stress across the depth marking the values at salient point.

3. Calculate the variation of shear stresses at various salient level near fixed support in the above problem and give the sketch. [**Ans.** See Fig. 10.26]

4. A cantilever beam of 1.2 m span is having cross-section as shown in Fig. 10.27. The permissible stresses in tension and compressions are 20 N/mm2 and 80 N/mm2 respectively. Determine the maximum concentrated load *W* it can carry at the free end. [**Ans.** *W* = 44.576 kN] [**Note:** In cantilever tension is at top.]

5. A timber beam is to be designed to carry a load of 6 kN/m over a simply supported span of 5 m. Permissible stress is 10 N/mm2. Keeping the depth twice the width, design the beam. If the permissible stress in shear is 1 N/mm2, check for shear.

[Ans. Required $b = 141.18$, $d = 282.31$. Hence select 150×300 mm section]

11 Principal Stresses and Strains

A structural member need not be always under simple (only one type) of stress. It may be subjected to direct stresses in different directions and may be subjected to shear stresses also. A beam is usually subjected to axial stresses due to bending and also for shear stresses. A shaft is subjected to shear stresses due to torsion and axial stresses due to bending/direct thrust. The stresses may vary from point to point in the member.

In a three-dimensional system, stresses acting at a point may be represented as shown in Fig. 11.1.

In many engineering problems, two dimensional idealization is made as shown in Fig. 11.2. In this chapter discussion is limited to two dimensional stress system. First general expression for stresses on a plane inclined at a selected axis is discussed. Then the terms principal stresses and planes are explained and the expressions to get them are presented. A number of problems are solved to make the concept clear. At the end analysis is given for principal strains also.

11.1 STRESSES ON INCLINED PLANES

Consider the element in a structural member which is under a general two dimensional state of stress as shown in Fig. 11.3(*a*). Note the sign convention:

+ve face +ve direction is +ve stress for p_x and p_y . In other words tensile stresses are +ve. Shearing stress as shown in the figure is +ve stress.

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Now we are interested in finding state of stress on plane *DE* which makes anticlockwise angle θ with the plane of stress p_x , in other words with *y*-axis.

For simplicity let us consider thickness of the element as unity. We are interested in finding normal and tangential stresses acting on the plane *DE*. Let normal stress be p_n and tangential stress p_t as shown in Fig. 11.3(*b*). Since the system is in equilibrium,

$$
\Sigma
$$
 Forces normal to $DE = 0$ gives

$$
P_x
$$

\nNow we are interested in finding state of stress on plane *DE* which makes anticlockwise angle
\nθ with the plane of stress *p_x*, in other words with y-axis.
\nFor simplicity let us consider thickness of the element as unity. We are interested in finding
\nnormal and tangential stresses acting on the plane *DE*. Let normal stress be *p_n* and tangential stress
\n*p_i* as shown in Fig. 11.3(b). Since the system is in equilibrium,
\n
$$
\sum \text{ Forces normal to } DE = 0 \text{ gives } CD + q \times CD \times 1 \times \sin \theta + q \times CE \times 1 \times \cos \theta
$$
\n
$$
P_n \times DE \times 1 = p_x \times CD \times 1 \times \cos \theta + q \times CD \times 1 \times \sin \theta + q \times CE \times 1 \times \cos \theta
$$
\n
$$
P_n = p_x \frac{CD}{DE} \cos \theta + p_y \frac{CE}{DE} \sin \theta + q \frac{CD}{DE} \sin \theta + q \frac{CE}{DE} \cos \theta
$$
\nSince\n
$$
\frac{CD}{DE} = \cos \theta \text{ and } \frac{CE}{DE} = \sin \theta, \text{ we get}
$$
\n
$$
p_n = p_x \cos^2 \theta + p_y \sin^2 \theta + q \cos \theta \sin \theta + q \sin \theta \cos \theta
$$
\n
$$
= p_x \left(\frac{1 + \cos 2\theta}{2}\right) + p_y \left(\frac{1 - \cos 2\theta}{2}\right) + 2q \sin \theta \cos \theta.
$$
\nsince\n
$$
\cos^2 \theta = \frac{1 + \cos 2\theta}{2} \text{ and } \sin^2 \theta = \frac{1 - \cos 2\theta}{2}
$$
\nThus\n
$$
p_x \times DE \times 1 = p_x \times CD \times 1 \times \sin \theta - q \times CD \times 1 \times \cos \theta
$$
\n
$$
= p_x \frac{CD}{DE} \cos \theta - q \frac{CD}{DE} \cos \theta + q \times CE \times 1 \times \sin \theta
$$
\n
$$
p_x \times DE \times 1 = p_x \times CD \times 1 \times \sin \theta - q \times CD \times 1 \times \cos \theta
$$
\n...(11.1)
\nSimilarly from equilibrium condition of forces tangential to plane *DE*, we get\n
$$
p_x \times DE \times 1 =
$$

Thus $p_n = \frac{p_x + p_y}{2} + \frac{p_x - p_y}{2} \cos 2\theta + q$ 2 2 $\cos 2\theta + q \sin 2\theta$...(11.1)

Similarly from equilibrium condition of forces tangential to plane *DE*, we get

 $p_t \times DE \times 1 = p_x \times CD \times 1 \times \sin \theta - q \times CD \times 1 \times \cos \theta$

 $-p_y \times CE \times 1 \times \cos \theta + q \times CE \times 1 \times \sin \theta$

$$
p_t = p_x \frac{CD}{DE} \sin \theta - p_y \frac{CE}{DE} \cos \theta - q \frac{CD}{DE} \cos \theta + q \frac{CE}{DE} \sin \theta
$$

But
$$
\frac{CD}{DE} = \cos \theta \quad \text{and} \quad \frac{CE}{DE} = \sin \theta.
$$

$$
\therefore \quad \begin{array}{c} DE & DE \\ p_t = p_x \sin \theta \cos \theta - p_y \sin \theta \cos \theta - q \cos^2 \theta + q \sin^2 \theta \end{array}
$$

$$
= (p_x - p_y) \sin \theta \cos \theta - q(\cos^2 \theta - \sin^2 \theta)
$$

$$
= \frac{p_x - p_y}{2} \sin 2\theta - q \cos 2\theta,
$$

since $2 \sin \theta \cos \theta = \sin 2\theta$ and $\cos^2 \theta - \sin^2 \theta = \cos 2\theta$

Thus

$$
p_t = \frac{p_x - p_y}{2} \sin 2\theta - q \cos 2\theta \qquad \qquad \dots (11.2)
$$

Thus in case of a general two dimension stress system the normal and tangential stresses acting on a plane making anticlockwise angle θ with the plane of $p_r(y)$ direction) are as given by equations (11.1) and (11.2).

Angle of Obliquity of Resultant Stress

The resultant of the stresses on the above inclined plane can be found as

$$
R = \sqrt{p_n^2 + p_t^2}
$$

which is inclined at 'α' to the plane (Ref. Fig. 11.3*c*).

 $\tan \alpha = \frac{p}{p}$ *n t*

In other words, the resultant is inclined at $\theta + \alpha$ to the plane of p_x .

11.2 PRINCIPAL STRESSES AND PLANES

The planes on which shearing stresses are zero are called *principal planes* and the stresses normal to principal planes are known as *principal stresses.*

Principal Plane

To locate principal plane

$$
p_t = 0
$$

From eqn. (11.2),
$$
\frac{p_x - p_y}{2} \sin 2\theta - q \cos 2\theta = 0
$$

$$
\therefore \qquad \tan 2\theta = \frac{q}{(p_x - p_y)/2} = \frac{2q}{p_x - p_y} \qquad \qquad \dots (11.3)
$$

There are two values for 2θ which differ by 180° for which eqn. (11.3) can be satisfied. Let $2\theta_1$ and $2\theta_2$ be the solution. Referring to Fig. 11.4, we find

Similarly,

$$
\sin 2\theta_2 = \frac{-2q}{\sqrt{(p_x - p_y)^2 + 4q^2}}
$$
...(11.5*a*)
and
$$
\cos 2\theta_2 = \frac{-(p_x - p_y)}{\sqrt{(p_x - p_y)^2 + 4q^2}}
$$
...(11.5*b*)

 $2θ_1$ and $2θ_2$ differ by 180°. Hence we can say $θ_1$ and $θ_2$ differ by 90°. Thus direction of principal planes to the plane of p_n are given by eqn. (11.4). Another principal plane is at right angles to it.

Principal Stresses

Principal stresses are the normal stresses on principal planes. Hence the values of principal stresses may be obtained by substitutes θ_1 and θ_2 values for θ in the expression for p_n . Denoting the values as p_1 and p_2 , we get

eqn. (11.3) can be satisfied. Let 2
$$
\theta_1
$$
 and 2 θ_2 be the solution.
\nReferring to Fig. 11.4, we find
\n
$$
\sin 2\theta_1 = \frac{2q}{\sqrt{(p_x - p_y)^2 + 4q^2}} \qquad ...(11.4a)
$$
\n
$$
\cos 2\theta_1 = \frac{p_x - p_y}{\sqrt{(p_x - p_y)^2 + 4q^2}} \qquad ...(11.4b)
$$
\nSimilarly,
\n
$$
\sin 2\theta_2 = \frac{-2q}{\sqrt{(p_x - p_y)^2 + 4q^2}}
$$
\nand
\n
$$
\cos 2\theta_2 = \frac{-(p_x - p_y)}{\sqrt{(p_x - p_y)^2 + 4q^2}}
$$
\nand
\n
$$
\cos 2\theta_2 = \frac{-(p_x - p_y)}{\sqrt{(p_x - p_y)^2 + 4q^2}}
$$
\nand
\n
$$
\theta_1
$$
 and 2 θ_2 differ by 180°. Hence we can say θ_1 and θ_2 differ by 90°. Thus direction on
\nplanes to the plane of p_n are given by eqn. (11.4). Another principal plane is at right ar
\nPrincipal stresses
\nPrincipal stresses
\nPrincipal stresses
\nPrincipal systems
\n
$$
p_1 = p_n
$$
 at $\theta = \theta_1$
\n
$$
= \frac{p_x + p_y}{2} + \frac{p_x - p_y}{2} \qquad \frac{p_x - p_y}{\sqrt{(p_x - p_y)^2 + 4q^2}} + \frac{2q}{\sqrt{(p_x - p_y)^2 + 4q^2}}
$$
\n
$$
= \frac{p_x + p_y}{2} + \frac{p_x - p_y}{2} \qquad \frac{p_x - p_y}{\sqrt{(p_x - p_y)^2 + 4q^2}} + \frac{2q}{\sqrt{(p_x - p_y)^2 + 4q^2}}
$$
\n
$$
= \frac{p_x + p_y}{2} + \frac{1}{2} \sqrt{(p_x - p_y)^2 + 4q^2}
$$
\n
$$
= \frac{p_x + p_y}{2} + \frac{1}{2} \sqrt{(p_x - p_y)^2 + 4q^2}
$$
\n
$$
= \frac{p_x + p_y}{2} + \frac{p_x - p_y}{2} \qquad \cos 2\theta_
$$

 $p_x - p_y^2 + 4q$

 $-p_v)^2 +$

2 $\sqrt{(p_x - p_y)^2 + 4q^2}$ $\sqrt{(p_x - p_y)^2 + 4q^2}$

 $\chi = \mu_y$ *y* $\pm 4q$ $\sqrt{v_x - \mu_y}$

 $(p_x - p_y)^2 + 4q^2 \sqrt{(p_x - p_y)^2 + 4q^2}$

 $p_x - p_y^2 + 4q$

 $(p_x - p_y)$

 $2+40$
$$
= \frac{p_x + p_y}{2} - \frac{1}{2} \frac{(p_x - p_y)^2 + 4q^2}{\sqrt{(p_x - p_y)^2 + 4q^2}}
$$

$$
= \frac{p_x + p_y}{2} - \frac{1}{2} \sqrt{(p_x - p_y)^2 + 4q^2}
$$

$$
= \frac{p_x + p_y}{2} - \sqrt{\left(\frac{p_x - p_y}{2}\right)^2 + q^2}
$$

Thus the principal stresses are given by

$$
= \frac{p_x + p_y}{2} - \frac{1}{2}\sqrt{(p_x - p_y)^2 + 4q^2}
$$

\n
$$
= \frac{p_x + p_y}{2} - \sqrt{\left(\frac{p_x - p_y}{2}\right)^2 + q^2}
$$

\nThus the principal stresses are given by
\n
$$
p_1 = \frac{p_x + p_y}{2} + \sqrt{\left(\frac{p_x - p_y}{2}\right)^2 + q^2}
$$
...(11.6*a*)
\nand
\n
$$
p_2 = \frac{p_x + p_y}{2} - \sqrt{\left(\frac{p_x - p_y}{2}\right)^2 + q^2}
$$
...(11.6*b*)
\nIt can be proved that principal stresses are maximum and minimum stresses also. To find
\nextreme value of normal stress p_n ,
\n
$$
\frac{dp_x}{d\theta} = 0
$$

\ni.e.,
\n
$$
\frac{p_x - p_y}{2} = 2(\text{– sin } 2\theta) + q \cdot 2 \text{ cos } 2\theta = 0
$$

\ni.e.,
\n
$$
p_i = 0
$$

\nThus the principal planes are the planes of maximum/minimum normal stresses also. Plane
\nconresponding to θ_1 gives minimum
\nnormal stress.
\nFor maximum shear
\nFrom all stress.
\n
$$
\text{Plane of Maximum State S} = 0
$$

\ni.e.,
\n
$$
\frac{p_x - p_y}{2} \cdot 2 \text{ cos } 2\theta' - q \cdot 2 (\text{– sin } 2\theta') = 0
$$

\n
$$
\text{ln } 2\theta = \frac{1 - (p_x - p_y)}{2q}
$$
...(11.7)
\nFrom eqn. (11.3) and eqn. (11.7), we get
\n
$$
\tan 2\theta' = \frac{-(p_x - p_y)}{2q}
$$
...(11.7)
\nFrom eqn. (11.3) and eqn. (11.7), we get
\n
$$
\tan 2\theta' = -1
$$

\nHence 2 θ' and 2 θ values differ by 90°. In other word planes of extreme shearing stresses are

It can be proved that principal stresses are maximum and minimum stresses also. To find extreme value of normal stress p_n ,

$$
\frac{dp_n}{d\theta} = 0
$$

i.e.,
$$
\frac{p_x - p_y}{2} \quad 2(-\sin 2\theta) + q \quad 2\cos 2\theta = 0
$$

i.e.,
$$
\frac{p_x - p_y}{2} \sin 2\theta - q \cos 2\theta = 0
$$

2

i.e., $p_t = 0$

Thus the *principal planes are the planes of maximum/minimum normal stresses also.* Plane corresponding to θ_1 gives the maximum value while plane corresponding to θ_2 gives minimum normal stress.

sin $2θ - q cos 2θ = 0$

Plane of Maximum Shear Stress

For maximum shear

$$
\left. \frac{dp_t}{d\theta} \right|_{\theta = \theta'} = 0
$$

i.e.,
$$
\frac{p_x - p_y}{2}
$$
 2 cos 2 θ' - q 2 (- sin 2 θ') = 0

or
$$
\tan 2\theta' = \frac{-(p_x - p_y)}{2q}
$$
 ...(11.7)

From eqn. (11.3) and eqn. (11.7), we get

tan $2θ \times \tan 2θ' = -1$

Hence 2θ′ and 2θ values differ by 90°. In other word *planes of extreme shearing stresses are at* $\frac{90}{2}$ = 45° to the principal planes.

Value of maximum shearing stress may be obtained by substituting θ' for θ in equation for p_t (eqn. 11.2)

Now
$$
\tan 2\theta' = \frac{-(p_x - p_y)}{2q}
$$

\nFrom Fig. 11.5, we get
\n
$$
\sin 2\theta' = \frac{(p_x - p_y)}{\sqrt{(p_x - p_y)^2 + 4q^2}}
$$
\n
$$
\cos 2\theta' = \frac{-2q}{\sqrt{(p_x - p_y)^2 + 4q^2}}
$$
\n
$$
\therefore \qquad p_{r \text{ (max)}} = \frac{\frac{p_x - p_y}{2}}{\frac{p_x - p_y}{2}} \sin 2\theta' - q \cos 2\theta'
$$
\n
$$
= \frac{\frac{p_x - p_y}{2}}{\frac{p_x - p_y}{2}} \frac{p_x - p_y}{\sqrt{(p_x - p_y)^2 + 4q^2}} \frac{q(-2q)}{\sqrt{(p_x - p_y)^2 + 4q^2}}
$$
\n
$$
= \frac{1}{2} \frac{(p_x - p_y)^2 + 4q^2}{\sqrt{(p_x - p_y)^2 + 4q^2}}
$$
\n
$$
= \frac{1}{2} \sqrt{(p_x - p_y)^2 + 4q^2}
$$
\n
$$
= \frac{1}{2} \sqrt{(p_x - p_y)^2 + 4q^2}
$$
\n
$$
= \frac{1}{2} \text{ the difference between } p_1 \text{ and } p_2
$$
\n
$$
= \frac{p_1 - p_2}{2}
$$
\n
$$
= \frac{p_1 - p_2}{2}
$$
\n
$$
\therefore (11.8a)
$$
\nas shown in Fig. 11.6(a) or 11.6(b).
\nas shown in Fig. 11.6(a) or 11.6(b).
\n
$$
p_1 = \text{plane}
$$
\n
$$
\therefore \qquad p_x = \text{plane}
$$
\n
$$
\therefore \qquad p_y = \text{plane}
$$
\n
$$
\therefore \qquad p_z = \text{plane}
$$
\n
$$
\therefore \qquad p_z = \text{plane}
$$
\n
$$
\therefore \qquad p_z = \text{plane}
$$
\n
$$
\therefore \q
$$

The direction of principal planes and the planes of maximum shearing stresses may be indicated as shown in Fig. 11.6(*a*) or 11.6(*b*).

Fig. 11.6

Example 11.1. Show that sum of normal stresses in any two mutually perpendicular directions is constant in case of a general two dimensional stress.

Solution: Let p_x and p_y as shown in Fig. 11.3(*a*) be normal stresses and *q* be the shearing stresses. Let the thickness of element be unity.

Taking element as shown in Fig. 11.3(*b*) and considering the equation of equilibrium in the direction normal to the plane *DE*, we get

 $p_n \times 1 \times DE = p_x \times 1 \times CD \cos \theta + p_y \times 1 \times CE \sin \theta$ $+ q \times 1 \times CD \sin \theta + q \times 1 \times CE \cos \theta$ $$ Noting that $\frac{CD}{DE} = \cos \theta$ and $\frac{CE}{DE} = \sin \theta$, we get Example the contents of control of equilibrity direction normal to the plane DE .

Taking element as shown in Fig. 11.3(b) and considering the equation of equilibrity
 $p_n \times 1 \times DE = p$, $\times 1 \times CD \cos \theta + p$, $\times 1 \times CE \sin \theta + \cos \theta + \cos \$

 $\left(\right)$

p_n = *p_x* cos² θ + *p_y* sin² θ + *q* cos θ sin θ + *q* sin θ cos θ

 $1 - \cos 2$ 2 $\left(\frac{1+\cos 2\theta}{2}\right) + p_y\left(\frac{1-\cos 2\theta}{2}\right) + 2q \sin \theta \cos \theta$

 $\frac{1}{2}$

 $=\frac{p_x + p_y}{2} + \frac{p_x - p_y}{2} \cos 2\theta + q \sin 2\theta$...(1)

$$
i.e.,
$$

i.e., $p_n = p_x \left(\frac{1 + \cos 2\theta}{2} \right) + p_y$

If p_n' is the stress on a plane at right angles to *CD*, $\theta' = \theta + 90^\circ$. Hence from the above general expression for stress on an inclined plane, we get

2

 $\left(\right)$

$$
p_n' = \frac{p_x + p_y}{2} + \frac{p_x - p_y}{2} \cos 2(\theta + 90) + q \sin 2(\theta + 90)
$$

=
$$
\frac{p_x + p_y}{2} - \frac{p_x - p_y}{2} \cos 2\theta - q \sin 2\theta \qquad ...(2)
$$

Adding eqns. (1) and (2), we get

$$
p_n + p_n' = p_x + p_y
$$

Thus *sum of normal stresses in any two mutually perpendicular planes is constant and is equal to* $p_x + p_y$.

Example 11.2. A material has strength in tension, compression and shear as 30N/mm², 90 N/mm² *and 25 N/mm² , respectively. If a specimen of diameter 25 mm is tested in tension and compression identity the failure surfaces and loads.*

Solution:

(*a*) In tension: Let axial direction be *x* direction.

Since it is uniaxial loading,

 $p_y = 0$, $q = 0$ and only p_x exists.

 \therefore when the material is subjected to full tensile stress, $p_x = 30 \text{ N/mm}^2$.

$$
p_1 = \frac{30 + 0}{2} + \sqrt{\left(\frac{30 - 0}{2}\right)^2 + 0^2} = 30 \text{ N/mm}^2
$$

$$
p_2 = \frac{30 + 0}{2} - \sqrt{\left(\frac{30 - 0}{2}\right)^2 + 0^2} = 0
$$

$$
q_{\text{max}} = \frac{p_x - p_y}{2} = \frac{30 - 0}{2} = 15 \text{ N/mm}^2 < 25 \text{ N/mm}^2.
$$

Hence failure criteria is normal stress $p_1 = 30$ N/mm². Corresponding load *P* is obtained by

$$
p = \frac{P}{A}
$$

30 = $\frac{P}{(\pi/4) \times 25^2}$ or **P** = **14726 N.**

Failure surface is given by plane of p_1 which is as shown in Fig. 11.7. (*b*) In case of compression test **Fig. 11.7**

Hence failure criteria is normal stress
$$
p_1 = 30 \text{ N/mm}^2 < 25 \text{ N/mm}^2
$$
.
\nHence failure criteria is normal stress $p_1 = 30 \text{ N/mm}^2$.
\nCorresponding load *P* is obtained by
\n $p = \frac{P}{A}$
\n $30 = \frac{P}{(\pi/4) \times 25^2}$ or $P = 14726 \text{ N}$.
\nFailure surface is given by plane of p_1 which is as shown in Fig. 11.7.
\n(b) In case of compression test
\n $p_x = -90 \text{ N/mm}^2$, $p_y = q = 0$
\n \therefore At failure when $p_x = -90 \text{ N/mm}^2$, load is
\n $-90 = \frac{P}{(\pi/4) \times 25^2}$
\nor
\n $p = -44178 \text{ N}$
\nAt this stage
\n $q_{\text{max}} = \sqrt{\frac{p_x - p_y}{2}}^2 + q^2$
\n $= \sqrt{\frac{-90 - 0}{2}}^2 + 0^2$
\n $= 45 \text{ N/mm}^2 > 25 \text{ N/mm}^2$.
\n \therefore Material fails because of maximum shear and not by axial compression.
\n $q_{\text{max}} = \sqrt{\frac{p_x - p_y}{2}}^2 + q^2 = \sqrt{\frac{p_x - 0}{2}}^2 + 0^2$
\n $= \frac{P_x}{2}$
\n \therefore Corresponding axial load is given by
\n $50 = \frac{P}{(\pi/4) \times 25^2}$
\nor
\n $p = 24543 \text{ N}$
\nThe plane of q_{max} is at 45° to the plane of p_x . This failure surface is shown in Fig. 11.6

 $= 45$ N/mm² > 25 N/mm².

∴ Material fails because of maximum shear and not by axial compression.

or $P = 24543 \text{ N}$

The plane of q_{max} is at 45° to the plane of p_x . This failure surface is shown in Fig. 11.8.

Plane of p_1

 \overline{p}

P

 p_{x}

Example 11.3. The direct stresses at a point in the strained material are 120 N/mm² compressive and 80 N/mm² tensile as shown in Fig. 11.9. There is no shear stress. Find the normal and tangential stresses on the plane AC. Also find the resultant stress on AC.

Solution: The plane *AC* makes 30° (anticlockwise) to the plane of p_x (*y*-axis). Hence $\theta = 30^\circ$. $p_x = 80 \text{ N/mm}^2$ $p_y = -120 \text{ N/mm}^2$ $q = 0$

$$
p_n = \frac{p_x + p_y}{2} + \frac{p_x - p_y}{2} \cos 2\theta + q \sin 2\theta
$$

= $\frac{80 - 120}{2} + \frac{80 - (-120)}{2} \cos (2 \times 30) + 0$
= $- 20 + 100 \cos 60$
Thus
 $p_n = 30 \text{ N/mm}^2$

$$
p_n = 30 \text{ N/mm}^2
$$

$$
p_t = \frac{p_x - p_y}{2} \sin 2\theta - q \cos 2\theta
$$

$$
= \frac{80 - (-120)}{2} \sin (2 \times 30) - 0
$$

Thus *pt* $= 86.6$ N/mm². The resultant of p_n and p_t is given by

$$
p = \sqrt{p_n^2 + p_t^2} = \sqrt{30^2 + 86.6^2}
$$

$$
p = 91.65 \text{ N/mm}^2
$$

Angle made by the resultant stress with p_t is given by tan $\alpha = \frac{p_t}{p_t}$ *n t* $=\frac{30}{25}$ 86.6 $\alpha = 19.1^\circ$ as shown in Fig. 11.9(*b*).

Example 11.4. The state of stress at a point in a strained material is as shown in Fig. 11.10. Determine

- (*i*) *the direction of principal planes*
- (*ii*) *the magnitude of principal stresses and*
- (*iii*) *the magnitude of maximum shear stress.*

Indicate the direction of all the above by a sketch.

Solution: $p_x = 200 \text{ N/mm}^2$ $p_y = 150 \text{ N/mm}^2$ $q = 100 \text{ N/mm}^2$ Let the principal plane make anticlockwise angle θ with the plane of p_x *i.e.* with *y*-axis. Then

(iii) the magnitude of maximum shear stress.
\nIndicate the direction of all the above by a sketch.
\n150 N/mm²
\n100 N/mm²
\n100 N/mm²
\n200 N/mm²
\n100 N/mm²
\n200 N/mm²
\nLet the principal plane make anticlockwise angle
$$
\theta
$$
 with the plane of p_x i.e. with y-
\ntan $2\theta = \frac{2q}{p_x - p_y} = \frac{2 \times 100}{200 - 150} = 4$
\n \therefore
\n20 = 75.96 and 75.96 + 180
\n20 = 75.96 and 75.96 + 180
\n20 = 75.96 and 75.96 + 180
\n20 = 37.98° and 127.98°
\n $p_1 = \frac{p_x + p_y}{2} + \sqrt{\left(\frac{p_x - p_y}{2}\right)^2 + q^2}$
\n $= \frac{200 + 150}{2} + \sqrt{\left(\frac{200 - 150}{2}\right)^2 + 100^2}$
\n $= 175 + 103.08$
\n $p_1 = 278.08$ N/mm²
\nand
\n $p_2 = \frac{p_x + p_y}{2} - \sqrt{\left(\frac{p_x - p_y}{2}\right)^2 + q^2}$
\n $p_2 = 175 - 103.08 = 71.92$ N/mm².
\n $q_{\text{max}} = \sqrt{\left(\frac{p_x - p_y}{2}\right)^2 + q^2}$
\ni.e.,
\n $q_{\text{max}} = 103.08$ N/mm²

$$
p_2 = 175 - 103.08 = 71.92 \text{ N/mm}^2.
$$

$$
q_{\max} = \sqrt{\left(\frac{p_x - p_y}{2}\right)^2 + q^2}
$$

i.e., $q_{\text{max}} = 103.08 \text{ N/mm}^2$

The planes of maximum shear are at 45° to the principal planes. Principal planes and planes of maximum shear are shown in Fig. 11.11.

x q_{max} plane q_{max} plane p_1 -plane q 45° 45° 45° p_2 - plane **Fig. 11.11** Example 11.5. The state of stress in a material stressed to 1970-there is
 P_{12} = plane

Example 11.5. The state of stress in a material stressed to 1970-dimensional state of stress

which they act.

Solution:
 $P_{12} =$

Example 11.5. The state of stress in a material stressed to two-dimensional state of stress is as shown in Fig. 11.12. Determine principal stresses and maximum shear stress and the planes on which they act.

Solution:

$$
p_{1, 2} = \frac{p_x + p_y}{2} \pm \sqrt{\left(\frac{p_x - p_y}{2}\right)^2 + q^2}
$$

In this problem,

$$
p_x = 80 \text{ N/mm}^2 \qquad p_y = -60 \text{ N/mm}^2 \qquad q = 20 \text{ N/mm}^2.
$$

$$
p_{1, 2} = \frac{80 + (-60)}{2} \pm \sqrt{\left(\frac{80 - (-60)}{2}\right)^2 + 20^2}
$$

$$
= 10 \pm \sqrt{70^2 + 20^2}
$$

$$
= 10 \pm 72.8
$$

$$
p_1 = 82.8 \text{ N/mm}^2
$$

and $p_2 = -62.8 \text{ N/mm}^2$

$$
q_{\text{max}} = \sqrt{\left(\frac{p_x - p_y}{2}\right)^2 + q^2}
$$

$$
= 72.8 \text{ N/mm}^2
$$

Let θ be the inclination of principal stress to the plane of p_x . Then, \odot

$$
\tan 2\theta = \frac{2q}{p_x - p_y} = \frac{2 \times 20}{80 - 60} = 2
$$
\n
$$
\therefore \quad 2\theta = 63.44^{\circ} \quad \text{or} \quad 63.44 + 180
$$
\n
$$
\therefore \quad \theta = 31.72^{\circ} \quad \text{or} \quad 121.72^{\circ}
$$

Planes of maximum shear make 45° to the above planes

$$
\therefore \qquad \theta' = 15.86^{\circ} \quad \text{and} \quad 60.86^{\circ}
$$

Example 11.6. The state of stress in two-dimensionally stressed body at a point is as shown in Fig. 11.13(a). Determine the principal planes, principal stresses, maximum shear stress and their planes

Fig. 11.13(a)

Solution: Let *x* and *y* directions be selected as shown in the figure. Then $p_x = -100$ N/mm², $p_y = -75$ N/mm²,

 $q = -50$ N/mm²

$$
p_1 = \frac{p_x + p_y}{2} + \sqrt{\left(\frac{p_x - p_y}{2}\right)^2 + q^2}
$$

= $\frac{-100 - 75}{2} + \sqrt{\left(\frac{-100 + 75}{2}\right)^2 + (-50)^2}$
= $-87.5 + 51.54$
 $p_1 = -35.96 \text{ N/mm}^2$
 $p_2 = -\frac{p_x + p_y}{2} - \sqrt{\left(\frac{p_x - p_y}{2}\right)^2 + q^2}$

i.e.,
$$
p_2 = -87.5 - 51.54
$$

\n $p_2 = -139.04 \text{ N/mm}^2$
\n $q_{\text{max}} = \sqrt{\left(\frac{p_x - p_y}{2}\right)^2 + q^2}$
\ni.e., $q_{\text{max}} = 51.51$.

Let principal plane of p_1 make angle θ with *x*-axis. Then

$$
\tan 2\theta = \frac{2q}{p_x - p_y} = \frac{2(-50)}{-100 + 75} = 4
$$

\n
$$
\therefore 2\theta = 75.96 \text{ and } 75.96 + 180
$$

\nor $\theta = 37.98^\circ \text{ and } 127.98^\circ.$

The planes of maximum shear stresses are at 45° to the principal planes. These planes are shown in Fig. 11.13(*b*).

Example 11.7. State of stress at a point in a material is as shown in the Fig. 11.14(a). Determine

- (*i*) *principal stresses*
- (*ii*) *maximum shear stress*
- (*iii*) *plane of maximum shear stress and*
- (*iv*) *the resultant stress on the plane of maximum shear stress.*

Solution: Selecting x and y-axis as shown in figure,
\n
$$
p_x = -50 \text{ N/mm}^2
$$
, $p_y = 100 \text{ N/mm}^2$, and $q = 75 \text{ N/mm}^2$.
\n
$$
\therefore p_1 = \frac{p_x + p_y}{2} + \sqrt{\frac{p_x - p_y}{2}}^2 + q^2
$$
\n
$$
= \frac{-50 + 100}{2} + \sqrt{\frac{-50 - 100}{2}}^2 + 75^2
$$
\n
$$
= 25 + 106.07
$$
\n
$$
= 131.07 \text{ N/mm}^2
$$
\n
$$
p_2 = \frac{p_x + p_y}{2} - \sqrt{\frac{p_x - p_y}{2}}^2 + q^2
$$
\n
$$
= 25 - 106.07
$$
\n
$$
= -81.07 \text{ N/mm}^2
$$
\n
$$
q_{\text{max}} = \sqrt{\frac{p_x - p_y}{2}}^2 + q^2
$$
\n
$$
= 106.07 \text{ N/mm}^2
$$
\nThe principal plane makes an angle θ to y-axis in anticlockwise direction. Then
\n
$$
\tan 2\theta = \frac{2q}{p_x - p_y} = \frac{2 \times 75}{\left(\frac{-50 - 100}{2}\right)} = -2
$$
\n
$$
\therefore \quad 2\theta = -63.43^{\circ}
$$
\nor
\n
$$
\theta = -31.72^{\circ} = 31.72^{\circ} \text{ clockwise.}
$$
\n
$$
\text{Plane of maximum shear mass of } 50^{\circ}
$$
\n
$$
\theta = -31.72^{\circ} + 45.00 = 13.28^{\circ}.
$$
\n
$$
\text{Normal stress on this plane is given by}
$$
\n
$$
p_x = \frac{p_x + p_y}{2} + \frac{p_x - p_y}{2} \cos 2\theta + q \sin 2\theta
$$
\n
$$
= \frac{-50 + 100}{2} + \frac{-50 - 100}{2} \cos 2(13.28) + 75 \sin (2 \times 13.28)
$$
\n
$$
= 25 - 67.08 + 33.54
$$
\n
$$
=
$$

or $θ = -31.72° = 31.72°$ clockwise.

Plane of maximum shear makes 45° to it

$$
\theta = -31.72 + 45.00 = 13.28^{\circ}.
$$

Normal stress on this plane is given by

∴ $2\theta = -63.43^\circ$

$$
p_x = \frac{p_x + p_y}{2} + \frac{p_x - p_y}{2} \cos 2\theta + q \sin 2\theta
$$

=
$$
\frac{-50 + 100}{2} + \frac{-50 - 100}{2} \cos 2(13.28) + 75 \sin (2 \times 13.28)
$$

=
$$
25 - 67.08 + 33.54
$$

=
$$
- 8.54 \text{ N/mm}^2
$$

$$
p_t = q_{\text{max}} = 106.07 \text{ N/mm}^2.
$$

Resultant stress

$$
p = \sqrt{(-8.54)^2 + 106.07^2}
$$

=
$$
106.41 \text{ N/mm}^2
$$

Let ' p ' make angle ϕ to tangential stress (maximum shear stress plane). Then referring to Fig. 11.14(*b*)

$$
\tan \phi = \frac{p_n}{p_t} = \frac{8.54}{106.07}
$$

 $\therefore \qquad \phi = 4.6^\circ \text{ as shown in Fig. 11.14}(b).$

Example 11.8. Show that when a material is subjected to shearing stresses and unidirectional direct stress, the major and minor principal stresses are of opposite nature.

Solution: Let p_x be uniaxial stress

$$
i.e.
$$

 $p_y = 0$

Let *q* be the shearing stresses.

Then from the equation

∴
\n
$$
\phi = 4.0
$$
 (b) 0.01
\nExample 11.8. Show that when a material is subjected to shearing
\nstresses and unidirectional direct stress, the major and minor
\nprincipal stresses are of opposite nature.
\nSolution: Let p_x be uniaxial stress.
\ni.e. p_y = 0
\nLet q be the shearing stresses.
\nThen from the equation
\n
$$
p_1 = \frac{p_x + p_y}{2} + \sqrt{\left(\frac{p_x - p_y}{2}\right)^2 + q^2}
$$
\nand
\nand
\n
$$
p_2 = \frac{p_x}{2} + \sqrt{\left(\frac{p_x}{2}\right)^2 + q^2}
$$
\nSince the second term is larger than the first term, naturally p₁ is +ve and p₂ is -ve. T
\nand minor principal stresses are having opposite nature.
\n11.3 PRINCIPAL STRESEES IN BEAMS
\nFigure 11.15 shows an element of a beam subjected to bending moment and shear force.
\nthe bending stress at point A, in the beam,
\nbending stress
\n
$$
p_x = \frac{M}{I}y
$$
\nand shear stress q = $\frac{F}{I}(a\bar{y})$
\nwhere M–Bending moment at the section
\nF–Shear force
\n
$$
I
$$
-Moment of area above the level of fibre at A about neutral axis.
\n
$$
b
$$
-width
\nand
\n $a\bar{y}$ -moment of area above the level of the path. The nature of bending stress
\nprinciple stress at point A is as shown in the Fig. 11.16 knowing p_x and q the
\nprincipal stress, maximum shear stress etc. may be found. The nature of bending stress
\ncompressive) should be carefully noted.

$$
Fig. 11.14(b)
$$

Since the second term is larger than the first term, naturally p_1 is +ve and p_2 is -ve. Thus major and minor principal stresses are having opposite nature.

11.3 PRINCIPAL STRESSES IN BEAMS

Figure 11.15 shows an element of a beam subjected to bending moment and shear force. We know, the bending stress at point *A*, in the beam,

 $bending$ stress

$$
v_x = \frac{1}{2}
$$

 $\frac{y}{I}$ *y*

and shear stress $q = \frac{F}{bl}(a\bar{y})$

where *M*—Bending moment at the section

F—Shear force

I—Moment of inertia

y—Distance of the point from neutral axis

b—width

and $a\bar{y}$ —moment of area above the level of fibre at *A* about neutral axis.

The state of stress at point *A* is as shown in the Fig. 11.16 knowing p_x and *q* the required principal stresses, maximum shear stress etc. may be found. The nature of bending stress (tensile or compressive) should be carefully noted.

Example 11.9. A shear force of 100 kN and a sagging moment of 80 kN-m act at a certain cross-section of rectangular beam 100 mm wide and 200 mm deep. Compute the principal stresses at a point 30 mm below the top surface.

Solution: Referring to Fig. 11.17,

$$
I = \frac{1}{12} \times 100 \times 200^3 = 66.667 \times 10^6
$$
 mm⁴

At point *A,* which is at 30 mm below top fibre

y = 100 - 30 = 70 mm
\n
$$
f_x = \frac{M}{I}y = \frac{80 \times 10^6}{66.667 \times 10^6} \times 70 = 84 \text{ N/mm}^2 \text{ (compressive)}
$$
\n
$$
p_x = -84 \text{ N/mm}^2
$$
\n
$$
p_x = -84 \text{ N/mm}^2
$$
\n
$$
p_x = \frac{100 \text{ mm}}{4} = \frac{1}{4} \text{ m}^2
$$
\n
$$
N = \frac{1}{
$$

Fig. 11.17

Shearing stress
\n
$$
q = \frac{F}{bt}(a\overline{y})
$$
\n
$$
= \frac{100 \times 10^3}{100 \times 66.667 \times 10^6} [100 \times 30 \times (100 - 15)]
$$
\nThus,
\n
$$
p_x = -84 \text{ N/mm}^2, p_y = 0, q = 3.82 \text{ N/mm}^2
$$
\n
$$
\therefore p_{1,2} = \frac{-84 + 0}{2} \pm \sqrt{\left(\frac{-84 + 0}{2}\right)^2 + (3.82)^2}
$$
\n
$$
= -42 \pm 42.17
$$
\n
$$
\therefore p_1 = 0.17 \text{ N/mm}^2
$$
\n
$$
p_2 = -84.17 \text{ N/mm}^2
$$
\nExample 11.10. A simply supported beam of 4 m span carries loads as shown in Fig. 12005. Section 6, the beam is 100 mm above the median is 100 mm above the mean is 100 mm. Use 120 kN = 20 kN

Example 11.10. A simply supported beam of 4 m span carries loads as shown in Fig. 11.18. The cross-section of the beam is 100 mm wide and 180 mm deep. At a section 1.5 m from left support, calculate the bending and shearing stresses at distances 0, 45 mm and 90 mm above the neutral axis. Find the principal planes and principal stresses at these points.

Solution: Due to symmetry,

$$
R_A = R_B = \frac{20 + 20 + 20}{2} = 30 \text{ kN}
$$

∴ At section 1.5 m from *A*,

$$
F = 30 - 20 = 10 \text{ kN}
$$

$$
M = 30 \times 1.5 - 20 \times 0.5 = 35 \text{ kN-m}
$$

$$
I = \frac{1}{12} \times 100 \times 180^3 = 48.6 \times 10^6 \text{ mm}^4
$$

g stress = $f = \frac{M}{12}y = \frac{35 \times 10^6}{100 \times 10^6}y$

$$
\therefore \quad \text{Bending stress} = f = \frac{M}{I}y = \frac{35 \times 10^6}{48.6 \times 10^6} y
$$

It varies linearly across the depth.

when $y = 0, f_1 = 0$

 $y = 45$ mm, $f_2 = 32.4$ N/mm² (comp). $y = 90$ mm, $f_3 = 64.8$ N/mm² (comp).

PRINCIPAL STRESSES AND STRAINS **363**

Shearing stress at a fibre '*y*' above *N*–*A* is

$$
q = \frac{F}{bI}(a\overline{y})
$$

$$
= \frac{10 \times 1000}{100 \times 48.6 \times 10^6} (a\overline{y})
$$

$$
\therefore \text{ At} \qquad y = 0, \quad q_1 = \frac{10 \times 1000}{100 \times 48.6 \times 10^6} \quad (100 \times 90 \times 45)
$$
\n
$$
= 0.833 \text{ N/mm}^2
$$
\nAt\n
$$
y = 45 \text{ mm}, \quad q_2 = \frac{10 \times 1000}{100 \times 48.6 \times 10^6} \quad 100 \times 45 \times \left(90 - \frac{45}{2}\right)
$$
\n
$$
= 0.625 \text{ N/mm}^2
$$
\nAt\n
$$
y = 90 \text{ mm}, \quad q_3 = 0
$$

×

The state of stresses on elements under considerations are as shown in Fig. 11.19.

$$
= \frac{10 \times 1000}{100 \times 48.6 \times 10^6} (a\bar{y})
$$

\n
$$
\therefore
$$
 At $y = 0$, $q_1 = \frac{10 \times 1000}{100 \times 48.6 \times 10^6} (100 \times 90 \times 45)$
\n
$$
= 0.833 \text{ N/mm}^2
$$

\nAt $y = 45 \text{ mm}$, $q_2 = \frac{10 \times 1000}{100 \times 48.6 \times 10^6} (100 \times 45 \times (90 - 100 \times 45 \times 100 \times 100 \times 45 \times 100 \times 10$

 $=$ – 16.2 \pm 16.212

∴ $p_1 = 0.012 \text{ N/mm}^2$

$$
p_2 = -32.412 \text{ N/mm}^2
$$

\ntan 2θ = $\frac{2 \times 0.625}{-32.4 - 0} = 0.09876$
\n2θ = 5.64° and 185.64°
\nθ = 2.82° and 92.82°
\n(c) At y = 90 mm,
\n $p_x = -64.8 \text{ N/mm}^2$, $p_y = 0$, $q = 0$
\n \therefore
\n $p_{1,2} = \frac{-64.8 + 0}{2} \pm \sqrt{\frac{(-64.8 - 0)}{2} + 0^2}$
\n= -32.4 ± 32.4
\n \therefore
\n $p_1 = 0 \text{ N/mm}^2$
\n $p_2 = -6.4.8 \text{ N/mm}^2$
\n $p_2 = 6.4.8 \text{ N mm}^2$
\n $p_2 = 6.4.8 \text{ N mm}^2$
\n $p_2 = 6.4.8 \text{ N mm}^2$
\n $p_2 = 6.4.$

Example 11.11. A simply supported beam of span 6 m has I-section as shown in Fig. 11.20(a). It carries uniformly distributed load (inclusive self weight) of 60 kN/m over entire span. Calculate the principal stresses and the maximum shearing stress at 100 mm above neutral axis of the beam at a section 1.5 m from support.

Fig. 11.20

Solution: $L = 6$ m, $w = 60$ kN/m

$$
\therefore \quad \text{Reaction at support} = \frac{6 \times 60}{2} = 180 \text{ kN}
$$

∴ Moment at 1.5 m from support

$$
M = 180 \times 1.5 - 60 \times \frac{1.5^2}{2} = 202.5 \text{ kN-m}
$$

Shear force at 1.5 m from support

$$
F = 180 - 1.5 \times 60 = 90 \text{ kN}
$$

Moment of inertia of the I-section,

$$
I = \frac{1}{12} \times 200 \times 400^3 - \frac{1}{12} \times 190 \times 380^3
$$

= 197.86 × 10⁶ mm⁴

∴ Bending stress at 100 mm above N–A

$$
f = \frac{M}{I}y = \frac{202.5 \times 10^6}{197.86 \times 10^6}
$$
 × 100 = 102.35 N/mm² (compressive)

Shear stress is given by

$$
q = \frac{F}{bI}(a\overline{y}) = \frac{90 \times 1000}{10 \times 197.86 \times 10^6} \times (200 \times 10 \times 195 + 10 \times 90 \times 145)
$$

= 23.68 N/mm²

Thus the state of stress on an element at $y = 100$ mm, as $p_x = f = -102.35$ N/mm², $p_y = 0$ $q = 23.68$ N/mm²

$$
p_{1, 2} = \frac{-102.35 + 0}{2} \pm \sqrt{\left(\frac{-102.35 - 0}{2}\right)^2 + (23.68)^2}
$$

= -51.175 ± 56.388

$$
p_1 = 5.21 \text{ N/mm}^2
$$

$$
q_{\text{max}} = \sqrt{\left(\frac{-102.35 - 0}{2}\right)^2 + (23.68)^2} = 56.39 \text{ N/mm}^2
$$

11.4 PRINCIPAL STRAINS

Let e_x be the strain in *x*-direction and e_y be the strain in *y*-direction where *x* and *y* are cartesian coordinate directions. Let shearing strain with respect to *x*-*y* coordinates be γ_{xy} . Then it is possible to find normal strain e_n , tangential strain e_t and shearing strain γ_{nt} on any plane inclined at θ to the plane of p_x . Then defining the plane with zero shearing strain as principal shearing plane, we can find the principal planes and principal strains.

Fig. 11.21

 B_1 $\overline{C_1}$

 B_1 C_2 C_3 C_4

 E_1

 θ

be $_{\mathsf{y}}$

A

xy

 $D \setminus D_1$ P

t

ae

Referring to Fig. $11.21(a)$, which represent state of stress at *A* with their positive senses it may be noted that direct strain e_x and e_y are tensile strains and shearing strain γ_{xy} reduces the angle θ .

n Figure 11.22 shows the element at *A* with its de-formed shape (shown with dotted lines) with e_x , *e_x* and γ_{xy} positive. Now we are interested in finding strains e_n , e_t and γ_{nt} on a plane at 'θ' to the plane of p_x forces. The size of element $(a \times b)$ is so selected that the diagonal AE is normal to the plane *DE i.e.* ∠*CAD* is also θ, for the convenience.

 (a) **To find e**_n

Drop $\perp C_1$ *, P* to *AD* \perp *BO* to C_1P . Now,

Let $AC = l$,

 $AD = l \cos \theta = a$

 $AB = l \sin \theta = b$

 $AP = AD + DD_1 + D_1P$ $= a + ae_x + b\gamma_{xy}$ $=$ *l* cos θ + *l* cos θ × e_x + *l* sin θ × γ_{xy} $=$ *l* [cos θ + e_x cos θ + γ_{xy} sin θ] *PC* = *PQ* + *QC* $= b + be_y$ $=$ *l* sin θ + *l* sin θ × e_y = *l* sin θ (1 + e_y) ∴ $AC_1 = \sqrt{AP^2 + PC^2}^{p_1}$ $= l \sqrt{(\cos \theta + e_x \cos \theta + \gamma_{xy} \sin \theta)^2 + \sin^2 \theta \times (1 + e_y)^2}$ **Fig. 11.22** e plane of p, b and p someoned that the material state in whom we have a matter of the set of the same of the state of define the convenience.

Let $AC = 1$ and $BC = 0$.
 $AD = 1 \cos \theta = a$
 $CD = 0$ and $CD = 1$ and $CD = 0$ b $AB = 1 \sin \$

Neglecting small quantities of higher order,

$$
AC_1 = l \sqrt{\cos^2 \theta + 2e_x \cos^2 \theta + 2\gamma_{xy} \sin \theta \cos \theta + \sin^2 \theta + 2e_y \sin^2 \theta}
$$

= $l \sqrt{1 + 2(e_x \cos^2 \theta + e_y \sin^2 \theta) + \gamma_{xy} \sin \theta \cos \theta}$

Expanding by the binomial theorem and neglecting small quantities of order 2 and above, we get $AC_1 = l [1 + e_x \cos^2 \theta + e_y \sin^2 \theta + \gamma_{xy} \sin \theta \cos \theta]$

$$
\therefore \qquad e_x = \frac{AC_1 - AC}{AC} = e_x \cos^2 \theta + e_y \sin^2 \theta + \gamma_{xy} \sin \theta \cos \theta \qquad \qquad \dots (11.9a)
$$

 $-\cos 2\theta$

But $\cos^2 \theta = \frac{1 + \cos 2\theta}{2}$

But
$$
\cos^2 \theta = \frac{1 + \cos 2\theta}{2}
$$
, $\sin^2 \theta = \frac{1 - \cos 2\theta}{2}$
and $\sin \theta \cos \theta = \frac{1}{2} \sin 2\theta$
 $e_x = \frac{e_x + e_y}{2} + \frac{e_x - e_y}{2} \cos 2\theta + \frac{1}{2} \gamma_{xy} \sin 2\theta$...(11.9b)

Strain
$$
e_t
$$
 can be found by replacing θ in the above equation by $90 + \theta$.

$$
\therefore \qquad e_t = \frac{e_x + e_y}{2} - \frac{e_x - e_y}{2} \cos 2\theta - \frac{1}{2} \gamma_{xy} \sin 2\theta
$$

To find γ*nt* :

After straining let point E_1 move to E_2 . Draw E_2R parallel to E_1 *D* [Ref. Fig. 11.23] ∴ Total shearing strain

$$
= \angle CAC_1 + \angle D_1E_2R = \phi_1 + \phi_2
$$

To find \upphi_1 : Now : $CC_3 = CC_2 \sin \theta$

$$
= (CQ – C2Q) sin θ
$$

=
$$
\left[(ae_x + bγ_{xy}) - be_y \frac{1}{tan θ} \right] sin θ
$$

=
$$
(ae_x + bγ_{xy}) sin θ - be_y cos θ
$$

Noting that $l \cos \theta = a$ and $l \sin \theta = b$, we get

$$
CC_3 = le_x \cos \theta \sin \theta + l \sin^2 \theta \gamma_{xy} - le_y \sin \theta \cos \theta
$$

= $l (e_x - e_y) \sin \theta \cos \theta + l \gamma_{xy} \sin^2 \theta$

$$
\therefore \phi_1 = \frac{CC_3}{l} = (e_x - e_y) \sin \theta \cos \theta + \gamma_{xy} \sin^2 \theta
$$
...(1)

To find ϕ_2 :

 E_3E_2 = Extension of AE_1 $= AE_1 \ e_n = a \cos \theta \ e_x = l \cos^2 \theta \times e_x$...(2) $\Phi_2 = \frac{RD_2}{E_1 D}$ $DD_2 - RD$ E_1D $DD_2 - E_3E$ E_1D \overline{c} 1 $2 - nD_2$ 1 $2 - L_3L_2$ 1 $=\frac{DD_2 - RD_2}{\sqrt{2}} = \frac{DD_2 - R}{\sqrt{2}}$ $=\frac{ae_x \cos \theta - l \cos^2 \theta e}{l}$ *a* χ cos θ – l cos² θ e_n sin $\theta - l \cos^2 \theta$ θ $-l \cos^2$ $=\frac{le_x \cos^2 \theta - le}{\sin \theta}$ *l* \int_{x} cos² θ – le_n cos $\sin \theta \cos$ 2 θ – le_{n} \cos^{2} θ θ cos θ − $= (e_x - e_n) \cot θ$...(3) To find \mathbf{y}_m :

To find \mathbf{y}_m

After straining tet point E, move to E, Daw E, $\mathbf{g}_{\mathbf{x}_k}$ best $\mathbf{g}_{\mathbf{x}_k}$.

After straining tet point E, move to E, Draw E, $\mathbf{g}_{\mathbf{x}_k}$ parallel to E, D Ref. Fig. 11.23

² ^γ*xy* sin 2^θ ...(11.10)

Substituting the value of
$$
e_n
$$
 from eqn. 11.9(*a*), we get
\n
$$
\Phi_2 = (e_x - e_y \cos^2 \theta - e_y \sin^2 \theta - \gamma_{xy} \sin \theta \cos \theta) \cot \theta
$$
\n
$$
= (e_x - e_y) \sin^2 \theta - e_y \sin^2 \theta - \gamma_{xy} \sin \theta \cos \theta \cot \theta
$$
\n
$$
= (e_x - e_y) \sin \theta \cos \theta + \gamma_{xy} \sin^2 \theta + (e_x - e_y) \sin \theta \cos \theta - \gamma_{xy} \cos^2 \theta
$$
\n
$$
\therefore \gamma_m = \Phi_1 + \Phi_2
$$
\n
$$
= (e_x - e_y) \sin \theta \cos \theta + \gamma_{xy} \sin^2 \theta + (e_x - e_y) \sin \theta \cos \theta - \gamma_{xy} \cos^2 \theta
$$
\n
$$
= 2(e_x - e_y) \sin \theta \cos \theta - \gamma_{xy} \cos 2\theta
$$
\n
$$
= (e_x - e_y) \sin 2\theta - \gamma_{xy} \cos 2\theta
$$
\n
$$
\therefore \frac{1}{2} \gamma_{\alpha} = \frac{e_x - e_y}{2} \sin 2\theta - \frac{1}{2} \gamma_{xy} \cos 2\theta
$$
\n
$$
\therefore \frac{1}{2} \gamma_{\alpha} = \frac{e_x - e_y}{2} \sin 2\theta - \frac{1}{2} \gamma_{xy} \cos 2\theta
$$
\n
$$
\therefore \text{ Defining principal strain as the par mth e plane where shearing strains is zero, the\ndirection of principal plane is obtained from 11.11 as\n
$$
\tan 2\theta = \frac{2(\frac{1}{2} \gamma_{xy})}{e_x - e_y} = \frac{1}{e_x - e_y}
$$
\n
$$
\therefore \text{Integrating principal strain so the plane where shearing strains is zero, the\ndirection of principal plane is obtained from 11.11 as\n
$$
\tan 2\theta = \frac{2(\frac{1}{2} \gamma_{xy})}{e_x - e_y} = \frac{1}{e_x - e_y}
$$
\n
$$
\therefore \frac{(11.12)}{2} \text{ The magnitude of principal planes may be obtained exactly on the same line as the principal
\nstress were obtained. The final result will be
\n
$$
e_{1
$$
$$
$$
$$

Note. Equation 11.1 is analogous to eqn. 11.9(*b*) and eqn. 11.2 is analogous to eqn. 11.11 in which p_1 or p_2 are replaced by e_1 and e_2 and q has been replaced by $\frac{1}{2} \gamma_{xy}$.

Defining **principal strain** as the *normal strains on the plane where shearing strains is zero*, the direction of principal plane is obtained from 11.11 as

$$
\tan 2\theta = \frac{2\left(\frac{1}{2}\gamma_{xy}\right)}{e_x - e_y} = \frac{\gamma_{xy}}{e_x - e_y} \tag{11.12}
$$

The magnitude of principal planes may be obtained exactly on the same line as the principal stresses were obtained. The final result will be

$$
e_{1, 2} = \frac{1}{2} (e_x + e_y) \pm \sqrt{\left(\frac{e_x - e_y}{2}\right)^2 + \left(\frac{1}{2} \gamma_{xy}\right)^2}
$$

= $\frac{1}{2} (e_x + e_y) \pm \frac{1}{2} \sqrt{(e_x - e_y)^2 + \gamma_{xy}^2}$...(11.13)

The maximum shearing strain occurs at 45° to the plane of principal plane and its magnitude is given by

$$
\gamma_{\text{max}} = \frac{1}{2} \sqrt{(e_x - e_y)^2 + \gamma_{xy}^2}
$$
...(11.14)

11.5 MEASUREMENT OF STRAIN

Strain is a physical quantity while stress is a concept. Hence it is possible to measure the strains in an experimental investigation. For this electrical strain gauges are used. Electrical strain gauge consists of a set of fine wires fixed at a predefined angles and then glued to the experimental model. After loading the model is strained, resulting into changes in the length and diameter of the wire. Hence the resistance of the wire also changes. Electrical strain meters, which work on the principal of Wheatstone Bridge measure the change in resistance and converts it to strains and then displays the readings. The set of strain gauges fixed at a point in different direction is called 'strain rosettes'. Commonly used strain rosettes are 45° rosette and 120° rosette which are as shown in Fig. 11.24. 45° rosette is also known as rectangular rosette.

Treating one of the strain gauge direction, strains $e_{\theta1}$ and $e_{\theta2}$ of the other two gauges is known writing $e_{\theta1}$ and $e_{\theta2}$ in terms of e_x , e_y and γ_{xy} , it is possible to find e_x , e_y and γ_{xy} .

After finding e_x , e_y and γ_{xy} it is possible to find the corresponding stresses. We know,

 $Ee_1 = p_1 - \mu p_2$...(1)

$$
E e_2 = p_2 - \mu p_1 \tag{2}
$$

Multiplying Eqn. (2) with μ and adding it to eqn.1, we get,

$$
E(e_1 + \mu e_2) = p_1 - \mu^2 p_1 = p_1 (1 - \mu^2)
$$

$$
p_1 = \frac{E(e_1 + \mu e_2)}{1 - \mu^2}
$$
...(11.15*a*)

Similarly $p_2 = \frac{E(e_2 + \mu e_1)}{1 - \mu^2}$

Example 11.12. At a point strains measured with rectangular rosettes are $e_A = 600$ microns, $e_B =$ *300 microns and ec = – 200 microns. Determine the principal strains. Find principal stresses also taking* $E = 2 \times 10^5$ *and* $\mu = 0.3$.

+ −

 μ µ

Solution: Taking the direction of strain gauge as *x*-axis,

$$
e_x = 600
$$
 microns
\n $e_B = e_{45} = 300$ microns and $e_c = e_{90} = -200$ microns $= e_y$
\n
$$
300 = \frac{1}{2} (e_x + e_y) + \frac{1}{2} (e_x - e_y) \cos 2 \times 45 + \frac{1}{2} \gamma_{xy} \sin 2 \times 45
$$
\n
$$
= \frac{1}{2} (600 - 200) + \frac{1}{2} (600 + 200) \times 0 + \frac{1}{2} \gamma_{xy}
$$
\n
$$
\gamma_{xy} = 200
$$
 microns.

∴ Principal strains are

$$
e_{1,2} = \frac{1}{2} (e_x + e_y) \pm \sqrt{\left(\frac{e_x - e_y}{2}\right)^2 + \left(\frac{1}{2} \gamma_{xy}\right)^2}
$$

...(11.15*b*)

$$
= \frac{1}{2}(600 - 200) \pm \sqrt{\frac{600 + 200}{2}}^2 + (\frac{1}{2} \times 200)^2
$$

\n
$$
= 200 \pm 412.3
$$

\n∴ $e_1 = 612.3 \text{ micron} = 612.3 \times 10^{-6}$
\n[Note: 1 micron = 1 × 10⁻⁶]
\n∴ $p_1 = \frac{E(e_1 + \mu e_2)}{1 - \mu^2} = \frac{2 \times 10^5 (612.3 - 0.3 \times 212.3) \times 10^{-6}}{1 - 0.3^2}$
\ni.e., $p_1 = 120.57 \text{ N/mm}^2$.
\n $p_2 = \frac{E(e_2 + \mu e_1)}{1 - \mu^2} = \frac{2 \times 10^5 (-212.3 + 0.3 \times 612.3)}{1 - 0.3^2}$
\ni.e., $p_2 = -6.29 \text{ N/mm}^2$
\nExample 11.13. In an experimental investigation strains observed with a 120° rosette are, $p_2 = -600 \text{ microns and } e_{240} = 100 \text{ microns. Determine the principal stresses, if } E$
\nN/mm² and $\mu = 0.3$.
\nSolution:
\n $e_x = 800$
\n $-600 = e_{120} = \frac{e_x + e_y}{2} + \frac{e_x - e_y}{2}$ (0.5) - $\frac{0.866}{2}$ Y_{xy} sin 240
\ni.e., $-600 = \frac{e_x + e_y}{2} + \frac{e_x - e_y}{2}$ (0.5) - $\frac{0.866}{2}$ Y_{xy}
\n100 = $e_{240} = \frac{e_x + e_y}{2} - \frac{e_x - e_y}{2}$ (0.5) + $\frac{0.866}{2}$ Y_{xy} sin 480
\n∴ $100 = \frac{e_x + e_y}{2} - \frac{e_x - e_y}{2}$ (0.5) + $\frac{0.866}{2}$ Y_{xy}

Example 11.13. In an experimental investigation strains observed with a 120° rosette are, $e_0 = 800$ *microns, e₁₂₀ = –600 microns and e₂₄₀ = 100 microns. Determine the principal stresses, if* $E = 2 \times 10^5$ *N/mm*² and $\mu = 0.3$.

Solution:
\n
$$
e_x = 800
$$
\n
$$
600 = e_{120} = \frac{e_x + e_y}{2} + \frac{e_x - e_y}{2} \cos 240 + \frac{1}{2} \gamma_{xy} \sin 240
$$
\n
$$
i.e.,
$$
\n
$$
600 = \frac{e_x + e_y}{2} - \frac{e_x - e_y}{2} \quad (0.5) - \frac{0.866}{2} \gamma_{xy} \qquad ...(1)
$$
\n
$$
100 = e_{240} = \frac{e_x + e_y}{2} + \frac{e_x - e_y}{2} \cos 480 + \frac{1}{2} \gamma_{xy} \sin 480
$$
\n
$$
i. \qquad 100 = \frac{e_x + e_y}{2} - \frac{e_x - e_y}{2} \quad (0.5) + \frac{0.866}{2} \gamma_{xy} \qquad ...(2)
$$

2

i.e.,

Adding eqns. (1) and (2) , we get

$$
-500 = e_x + e_y - \frac{e_x - e_y}{2}
$$

i.e., $-1000 = e_x + 3e_y$ But $e_x = 800$ microns ∴ $3e_y = -1000 - 800 = -1800$ $e_y = -600$ microns Substituting the values of e_x and e_y in eqn. (2), we get

> $100 = \frac{800 - 600}{3}$ 2 $800 + 600$ 2 $\frac{-600}{2} - \frac{800 + 600}{2}$ (0.5) + $\frac{0.866}{2}$ $rac{0.866}{2}$ γ_{xy} γ_{xy} = 808.3 microns

Thus
$$
e_x = 800
$$
 microns, $e_y = -600$ microns and $\gamma_{xy} = 808.3$ microns
\n
$$
\therefore \qquad e_{1,2} = \frac{e_x + e_y}{2} \pm \sqrt{\left(\frac{e_x - e_y}{2}\right)^2 + \left(\frac{1}{2} \gamma_{xy}\right)^2}
$$
\n
$$
= \frac{800 - 600}{2} \pm \sqrt{\left(\frac{800 + 600}{2}\right)^2 + \left(\frac{1}{2} \times 808.3\right)^2}
$$
\n
$$
= 100 \pm 808.29
$$
\n
$$
e_1 = 908.29 \text{ microns} = 908.29 \times 10^{-6}
$$
\nand\n
$$
e_2 = -708.29 \text{ microns} = -708.29 \times 10^{-6}
$$
\n
$$
\therefore \qquad p_1 = \frac{E(e_1 + \mu e_2)}{1 - \mu^2} = \frac{2 \times 10^3 (908.29 - 0.3 \times 708.29) \times 10^{-6}}{1 - 0.3^2}
$$
\ni.e.,\n
$$
p_1 = 152.9 \text{ Nmm}^2
$$
\ni.e.,\n
$$
p_2 = \frac{E(e_1 + \mu e_1)}{1 - \mu^2} = \frac{2 \times 10^3 (-708.29 + 0.3 \times 908.29) \times 10^{-6}}{1 - 0.3^2}
$$
\ni.e.,\n
$$
p_2 = -95.8 \text{ N/mm}^2
$$
\n**IMPORTANT FORMULAR**\n1. $p_n = \frac{p_x + p_y}{2} + \frac{p_z - p_y}{2} \cos 2\theta + q \sin 2\theta$.
\n2. $p_i = \frac{p_x - p_y}{2} \sin 2\theta - q \cos 2\theta$.
\n3. Principal planes are given by
\n
$$
\tan 2\theta = \frac{2q}{p_x - p_y}.
$$
\n4. $p_{1,2} = \frac{p_x + p_y}{2} \pm \sqrt{\left(\frac{p_x - p_y}{2}\right)^2 + q^2}.$
\n5. Maximum shear stress = $\frac{p_1 - p_2}{2} = \$

8.
$$
e_n = \frac{e_x + e_y}{2} + \frac{e_x - e_y}{2} \cos 2\theta + \frac{1}{2} \gamma_{xy} \sin 2\theta
$$

\n $\frac{1}{2} r_m = \frac{(e_x - e_y)}{2} \sin 2\theta - \frac{1}{2\gamma_{xy}} \cos 2\theta$
\n $e_{1,2} = \frac{1}{2} (e_x + e_y) \pm \sqrt{\frac{(e_x - e_y)}{2}^2 + (\frac{1}{2} \gamma_{xy})^2}$
\n $\gamma_{max} = \frac{e_1 - e_2}{2} = \sqrt{\frac{(e_x - e_y)^2}{2} + (\frac{1}{2} \gamma_{xy})^2}$
\n**THEORY QUESTIONS**
\n1. Derive the expressions for normal and tangential stresses on a plane inclined at θ to the p_x force. Take a general two dimensional state of stress.
\n2. State the equations for normal and tangential stresses on an inclined plane, in an element was
\ngeneral two dimensional stress system. Derive the expressions for principal planes,
\nstresses and maximum shear stress.
\n3. Explain the terms principal stresses and principal strains.
\n1. A point in a strained material is subjected to tensile stresses $p_x = 180 \text{ N/mm}^2$ and $p_y = 12$
\nDetermine the intensities of normal, tangential and resultant stresses on a plane incli
\nanticlockwise to the axis of minor stress.
\n[**Ans.** $p_n = 165 \text{ N/mm}^2$, $p_y = 25.98 \text{ N/mm}^2$, $p = 167.03 \text{ N/mm}^2$, $\alpha = 111.05^\circ$ anticlock
\naxis of minor strings.
\n[**Ans.** $p_n = 165 \text{ N/mm}^2$, $p_y = 25.98 \text{ N/mm}^2$, $p = 167.03 \text{ N/mm}^2$, $\alpha = 111.05^\circ$ anticlock
\naxis of minor principal plane]
\n2. The state of stress at a point in a strained material is as shown in Fig. 11.25. Det
\nnormal, tangential and the resultant stress on plane *DE*. Determine the direction of result
\n θ
\n θ
\n θ
\n θ
\n θ
\n θ
\n $\$

1. Derive the expressions for normal and tangential stresses on a plane inclined at θ to the plane of p_x force. Take a general two dimensional state of stress.

THEORY QUESTIONS

- **2.** State the equations for normal and tangential stresses on an inclined plane, in an element under general two dimensional stress system. Derive the expressions for principal planes, principal stresses and maximum shear stress.
- **3.** Explain the terms principal stresses and principal strains.

PROBLEMS FOR EXERCISE

1. A point in a strained material is subjected to tensile stresses $p_x = 180 \text{ N/mm}^2$ and $p_y = 120 \text{ N/mm}^2$. Determine the intensities of normal, tangential and resultant stresses on a plane inclined at 30° anticlockwise to the axis of minor stress.

[**Note :** Axis of minor stress means the plane of major stress]

[Ans. $p_n = 165 \text{ N/mm}^2$, $p_t = 25.98 \text{ N/mm}^2$, $p = 167.03 \text{ N/mm}^2$, $\alpha = 111.05^{\circ}$ anticlockwise to the axis of minor principal plane]

2. The state of stress at a point in a strained material is as shown in Fig. 11.25. Determine the normal, tangential and the resultant stress on plane *DE*. Determine the direction of resultant also.

Fig. 11.25

[Ans. $p_n = 0$, $p_t = 69.28 \text{ N/mm}^2$, $p = 69.28 \text{ N/mm}^2$ acts tangential to *DE*.]

3. The state of stress at a point is as shown in Fig. 11.26. Determine the principal stresses and maximum shear stress. Indicate their planes on a separate sketch.

[Ans. $p_1 = 133.13 \text{ N/mm}^2$, $p_2 = -93.13 \text{ N/mm}^2$, $q_{\text{max}} = 113.13 \text{ N/mm}^2$, $\theta = -22.5^\circ$ and 67.5° , $\theta' = 22.5$ and 112.5°]

Fig. 11.26

4. The state of stress in a two dimensionally stressed material is as shown in Fig. 11.27. Determine the principal stresses, principal planes and the maximum shear stress. Determine normal and tangential stresses on plane *AC* also.

Fig. 11.27

[Ans. $p_1 = -9.18 \text{ N/mm}^2$, $p_2 = -130.82 \text{ N/mm}^2$, $q_{\text{max}} = 60.82 \text{ N/mm}^2$, $\theta = -40.27^{\circ}$ and 49.73°, $p_n = -116.96$ N/mm², $p_t = 21.34$ N/mm²]